

$\bar{K}N$ 相互作用および関連する 物理、その最近の進展



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Contents



$\Lambda(1405)$ and $\bar{K}N$ interactions

[Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 \(2011\); NPA 881, 98 \(2012\);](#)

[T. Hyodo, M. Niiyama, PPNP 120, 103868 \(2021\);](#)

[T. Hyodo, W. Weise, arXiv:2202.06181 \[nucl-th\] \(Handbook of Nuclear Physics\)](#)

- Recent developments

[J.-X. Lu, L.S. Geng, M. Doering, M. Mai, PRL 130, 071902 \(2023\);](#)

[J. Bulava, *et al.* \(BaSc\), arXiv:2307.10413 \[hep-lat\]; arXiv:2307.13471 \[hep-lat\]](#)



K^-p femtoscopy

[Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise. PRL 124, 132501 \(2020\)](#)

- Experimental data

[ALICE collaboration, PRL 124, 092301 \(2020\); PLB 822, 136708 \(2021\);](#)

[EPJC 83, 340 \(2023\)](#)

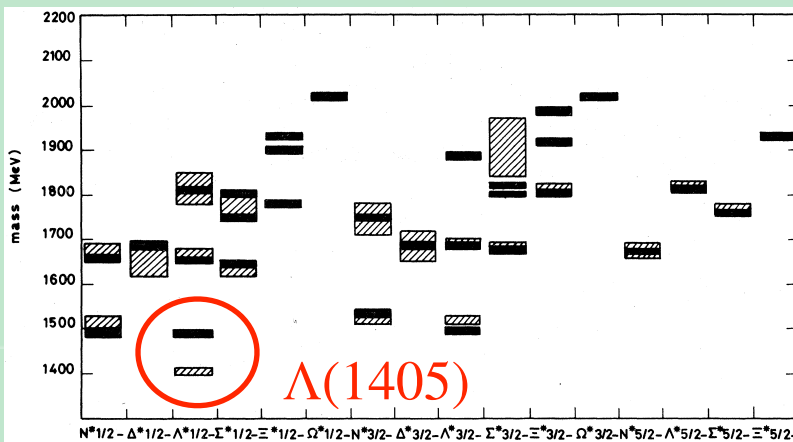
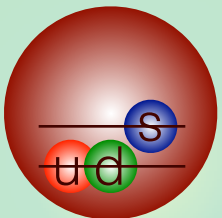


Summary

$\Lambda(1405)$ and $\bar{K}N$ scattering

$\Lambda(1405)$ does not fit in standard picture \rightarrow exotic candidate

N. Isgur and G. Karl, PRD 18, 4187 (1978)

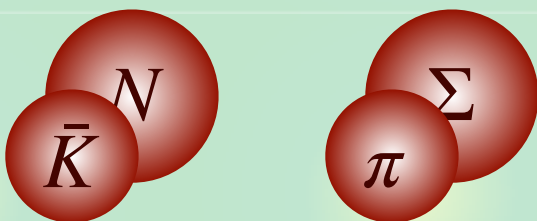


— : theory

▨ : experiment

Resonance in coupled-channel scattering

- Coupling to MB states



energy ↑

— $\bar{K}N$ threshold

▭ $\Lambda(1405)$

— $\pi\Sigma$ threshold

Detailed analysis of $\bar{K}N$ - $\pi\Sigma$ scattering is necessary

Strategy for $\bar{K}N$ interaction

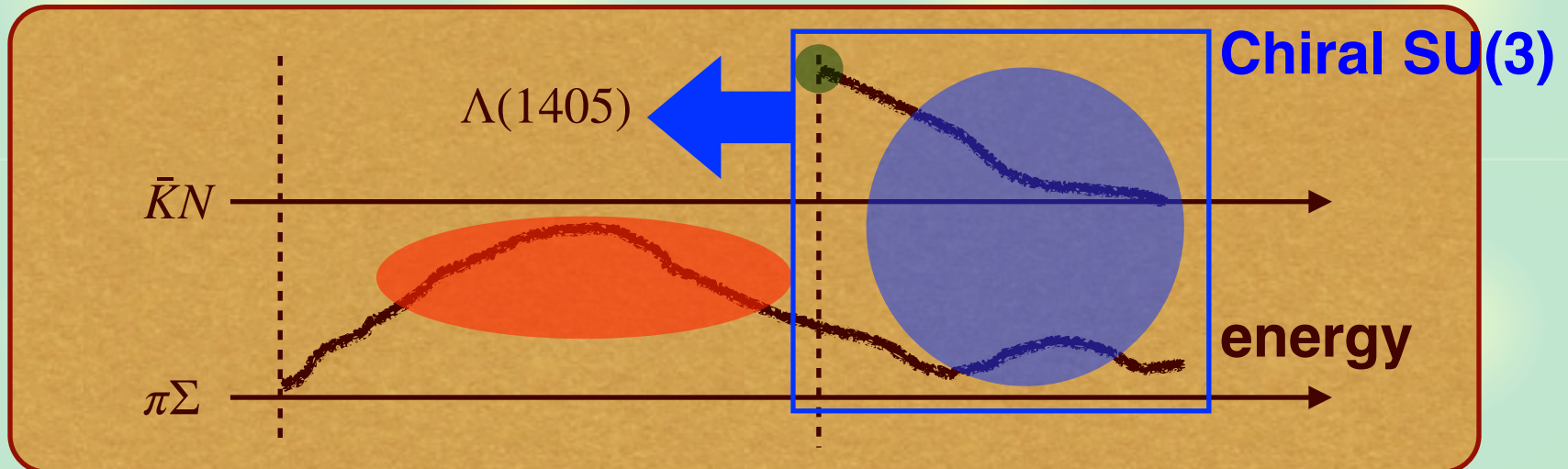
Above the $\bar{K}N$ threshold : direct constraints

- K^-p total cross sections (old data)
- $\bar{K}N$ threshold branching ratios (old data)
- K^-p scattering length (new data : SIDDHARTA)

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881, 98 (2012)

Below the $\bar{K}N$ threshold: indirect (reaction model needed)

- $\pi\Sigma$ mass spectra (LEPS, CLAS, HADES, J-PARC, ...)



Best-fit results by chiral SU(3) dynamics

K at rest

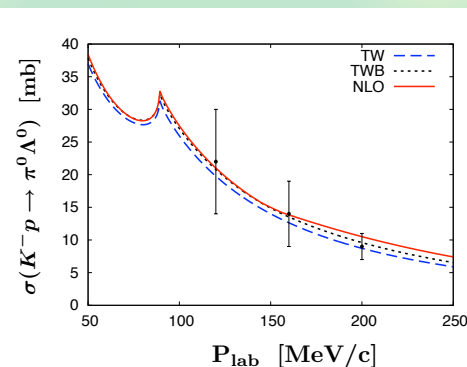
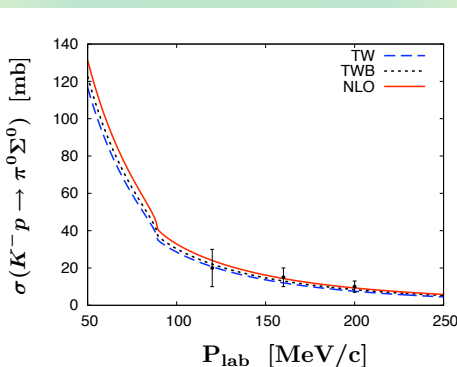
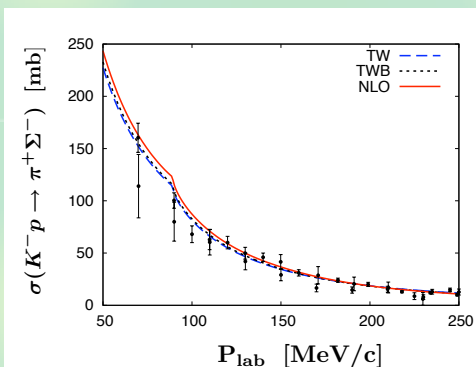
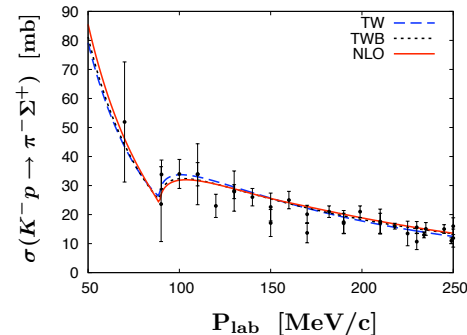
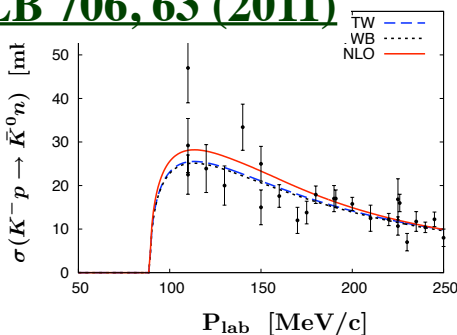
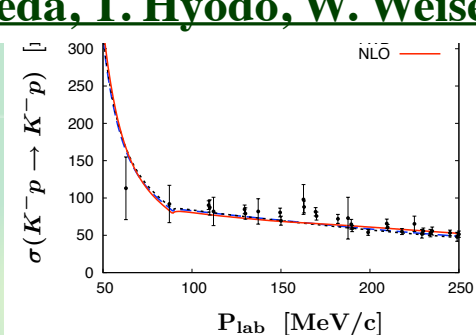
	TW	TWB	NLO	Experiment
ΔE [eV]	373	377	306	$283 \pm 36 \pm 6$ [10]
Γ [eV]	495	514	591	$541 \pm 89 \pm 22$ [10]
γ	2.36	2.36	2.37	2.36 ± 0.04 [11]
R_n	0.20	0.19	0.19	0.189 ± 0.015 [11]
R_c	0.66	0.66	0.66	0.664 ± 0.011 [11]
$\chi^2/\text{d.o.f}$	1.12	1.15	0.96	

} **SIDDHARTA**

} **Branching ratios**

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011)

K^-p cross sections



Accurate description of all existing data ($\chi^2/\text{d.o.f} \sim 1$)

PDG has changed

2020 update of PDG

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881, 98 (2012); ▲

Z.H. Guo, J.A. Oller, PRC 87, 035202 (2013); ✕

M. Mai, U.G. Meißner, EPJA 51, 30 (2015) ■ ○

- Particle Listing section:

Citation: P.A. Zyla *et al.* (Particle Data Group), Prog. Theor. Exp. Phys. **2020**, 083C01 (2020)

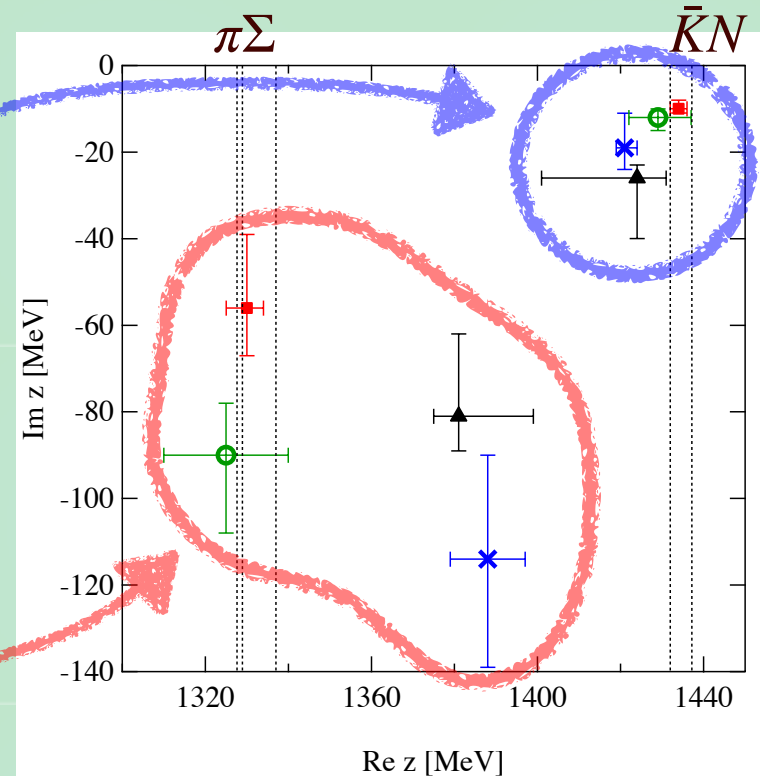
$\Lambda(1405) 1/2^-$

$I(J^P) = 0(\frac{1}{2}^-)$ Status: ****

Citation: P.A. Zyla *et al.* (Particle Data Group), Prog. Theor. Exp. Phys. **2020**, 083C01 (2020)

$\Lambda(1380) 1/2^-$

$J^P = \frac{1}{2}^-$ Status: **
new!



T. Hyodo, M. Niyama, PPNP 120, 103868 (2021)

- “ $\Lambda(1405)$ ” is no longer at 1405 MeV but ~ 1420 MeV.
- Lower pole : two-star resonance $\Lambda(1380)$

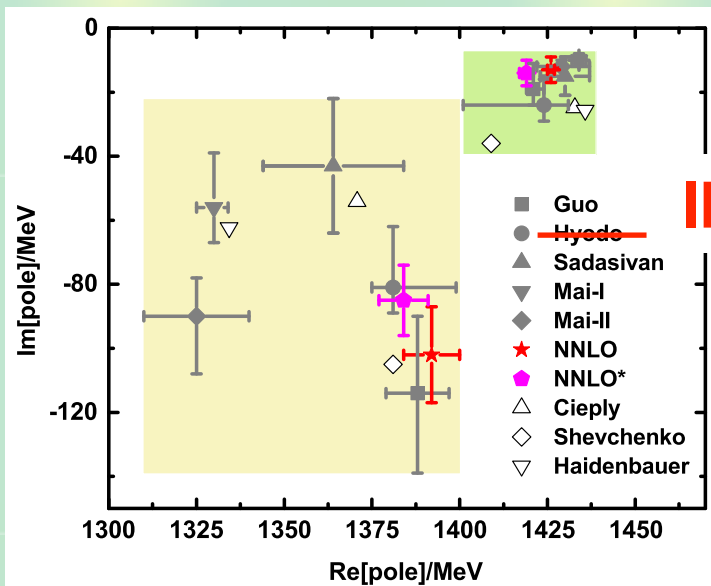
NNLO analysis and lattice QCD

Analysis at NNLO chiral SU(3) dynamics ($\bar{K}N$ and πN included)

J.-X. Lu, L.S. Geng, M. Doering, M. Mai, PRL 130, 071902 (2023)

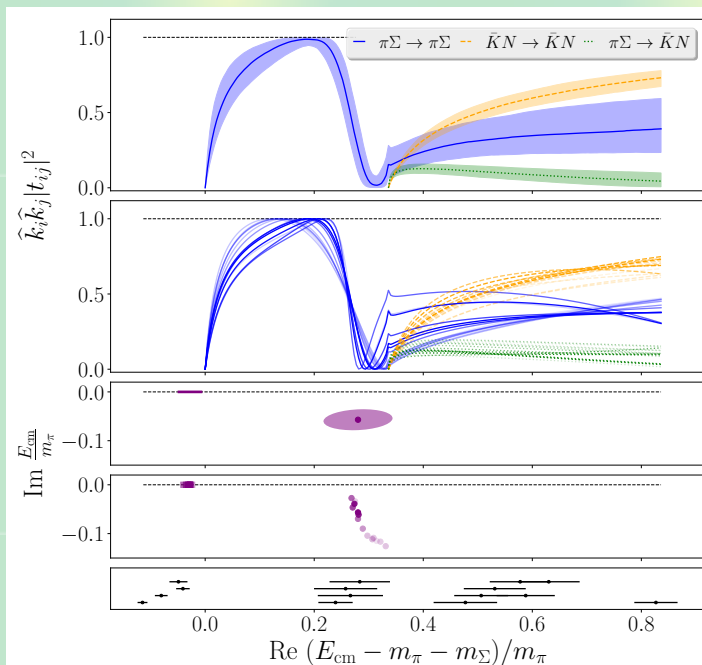
Lattice QCD calculation of $\bar{K}N$ - $\pi\Sigma$ scattering ($m_\pi \sim 200$ MeV)

J. Bulava, *et al.* (BaSc), arXiv:2307.10413 [hep-lat]; arXiv:2307.13471 [hep-lat]



Pole positions (MeV)

$\Lambda(1380)$	$1392 \pm 8 - i(102 \pm 15)$
$\Lambda(1405)$	$1425 \pm 1 - i(13 \pm 4)$



$$E_1 = 1392(9)(2)(16) \text{ MeV},$$

$$E_2 = [1455(13)(2)(17) - i11.5(4.4)(4)(0.1)] \text{ MeV},$$

Two states are confirmed at NNLO and lattice QCD

Construction of $\bar{K}N$ potentials

Local $\bar{K}N$ potential is useful for various applications

meson-baryon amplitude
(chiral SU(3) EFT)

T. Hyodo, W. Weise, PRC 77, 035204 (2008)

Kyoto $\bar{K}N$ potential
(single-channel, complex)

K. Miyahara, T. Hyodo,
PRC 93, 015201 (2016)

Kyoto $\bar{K}N-\pi\Sigma-\pi\Lambda$ potential
(coupled-channel, real)

K. Miyahara, T. Hyodo, W. Weise,
PRC 98, 025201 (2018)

Kaonic nuclei

Kaonic deuterium

K^-p correlation function



$\Lambda(1405)$ and $\bar{K}N$ interactions

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881, 98 (2012);

T. Hyodo, M. Niiyama, PPNP 120, 103868 (2021);

T. Hyodo, W. Weise, arXiv:2202.06181 [nucl-th] (Handbook of Nuclear Physics)

- Recent developments

J.-X. Lu, L.S. Geng, M. Doering, M. Mai, PRL 130, 071902 (2023);

J. Bulava, *et al.* (BaSc), arXiv:2307.10413 [hep-lat]; arXiv:2307.13471 [hep-lat]



K^-p femtoscopy

Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise. PRL 124, 132501 (2020)

- Experimental data

ALICE collaboration, PRL 124, 092301 (2020); PLB 822, 136708 (2021);

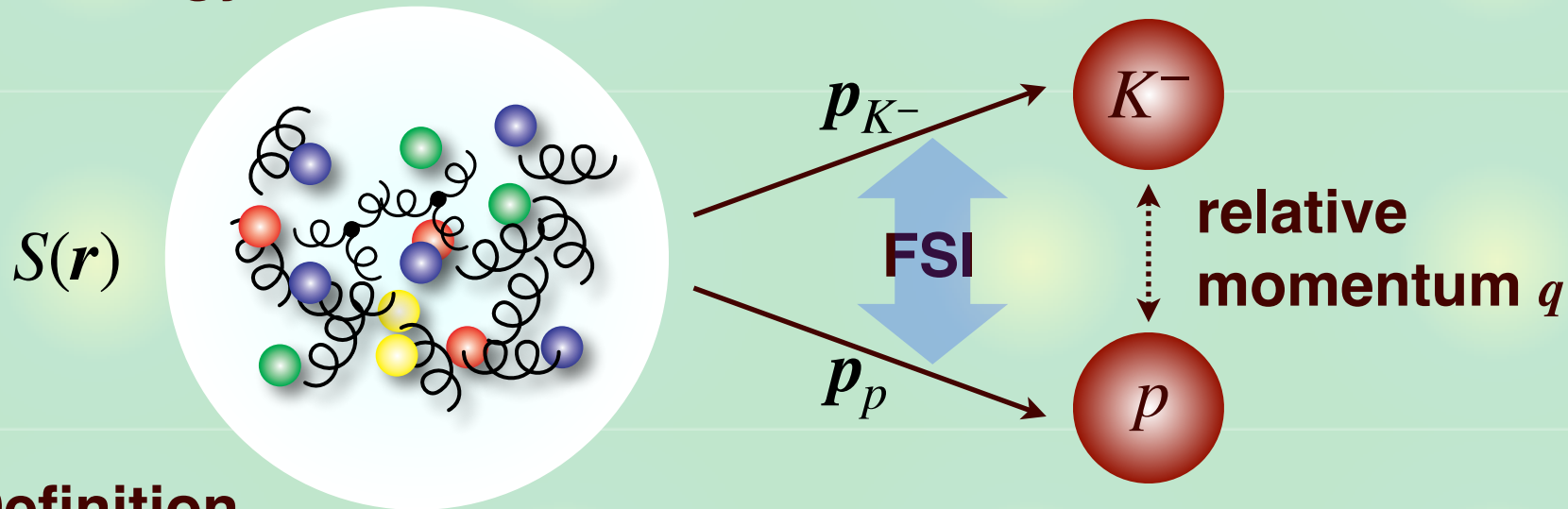
EPJC 83, 340 (2023)



Summary

Correlation function and hadron interaction

High-energy collision: chaotic source $S(\mathbf{r})$ of hadron emission



- Definition

$$C(q) = \frac{N_{K^-p}(\mathbf{p}_{K^-}, \mathbf{p}_p)}{N_{K^-}(\mathbf{p}_{K^-})N_p(\mathbf{p}_p)} \quad (= 1 \text{ in the absence of FSI/QS})$$

- Theory (Koonin-Pratt formula)

S.E. Koonin PLB 70, 43 (1977); S. Pratt, PRD 33, 1314 (1986)

$$C(q) \simeq \int d^3\mathbf{r} S(\mathbf{r}) |\Psi_q^{(-)}(\mathbf{r})|^2$$

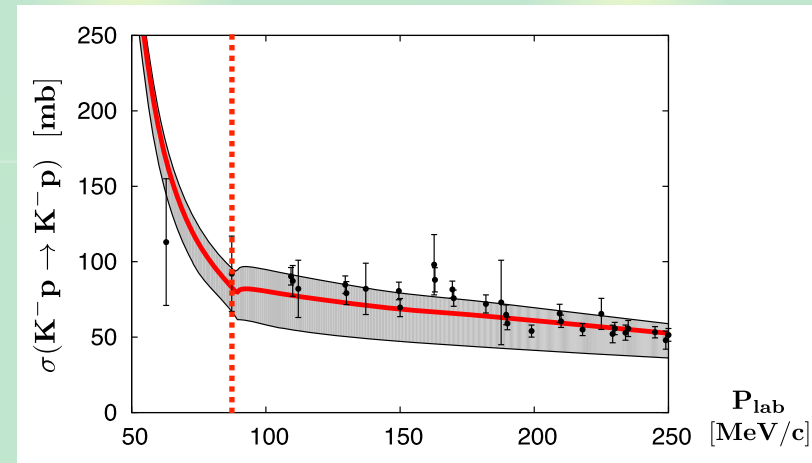
Source function $S(\mathbf{r}) \leftrightarrow$ wave function $\Psi_q^{(-)}(\mathbf{r})$ (FSI)

Experimental data of K^-p correlation

K^-p total cross sections

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011)

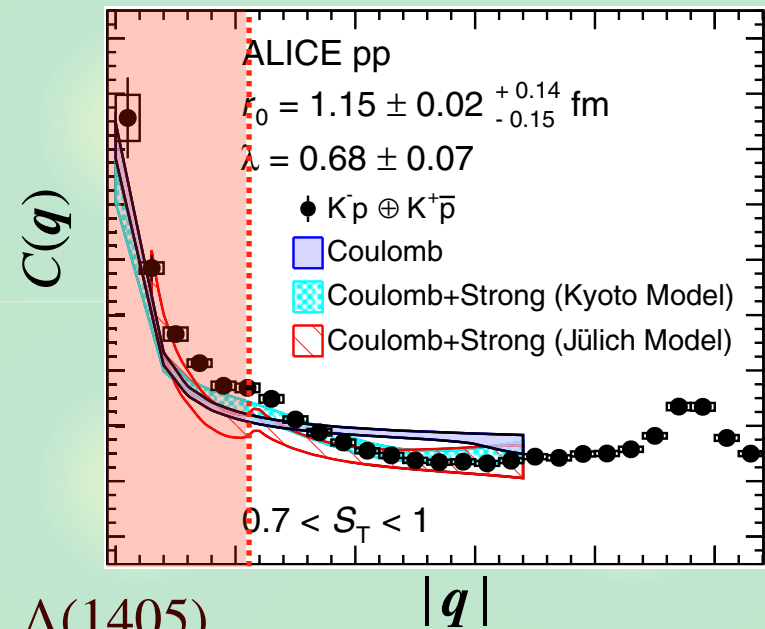
- Old bubble chamber data
- Resolution is not good
- Threshold cusp is not visible



K^-p correlation function

ALICE collaboration, PRL 124, 092301 (2020)

- Excellent **precision** (\bar{K}^0_n cusp)
- Low-energy data **below** \bar{K}^0_n



—> Important constraint on $\bar{K}N$ and $\Lambda(1405)$

Coupled-channel effects

Schrödinger equation (s-wave)

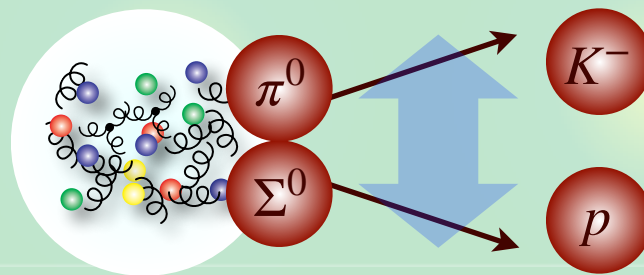
$$\begin{pmatrix} -\frac{\nabla^2}{2\mu_1} + V_{11}(r) + V_C(r) & V_{12}(r) & \dots \\ V_{21}(r) & -\frac{\nabla^2}{2\mu_2} + V_{22}(r) + \Delta_2 & \dots \\ \vdots & \vdots & \ddots \end{pmatrix} \begin{pmatrix} \psi_{K^-p}(r) \\ \psi_{\bar{K}^0n}(r) \\ \vdots \end{pmatrix} = E \begin{pmatrix} \psi_{K^-p}(r) \\ \psi_{\bar{K}^0n}(r) \\ \vdots \end{pmatrix}$$

Coulomb **threshold energy difference**

Asymptotic ($r \rightarrow \infty$) wave function

$$\begin{pmatrix} \psi_{K^-p}(r) \\ \psi_{\bar{K}^0n}(r) \\ \vdots \end{pmatrix} \propto \begin{pmatrix} \#e^{-iqr} + \#e^{iqr} \\ \#e^{-iq_2r} + \#e^{iq_2r} \\ \vdots \end{pmatrix} \quad \text{incoming + outgoing}$$

- **Transition** from $\bar{K}^0n, \pi^+\Sigma^-, \pi^0\Sigma^0, \pi^-\Sigma^+, \pi^0\Lambda$ is in $\psi_i(r)$ with $i \neq K^-p$



Coupled-channel correlation function

Coupled-channel Koonin-Pratt formula

R. Lednicky, V.V. Lyuboshitz, V.L. Lyuboshitz, Phys. Atom. Nucl. 61, 2950 (1998);

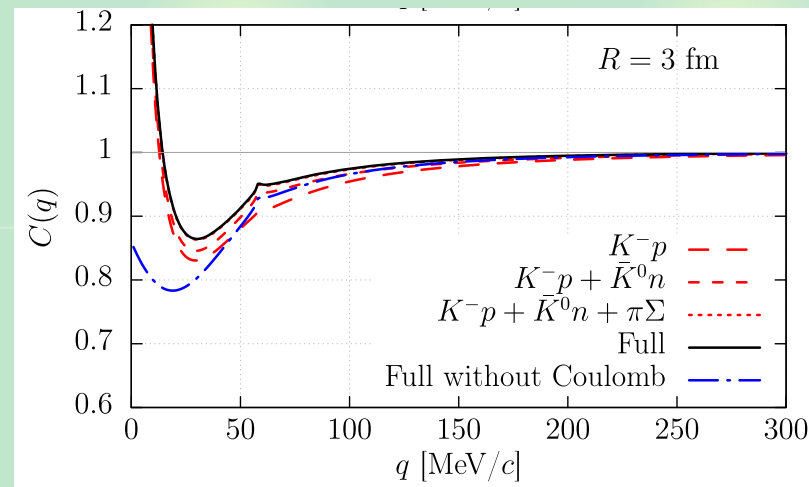
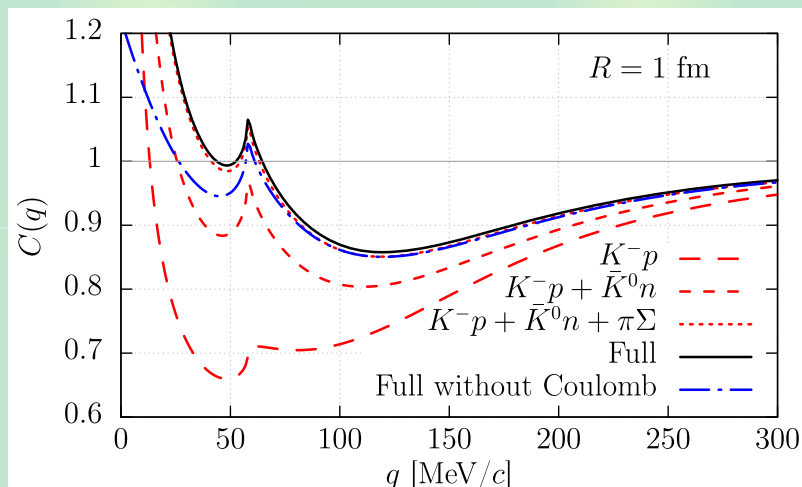
J. Haidenbauer, NPA 981, 1 (2019);

Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise, PRL124, 132501 (2020)

$$C_{K^-p}(q) \simeq \int d^3r S_{K^-p}(r) |\Psi_{K^-p,q}^{(-)}(r)|^2 + \sum_{i \neq K^-p} \omega_i \int d^3r S_i(r) |\Psi_{i,q}^{(-)}(r)|^2$$

- **Transition** from $\bar{K}^0 n, \pi^+ \Sigma^-, \pi^0 \Sigma^0, \pi^- \Sigma^+, \pi^0 \Lambda$

- ω_i : weight of source channel i relative to K^-p



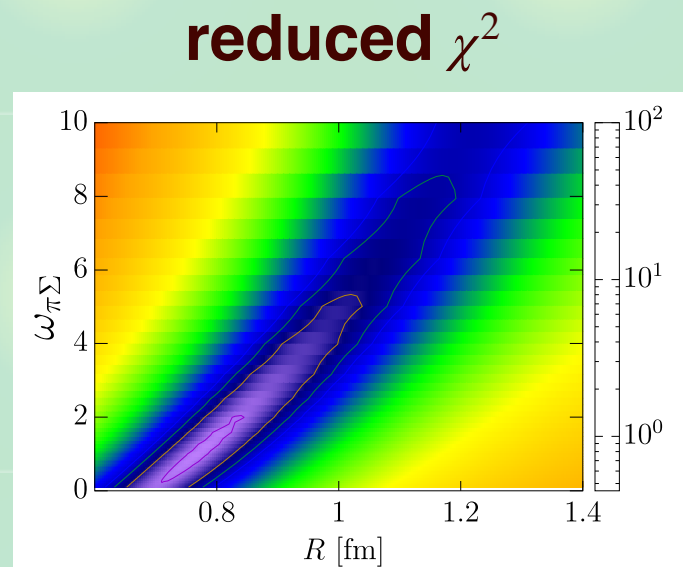
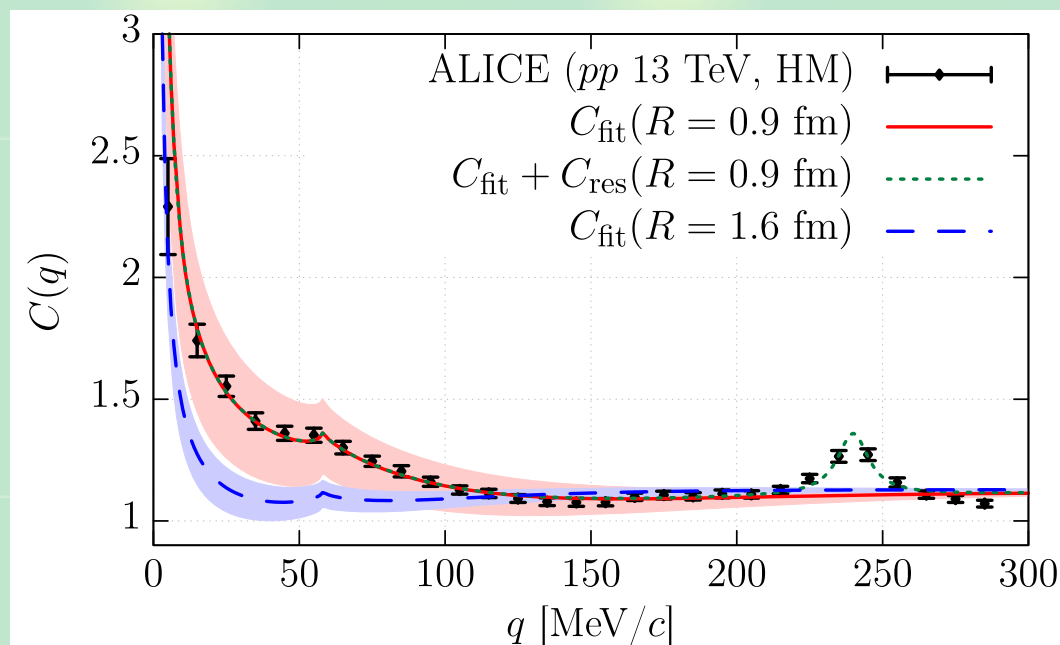
Coupled-channel effect is enhanced for small sources

Correlation from chiral SU(3) dynamics

Wave function $\Psi_{i,q}^{(-)}(r)$: coupled-channel $\bar{K}N$ - $\pi\Sigma$ - $\pi\Lambda$ potential

K. Miyahara, T. Hyodo, W. Weise, PRC 98, 025201 (2018)

- Source function $S(r)$: Gaussian, $R \sim 1$ fm in K^+p data
- Source weight $\omega_{\pi\Sigma} \sim 2$ by simple statistical model estimate



Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise, PRL 124, 132501 (2020)

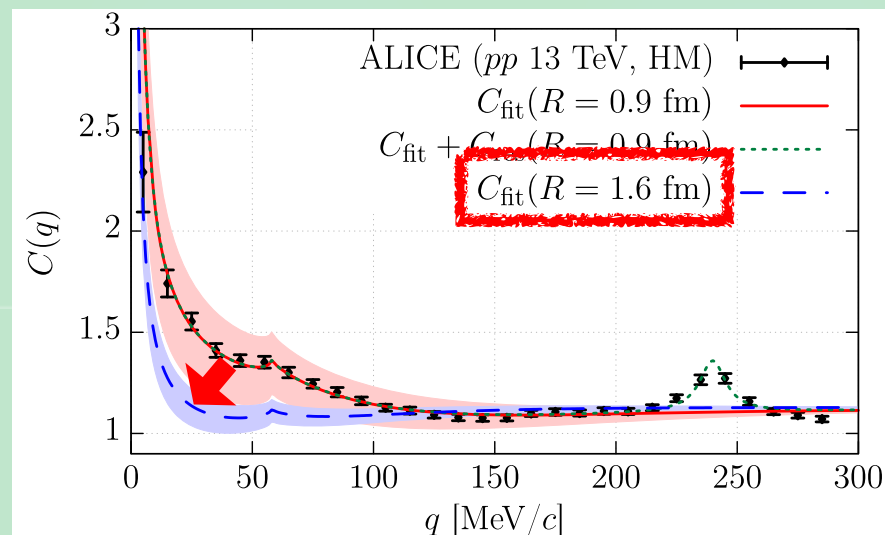
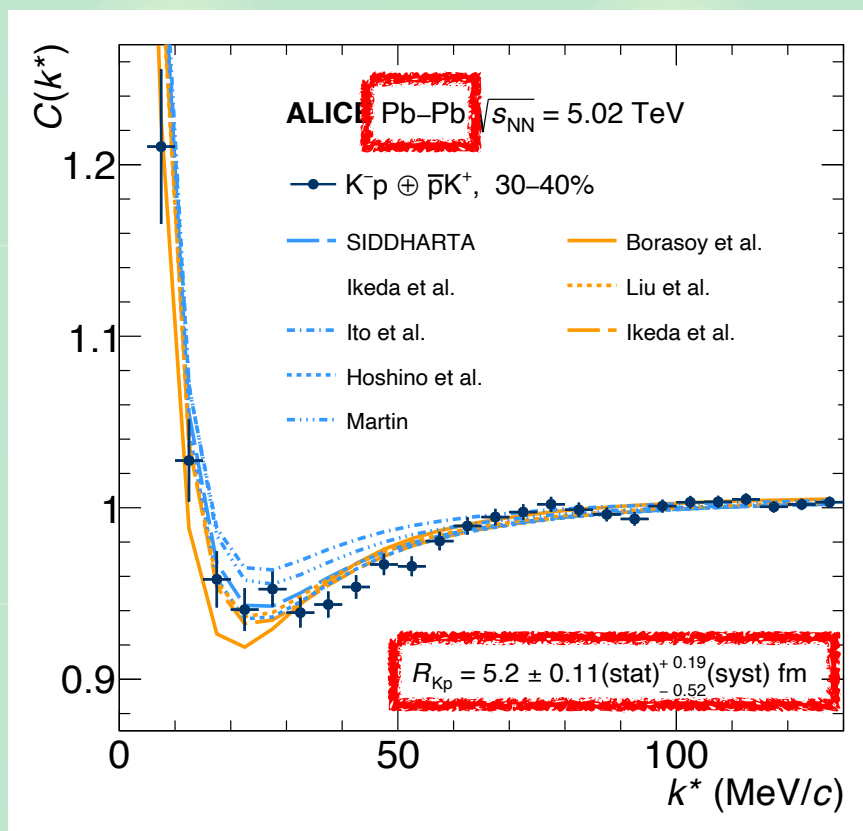
Correlation function by ALICE is well reproduced

Source size dependence

New data with Pb-Pb collisions at 5.02 TeV

ALICE collaboration, PLB 822, 136708 (2021)

- Scattering length $a_{K^-p} = -0.91 + 0.92i$ fm



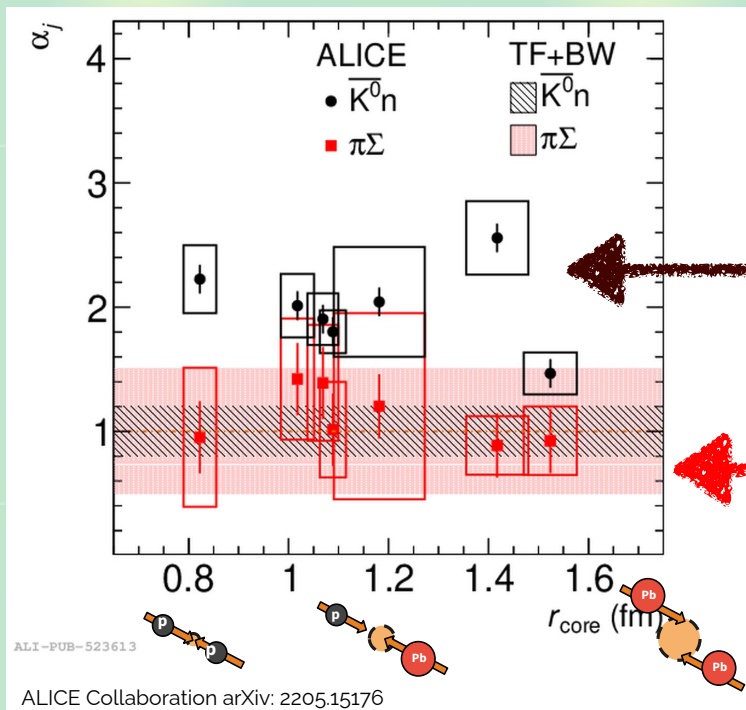
Correlation is suppressed at larger R , as predicted

Systematic study of source size dependence

Correlations in pp , p -Pb, Pb-Pb by Kyoto $\bar{K}N$ - $\pi\Sigma$ - $\pi\Lambda$ potential

ALICE collaboration, EPJC 83, 340 (2023)

$$C_{K^-p}(q) \simeq \int d^3r S_{K^-p}(r) |\Psi_{K^-p,q}^{(-)}(r)|^2 + \sum_{i \neq K^-p} \omega_i \int d^3r S_i(r) |\Psi_{i,q}^{(-)}(r)|^2$$



enhancement needed to explain data

Expected weight ω_i by Thermal Fist + Blast Wave

More strength is needed in the \bar{K}^0n channel

Summary



K^-p scattering and kaonic hydrogen are well described by NLO chiral SU(3) dynamics.

[Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 \(2011\); NPA 881, 98 \(2012\)](#)

- NNLO, scattering on the lattice, ...

[J.-X. Lu, L.S. Geng, M. Doering, M. Mai, PRL 130, 071902 \(2023\);](#)

[J. Bulava, *et al.* \(BaSc\), arXiv:2307.10413 \[hep-lat\]; arXiv:2307.13471 \[hep-lat\]](#)



Global structures of K^-p correlation functions are reproduced by Kyoto $\bar{K}N-\pi\Sigma-\pi\Lambda$ potential.

[Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise. PRL124, 132501 \(2020\)](#)

- Source size dependence

[ALICE collaboration, PRL 124, 092301 \(2020\); PLB 822, 136708 \(2021\);](#)

[EPJC 83, 340 \(2023\)](#)