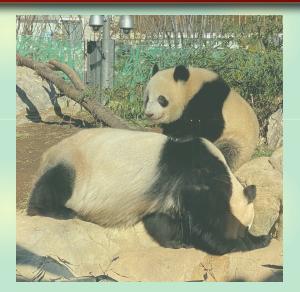
Femtoscopy for exotic hadrons and nuclei





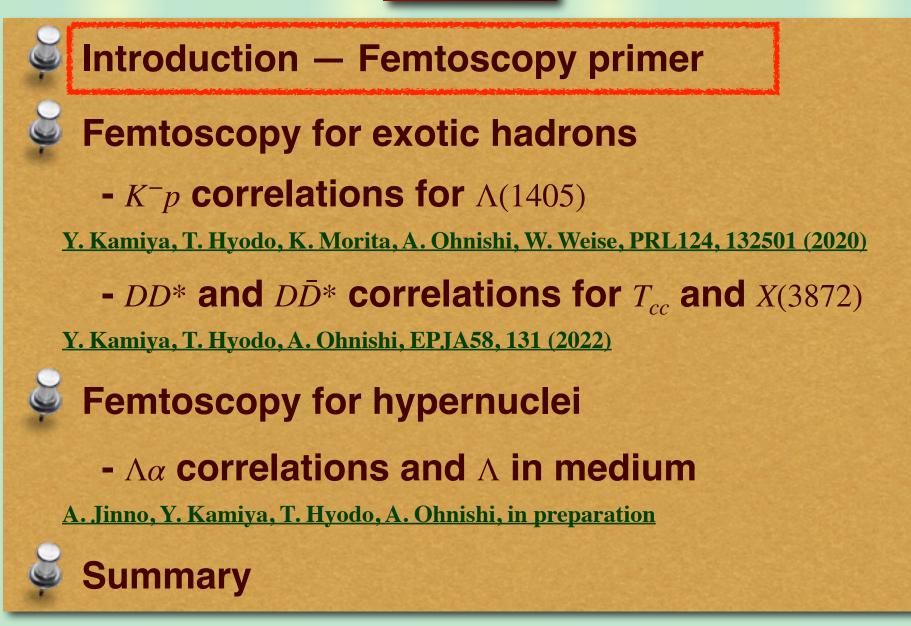
Tetsuo Hyodo

Tokyo Metropolitan Univ.



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Introduction — Femtoscopy primer

Scattering experiments and femtoscopy

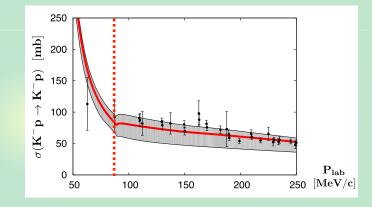
Traditional methods: scattering experiments

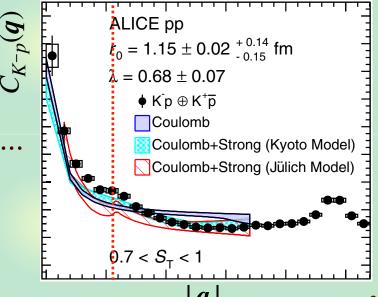
Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011)

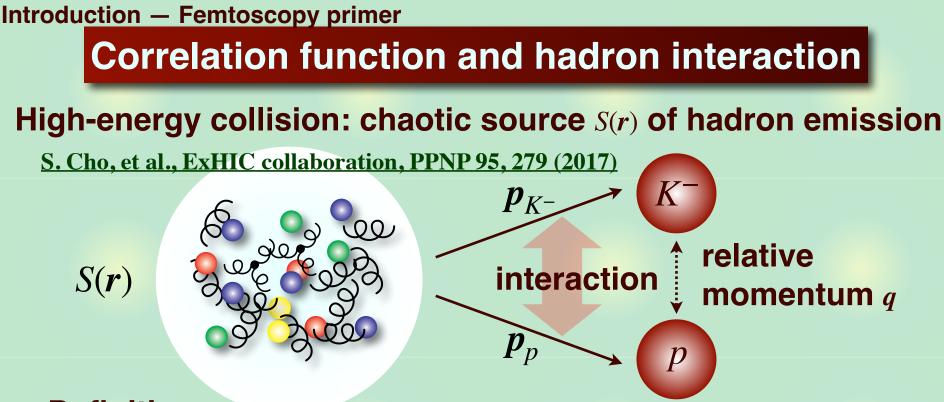
- limited channels: NN, YN, πN , KN, $\bar{K}N$, ...
- limited statistics (low-energy)
- heavy (c, b) hadrons: impossible

Femtoscopy: correlation function ALICE collaboration, PRL 124, 092301 (2020)

- various systems: $\Lambda\Lambda, N\Omega, \phi N, \overline{K}\Lambda, DN, \cdots$
- Excellent precision (\bar{K}^0n cusp)
- heavy hadrons: possible!







- Definition

$$C(\boldsymbol{q}) = \frac{N_{K^-p}(\boldsymbol{p}_{K^-}, \boldsymbol{p}_p)}{N_{K^-}(\boldsymbol{p}_{K^-})N_p(\boldsymbol{p}_p)} \quad \text{(= 1 in the absence of FSI/QS)}$$

- Theory (Koonin-Pratt formula)

S.E. Koonin PLB 70, 43 (1977); S. Pratt, PRD 33, 1314 (1986) $C(q) \simeq \int d^3 r S(r) |\Psi_q^{(-)}(r)|^2$

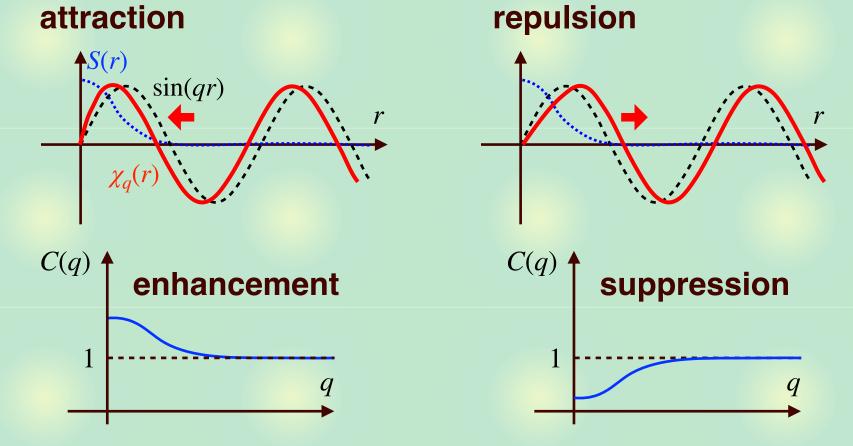
Source function S(r) < -> wave function $\Psi_q^{(-)}(r)$ (interaction)

Introduction — <u>Femtoscopy primer</u>

Wave functions and correlations

Spherical source with s-wave interaction dominance

$$C(q) \simeq 1 + \int_0^\infty dr \, S(r) \{ |\chi_q(r)|^2 - \sin^2(qr) \}$$



Qualitative behavior reflects nature of interaction

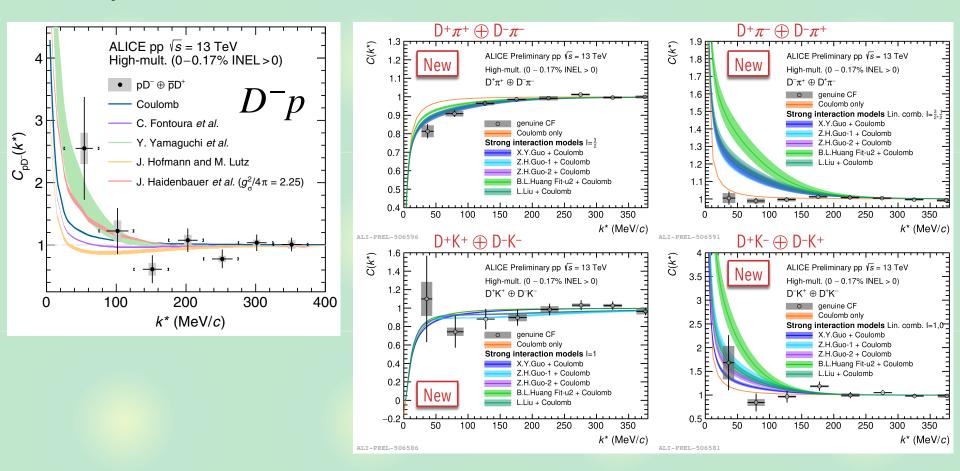
Introduction — Femtoscopy primer

Experimental data in charm sector

Observed correlation functions with charm: DN, $D\pi$, DK

ALICE collaboration, PRD 106, 052010 (2022);

Talk by F. Grosa @ Quark Matter 2022



Unique way to obtain data in charm sector (yet low statistics),



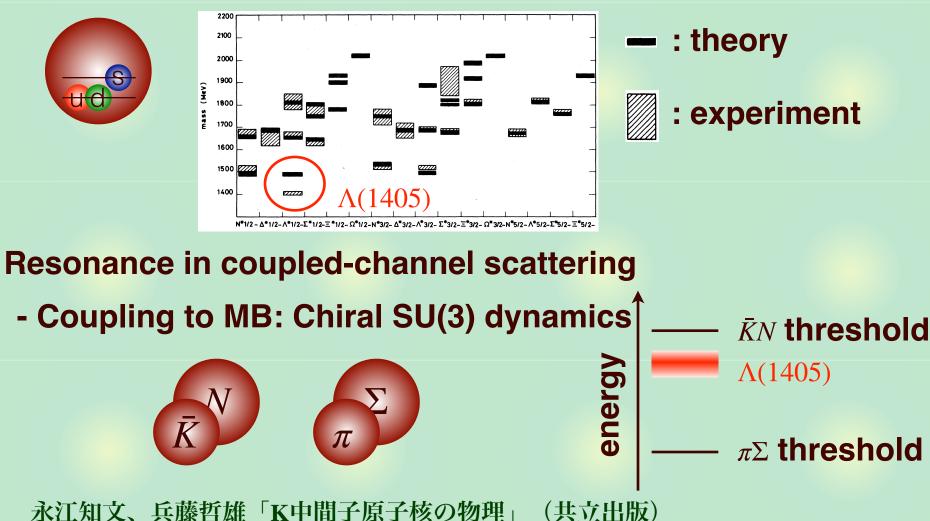
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Introduction — Femtoscopy primer Femtoscopy for exotic hadrons - K^-p correlations for $\Lambda(1405)$ Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise, PRL124, 132501 (2020) - DD^* and $D\bar{D}^*$ correlations for T_{cc} and X(3872)Y. Kamiya, T. Hyodo, A. Ohnishi, EPJA58, 131 (2022) **Femtoscopy for hypernuclei** - $\Lambda \alpha$ correlations and Λ in medium A. Jinno, Y. Kamiya, T. Hyodo, A. Ohnishi, in preparation **Summary**

$\Lambda(1405)$ and $\bar{K}N$ scattering

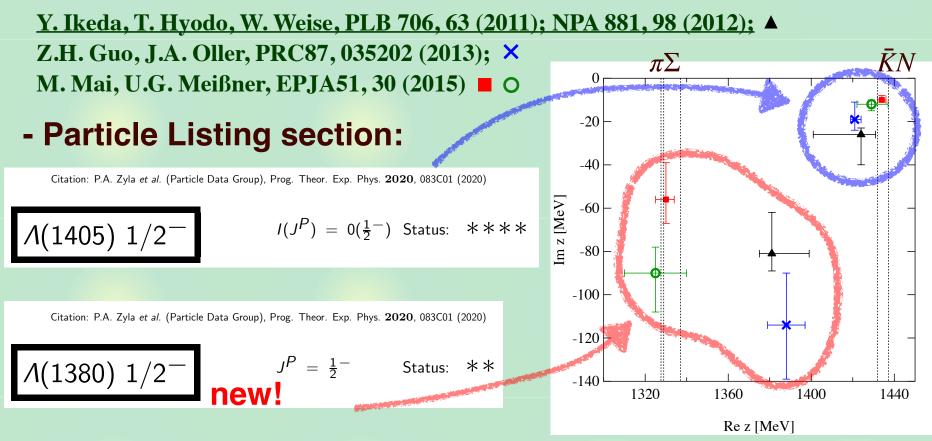
$\Lambda(1405)$ does not fit in standard picture —> exotic candidate

N. Isgur and G. Karl, PRD18, 4187 (1978)



Pole positions are determined

2020 update of PDG



T. Hyodo, M. Niiyama, Prog. Part. Nucl. Phys. 120, 103868 (2021)

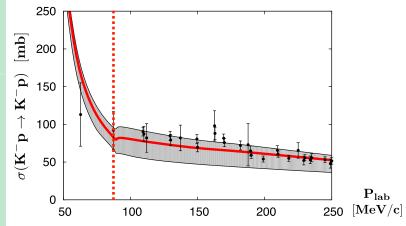
- "Λ(1405)" is no longer at 1405 MeV but ~ 1420 MeV.
- Lower pole : two-star resonance $\Lambda(1380)$

Experimental data of *K*⁻*p* **correlation**

K⁻*p* total cross sections

<u>Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011)</u>

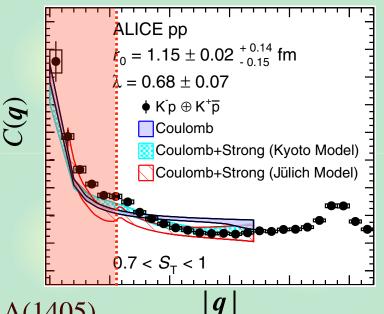
- Old bubble chamber data
- Resolution is not good
- Threshold cusp is not visible



K⁻p correlation function

ALICE collaboration, PRL 124, 092301 (2020)

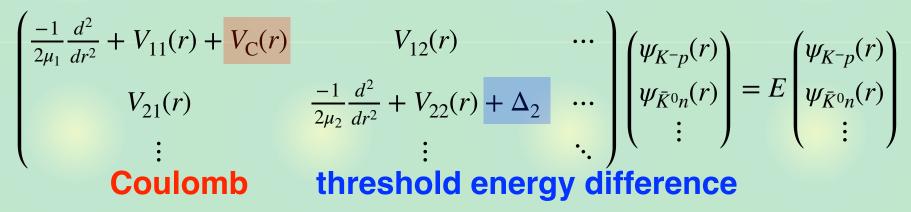
- Excellent precision ($\bar{K}^0 n$ cusp)
- Low-energy data below $\bar{K}^0 n$



-> Important constraint on $\bar{K}N$ and $\Lambda(1405)$

Coupled-channel effects

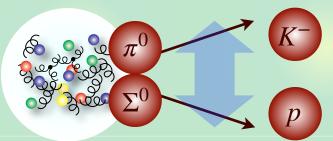
Schrödinger equation (s-wave)



Asymptotic $(r \rightarrow \infty)$ wave function

 $\begin{pmatrix} \psi_{K^-p}(r) \\ \psi_{\bar{K}^0n}(r) \\ \cdot \end{pmatrix} \propto \begin{pmatrix} \#e^{-iqr} + \#e^{iqr} \\ \#e^{-iq_2r} + \#e^{iq_2r} \\ \vdots \end{pmatrix}$ incoming + outgoing

- Transition from $\bar{K}^0 n, \pi^+ \Sigma^-, \pi^0 \Sigma^0, \pi^- \Sigma^+, \pi^0 \Lambda$ is in $\psi_i(r)$ with $i \neq K^- p$



Coupled-channel correlation function

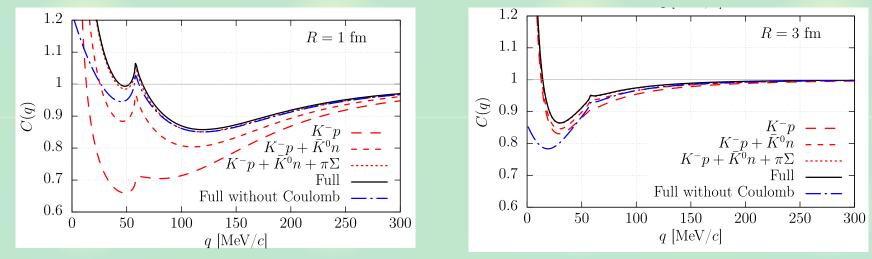
Coupled-channel Koonin-Pratt formula

- R. Lednicky, V.V. Lyuboshitz, V.L. Lyuboshitz, Phys. Atom. Nucl. 61, 2950 (1998); J. Haidenbauer, NPA 981, 1 (2019);
- Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise, PRL124, 132501 (2020)

$$C_{K^{-p}}(\boldsymbol{q}) \simeq \int d^3 \boldsymbol{r} \, S_{K^{-p}}(\boldsymbol{r}) \, |\Psi_{K^{-p},\boldsymbol{q}}^{(-)}(\boldsymbol{r})|^2 + \sum_{i \neq K^{-p}} \omega_i \int d^3 \boldsymbol{r} \, S_i(\boldsymbol{r}) \, |\Psi_{i,\boldsymbol{q}}^{(-)}(\boldsymbol{r})|^2$$
Francition from $\bar{K}^0 \boldsymbol{v} \, \pi^+ \Sigma^- \, \pi^0 \Sigma^0 \, \pi^- \Sigma^+ \, \pi^0 \Lambda$

- ω_i : weight of source channel *i* relative to K^-p

11.



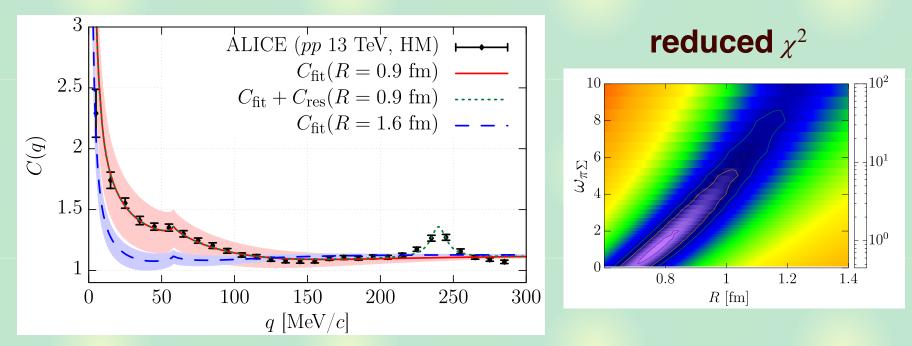
Coupled-channel effect is enhanced for small sources

Correlation from chiral SU(3) dynamics

Wave function $\Psi_{i,q}^{(-)}(r)$: Kyoto $\bar{K}N$ - $\pi\Sigma$ - $\pi\Lambda$ potential

K. Miyahara, T. Hyodo, W. Weise, PRC98, 025201 (2018)

- Source function S(r): Gaussian, $R \sim 1$ fm from K^+p data
- Source weight $\omega_{\pi\Sigma} \sim 2$ by simple statistical model estimate



Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise, PRL124, 132501 (2020)

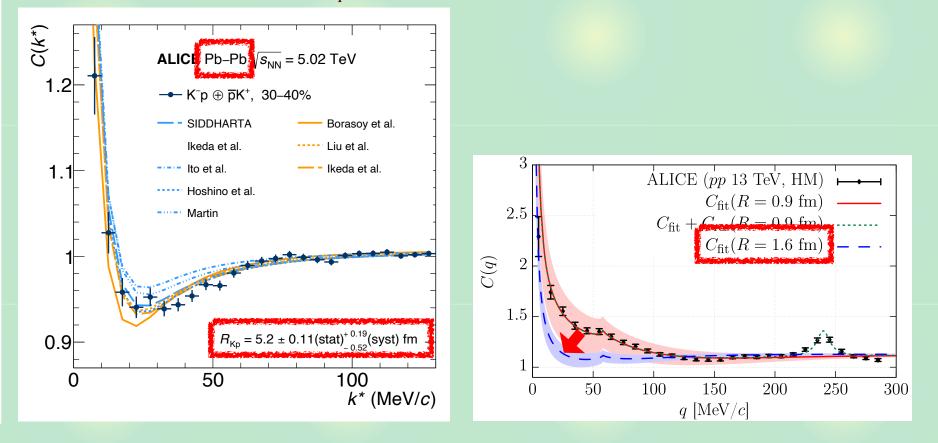
Correlation function by ALICE is well reproduced

Source size dependence

New data with Pb-Pb collisions at 5.02 TeV

ALICE collaboration, PLB 822, 136708 (2021)

- Scattering length $a_{K^-p} = -0.91 + 0.92i$ fm



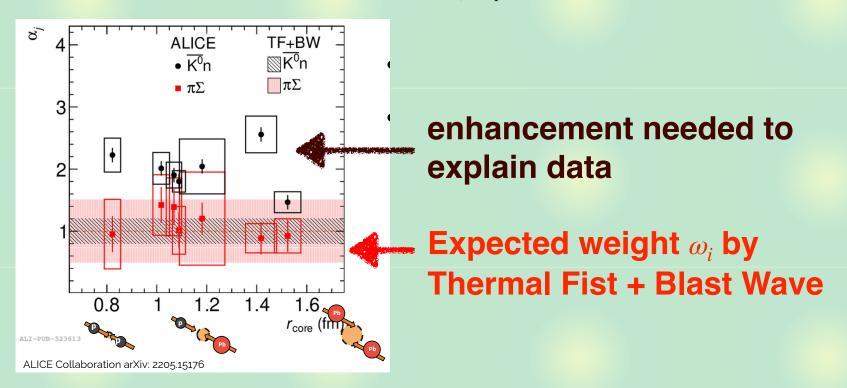
Correlation is suppressed at larger *R***, as predicted**

Systematic study of source size dependence

Correlations in *pp*, *p*-Pb, Pb-Pb **by Kyoto** $\bar{K}N$ - $\pi\Sigma$ - $\pi\Lambda$ **potential**

ALICE collaboration, EPJC 83, 340 (2023)

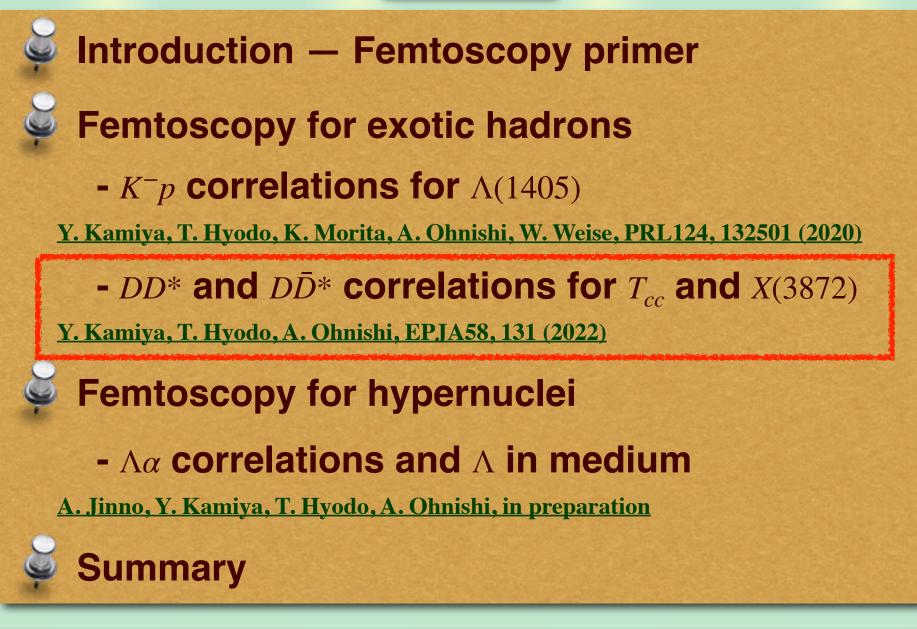
$$C_{K^{-p}}(\boldsymbol{q}) \simeq \int d^3 \boldsymbol{r} \, S_{K^{-p}}(\boldsymbol{r}) \, |\Psi_{K^{-p},\boldsymbol{q}}^{(-)}(\boldsymbol{r})|^2 + \sum_{i \neq K^{-p}} \omega_i \int d^3 \boldsymbol{r} \, S_i(\boldsymbol{r}) \, |\Psi_{i,\boldsymbol{q}}^{(-)}(\boldsymbol{r})|^2$$



More strength is needed in the $\bar{K}^0 n$ channel



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Observation of *T_{cc}*

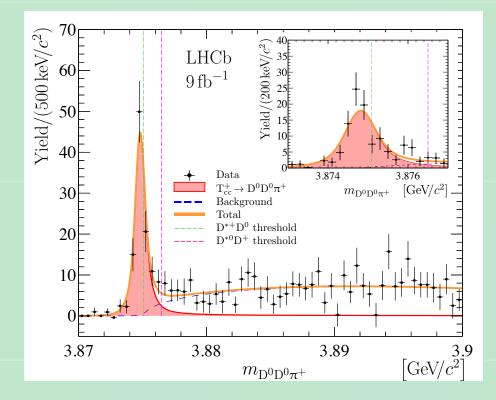
 T_{cc} observed in $D^0D^0\pi^+$ spectrum

LHCb collaboration, Nature Phys., 18, 751 (2022); Nature Comm., 13, 3351 (2022)

- Signal near DD* threshold
- Charm $C = +2 : \sim cc\bar{u}\bar{d}$
- Level structure

3870

 $3875 \begin{bmatrix} Energy (MeV) \\ ---- D^+ D^{*0} (3876.51) \\ ---- D^0 D^{*+} (3875.10) \\ \hline T_{cc} \end{bmatrix}$

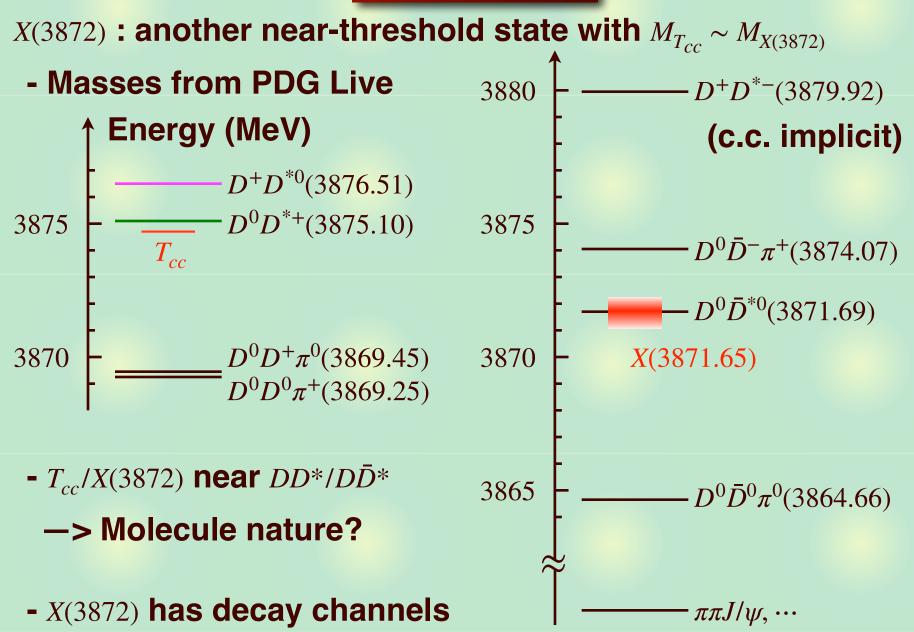


Very small (few MeV ~ keV) energy scale involved

 $D^0 D^+ \pi^0 (3869.45)$

 $D^0 D^0 \pi^+ (3869.25)$

 T_{cc} and X(3872)



18

*DD**,*DD** **potentials**

Coupled-channel potentials

$$V_{DD^*/D\bar{D}^*} = \frac{1}{2} \begin{pmatrix} V_{I=1} + V_{I=0} & V_{I=1} - V_{I=0} \\ V_{I=1} - V_{I=0} & V_{I=1} + V_{I=0} + V_c \end{pmatrix} \frac{D^0 D^{*+} / \{D^0 \bar{D}^{*0}\}}{D^+ D^{*0} / \{D^+ D^{*-}\}}$$

 \uparrow **Coulomb for** $\{D^+D^{*-}\}$

- I = 0 : one-range gaussian potentials, I = 1 neglected $V_{I=0} = V_0 \exp\{-m_{\pi}^2 r^2\}, \quad V_{I=1} = 0$ \uparrow range by π exchange

 $V_0 \in \mathbb{C}$ <-- scattering lengths (molecule picture)

- T_{cc} : $a_0^{D^0 D^{*+}} = -7.16 + i1.85$ fm (LHCb analysis)

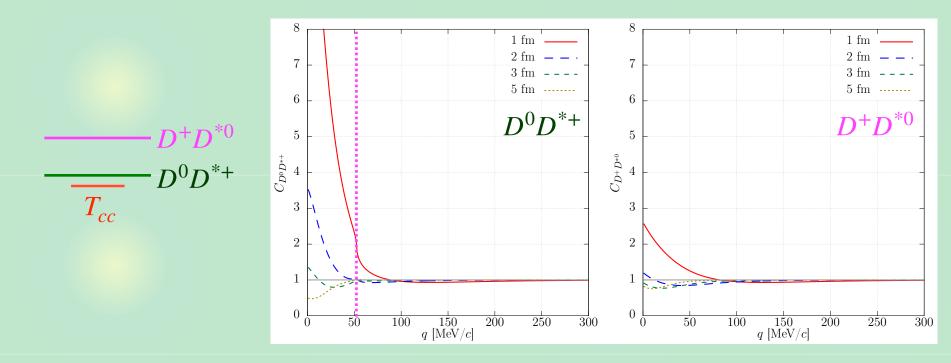
LHCb collaboration, Nature Comm., 13, 3351 (2022)

- X(3872) : $a_0^{D^0 \bar{D}^{*0}} = -4.23 + i3.95 \text{ fm} (a_0 = -i/\sqrt{2\mu E_h} \text{ with PDG } E_h)$ 19

 $DD^* \sim T_{cc}$ sector

D^0D^{*+} and D^+D^{*0} correlation functions ($cc\bar{u}d$, exotic)

Y. Kamiya, T. Hyodo, A. Ohnishi, EPJA58, 131 (2022)



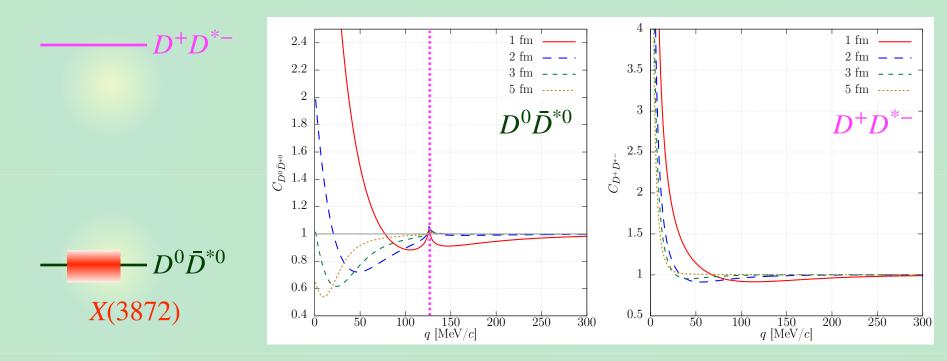
- Bound state feature (source size dep.) in both channels

- Strong signal in D^0D^{*+} , weaker one in D^+D^{*0}
- D^+D^{*0} cusp in D^0D^{*+} ($q \sim 52$ MeV) is not very prominent

$D\bar{D}^* \sim X(3872)$ sector

$D^0 \overline{D}^{*0}$ and $D^+ \overline{D}^{*-}$ correlation functions ($c \overline{c} q \overline{q}$)

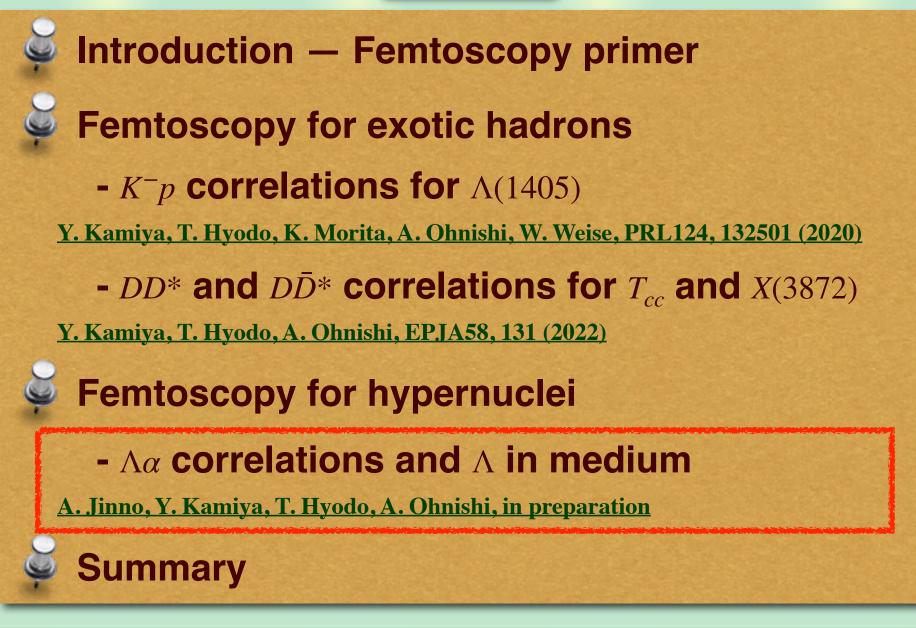
Y. Kamiya, T. Hyodo, A. Ohnishi, EPJA58, 131 (2022)



- Bound state feature in $D^0 \overline{D}^{*0}$ correlation
- Sizable D^+D^{*-} cusp in $D^0\overline{D}^{*0}$ ($q \sim 126 \text{ MeV}$)
- D+D*- correlation : Coulomb attraction dominance



Contents



Motivation

Hyperon puzzle in neutron stars

- ANN three-body force for repulsion at high density

D. Gerstung, N. Kaiser, W. Weise, EPJA 55, 175 (2020)

How to verify this in experiments?

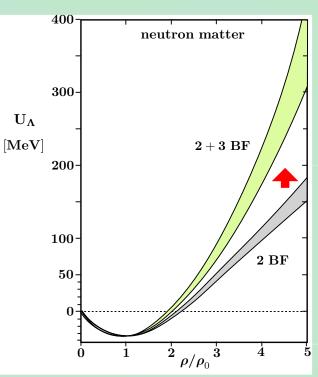
- Λ directed flow in heavy ion collisions

Y. Nara, A. Jinno, K. Murase, A. Ohnishi, PRC 106, 044902 (2022)

A-nucleus correlation function?

- Heavy nuclei are difficult to produce
- Strong binding of α —> high central density $\gtrsim 2\rho_0$

Repulsion at high density by $\Lambda \alpha$ correlation function?



$\Lambda \alpha$ potentials

Skyrme-Hartree Fock methods for Λ hypernuclei

- LY4 : empirical potential

D.E. Lanskoy, Y. Yamamoto, PRC 55, 2330 (1997)

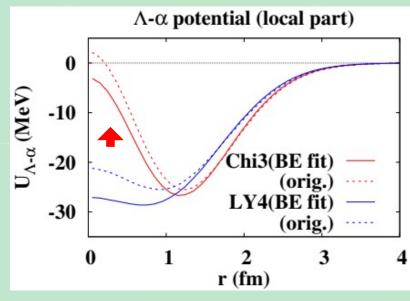
- Chi3 : based on chiral EFT with ΛNN force

A. Jinno, K. Murase, Y. Nara, A. Ohnishi, in preparation

- Both potentials reproduce hypernuclear data from C to Pb

$\Lambda \alpha$ potentials

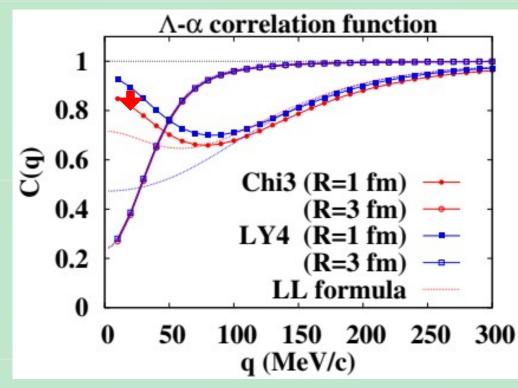
- overestimate ⁵_AHe binding energy
 adjustment of parameters
- LY4 : Woods-Saxon like
- Chi3 : central repulsion



$\Lambda \alpha$ correlation functions

Results of correlation function

A. Jinno, Y. Kamiya, T. Hyodo, A. Ohnishi, in preparation



- Bound state signature (dip at small q)
- Central repulsion: slightly stronger correlation for R = 1 fm
- Int. range ~ a few fm —> LL does not work for R = 1 fm

Summary

