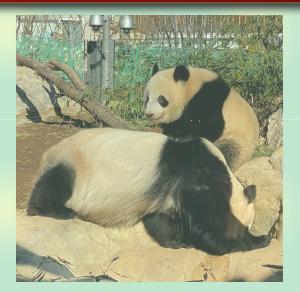
# Femtoscopy for exotic hadrons and nuclei





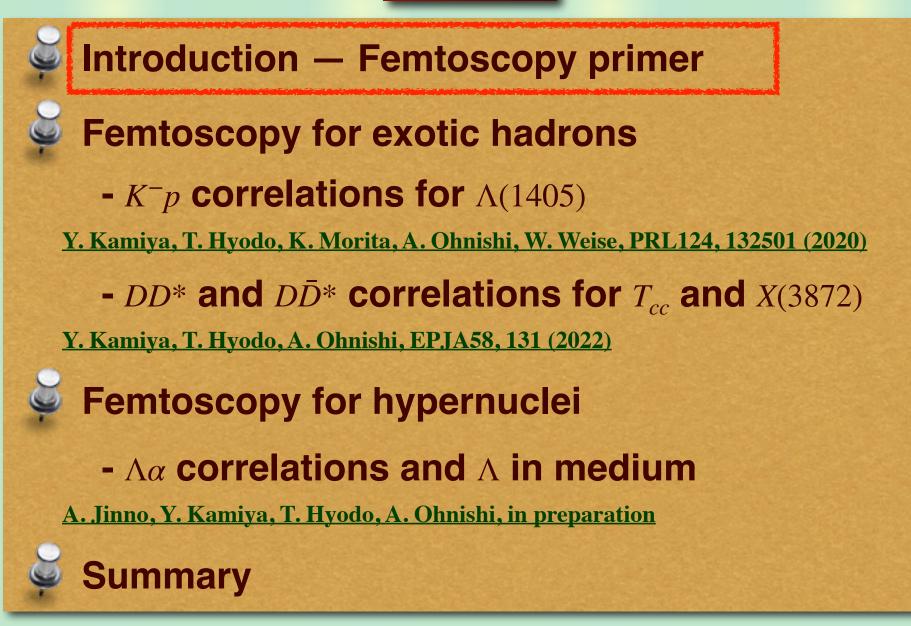
## **Tetsuo Hyodo**

Tokyo Metropolitan Univ.



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#### Introduction — Femtoscopy primer

## **Scattering experiments and femtoscopy**

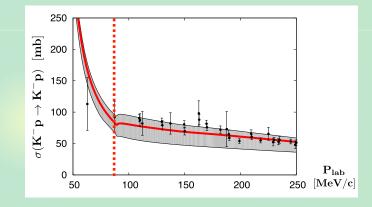
## Traditional methods: scattering experiments

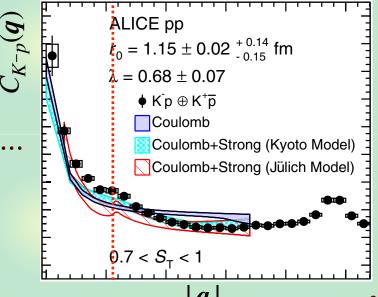
Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011)

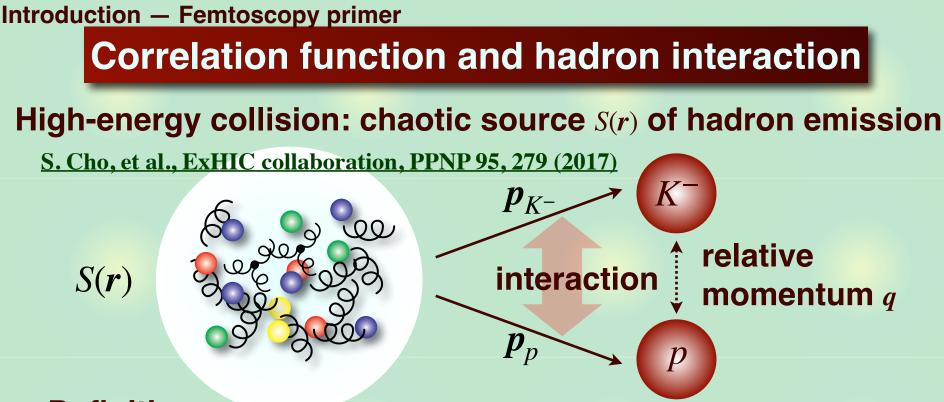
- limited channels: NN, YN,  $\pi N$ , KN,  $\bar{K}N$ , ...
- limited statistics (low-energy)
- heavy (c, b) hadrons: impossible

### **Femtoscopy: correlation function** ALICE collaboration, PRL 124, 092301 (2020)

- various systems:  $\Lambda\Lambda, N\Omega, \phi N, \overline{K}\Lambda, DN, \cdots$
- Excellent precision ( $\bar{K}^0n$  cusp)
- heavy hadrons: possible!







### - Definition

$$C(\boldsymbol{q}) = \frac{N_{K^-p}(\boldsymbol{p}_{K^-}, \boldsymbol{p}_p)}{N_{K^-}(\boldsymbol{p}_{K^-})N_p(\boldsymbol{p}_p)} \quad \text{(= 1 in the absence of FSI/QS)}$$

### - Theory (Koonin-Pratt formula)

S.E. Koonin PLB 70, 43 (1977); S. Pratt, PRD 33, 1314 (1986)  $C(q) \simeq \int d^3 r S(r) |\Psi_q^{(-)}(r)|^2$ 

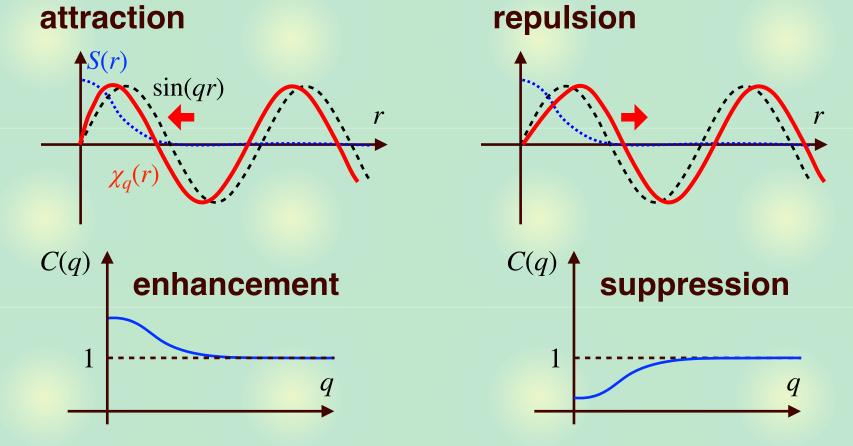
### Source function S(r) < -> wave function $\Psi_q^{(-)}(r)$ (interaction)

#### Introduction — <u>Femtoscopy primer</u>

## **Wave functions and correlations**

### Spherical source with s-wave interaction dominance

$$C(q) \simeq 1 + \int_0^\infty dr \, S(r) \{ |\chi_q(r)|^2 - \sin^2(qr) \}$$



### **Qualitative behavior reflects nature of interaction**

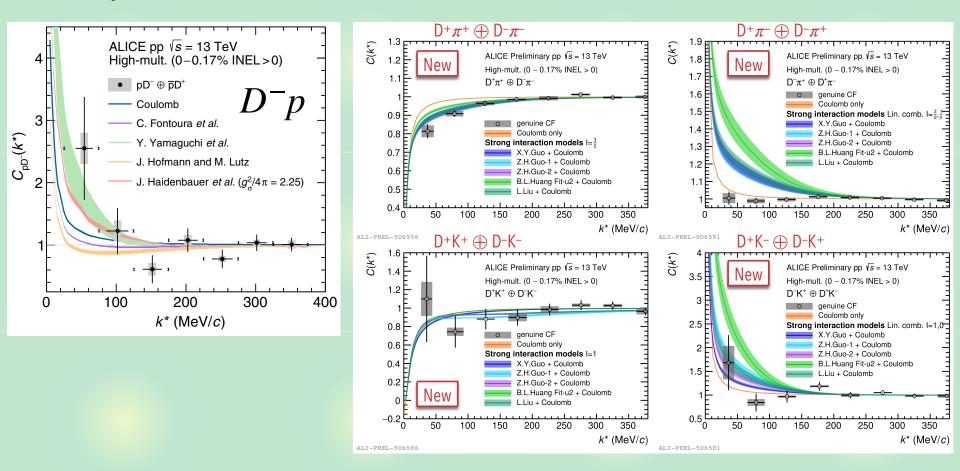
#### Introduction — Femtoscopy primer

## **Experimental data in charm sector**

### **Observed correlation functions with charm:** DN, $D\pi$ , DK

#### ALICE collaboration, PRD 106, 052010 (2022);

Talk by F. Grosa @ Quark Matter 2022



### Unique way to obtain data in charm sector (yet low statistics),



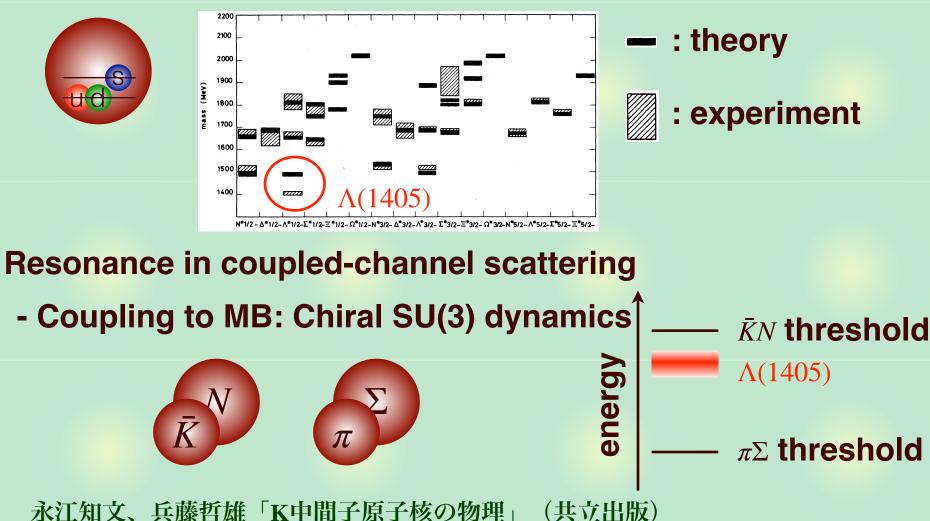
## Contents

**Introduction — Femtoscopy primer** Femtoscopy for exotic hadrons -  $K^-p$  correlations for  $\Lambda(1405)$ Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise, PRL124, 132501 (2020) -  $DD^*$  and  $D\bar{D}^*$  correlations for  $T_{cc}$  and X(3872)Y. Kamiya, T. Hyodo, A. Ohnishi, EPJA58, 131 (2022) **Femtoscopy for hypernuclei** -  $\Lambda \alpha$  correlations and  $\Lambda$  in medium A. Jinno, Y. Kamiya, T. Hyodo, A. Ohnishi, in preparation **Summary** 

## $\Lambda(1405)$ and $\bar{K}N$ scattering

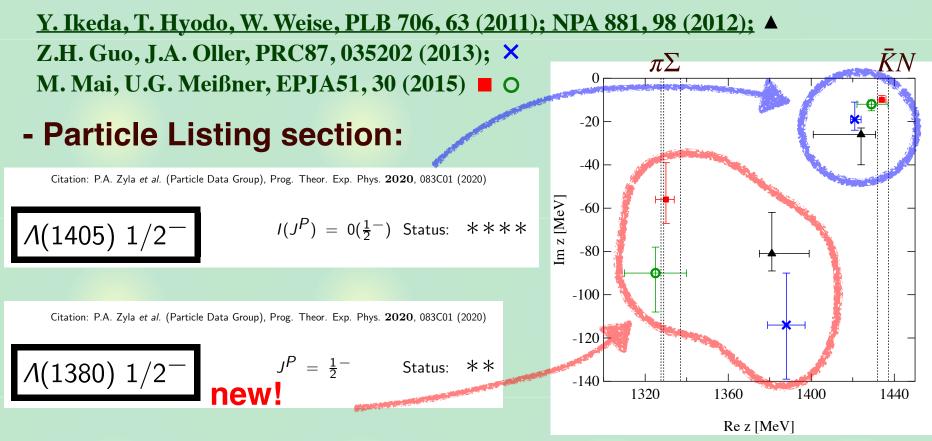
## $\Lambda(1405)$ does not fit in standard picture —> exotic candidate

N. Isgur and G. Karl, PRD18, 4187 (1978)



## **Pole positions are determined**

## 2020 update of PDG



**T. Hyodo, M. Niiyama, Prog. Part. Nucl. Phys. 120, 103868 (2021)** 

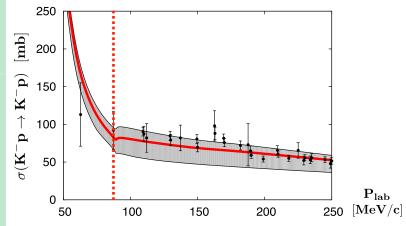
- "Λ(1405)" is no longer at 1405 MeV but ~ 1420 MeV.
- Lower pole : two-star resonance  $\Lambda(1380)$

## **Experimental data of** *K*<sup>-</sup>*p* **correlation**

*K*<sup>-</sup>*p* total cross sections

<u>Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011)</u>

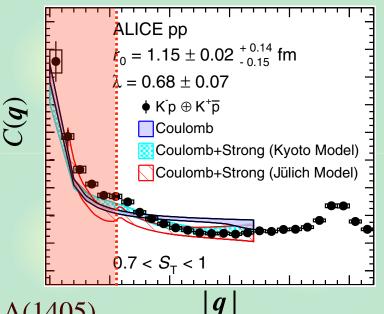
- Old bubble chamber data
- Resolution is not good
- Threshold cusp is not visible



## *K<sup>-</sup>p* correlation function

ALICE collaboration, PRL 124, 092301 (2020)

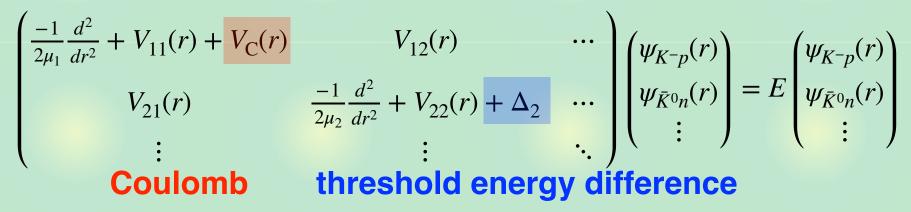
- Excellent precision ( $\bar{K}^0 n$  cusp)
- Low-energy data below  $\bar{K}^0 n$



-> Important constraint on  $\bar{K}N$  and  $\Lambda(1405)$ 

## **Coupled-channel effects**

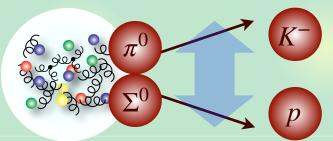
Schrödinger equation (s-wave)



## Asymptotic $(r \rightarrow \infty)$ wave function

 $\begin{pmatrix} \psi_{K^-p}(r) \\ \psi_{\bar{K}^0n}(r) \\ \cdot \end{pmatrix} \propto \begin{pmatrix} \#e^{-iqr} + \#e^{iqr} \\ \#e^{-iq_2r} + \#e^{iq_2r} \\ \vdots \end{pmatrix}$  incoming + outgoing

- Transition from  $\bar{K}^0 n, \pi^+ \Sigma^-, \pi^0 \Sigma^0, \pi^- \Sigma^+, \pi^0 \Lambda$  is in  $\psi_i(r)$  with  $i \neq K^- p$ 



## **Coupled-channel correlation function**

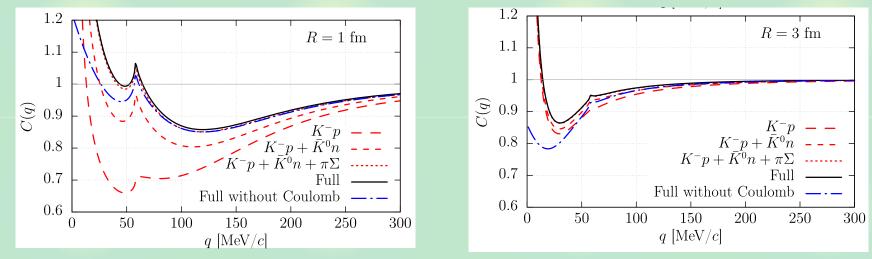
## **Coupled-channel Koonin-Pratt formula**

- R. Lednicky, V.V. Lyuboshitz, V.L. Lyuboshitz, Phys. Atom. Nucl. 61, 2950 (1998); J. Haidenbauer, NPA 981, 1 (2019);
- Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise, PRL124, 132501 (2020)

$$C_{K^{-p}}(\boldsymbol{q}) \simeq \int d^3 \boldsymbol{r} \, S_{K^{-p}}(\boldsymbol{r}) \, |\Psi_{K^{-p},\boldsymbol{q}}^{(-)}(\boldsymbol{r})|^2 + \sum_{i \neq K^{-p}} \omega_i \int d^3 \boldsymbol{r} \, S_i(\boldsymbol{r}) \, |\Psi_{i,\boldsymbol{q}}^{(-)}(\boldsymbol{r})|^2$$
Francition from  $\bar{K}^0 \boldsymbol{v} \, \pi^+ \Sigma^- \, \pi^0 \Sigma^0 \, \pi^- \Sigma^+ \, \pi^0 \Lambda$ 

-  $\omega_i$  : weight of source channel *i* relative to  $K^-p$ 

11.



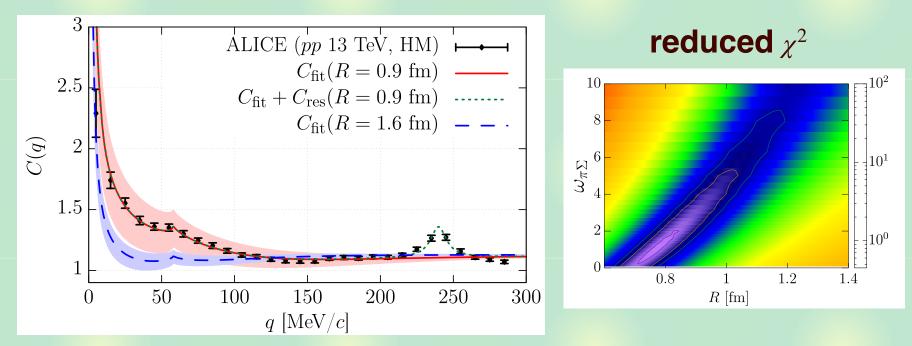
### **Coupled-channel effect is enhanced for small sources**

## **Correlation from chiral SU(3) dynamics**

Wave function  $\Psi_{i,q}^{(-)}(r)$ : Kyoto  $\bar{K}N$ - $\pi\Sigma$ - $\pi\Lambda$  potential

K. Miyahara, T. Hyodo, W. Weise, PRC98, 025201 (2018)

- Source function S(r): Gaussian,  $R \sim 1$  fm from  $K^+p$  data
- Source weight  $\omega_{\pi\Sigma} \sim 2$  by simple statistical model estimate



Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise, PRL124, 132501 (2020)

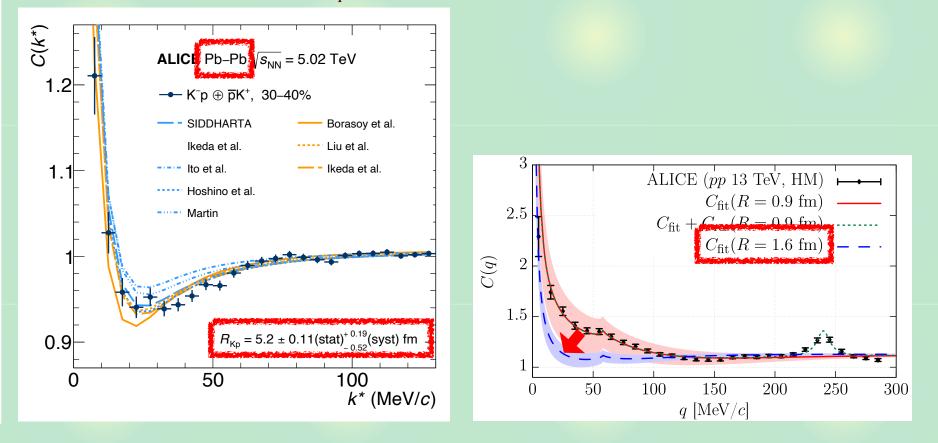
## **Correlation function by ALICE is well reproduced**

## Source size dependence

## New data with Pb-Pb collisions at 5.02 TeV

ALICE collaboration, PLB 822, 136708 (2021)

### - Scattering length $a_{K^-p} = -0.91 + 0.92i$ fm



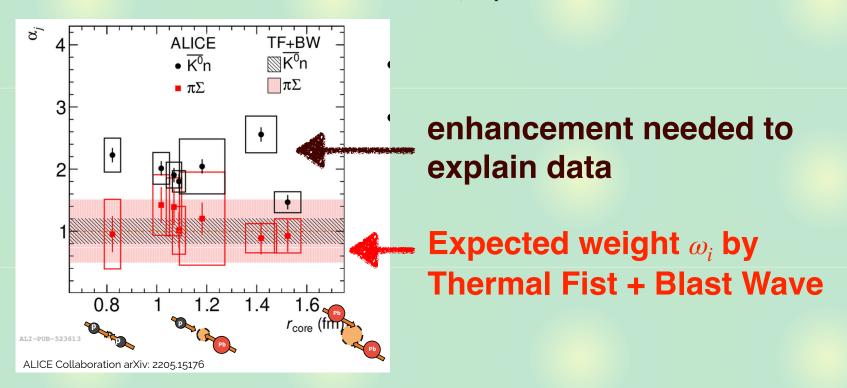
### **Correlation is suppressed at larger** *R***, as predicted**

## Systematic study of source size dependence

## **Correlations in** *pp*, *p*-Pb, Pb-Pb **by Kyoto** $\bar{K}N$ - $\pi\Sigma$ - $\pi\Lambda$ **potential**

ALICE collaboration, EPJC 83, 340 (2023)

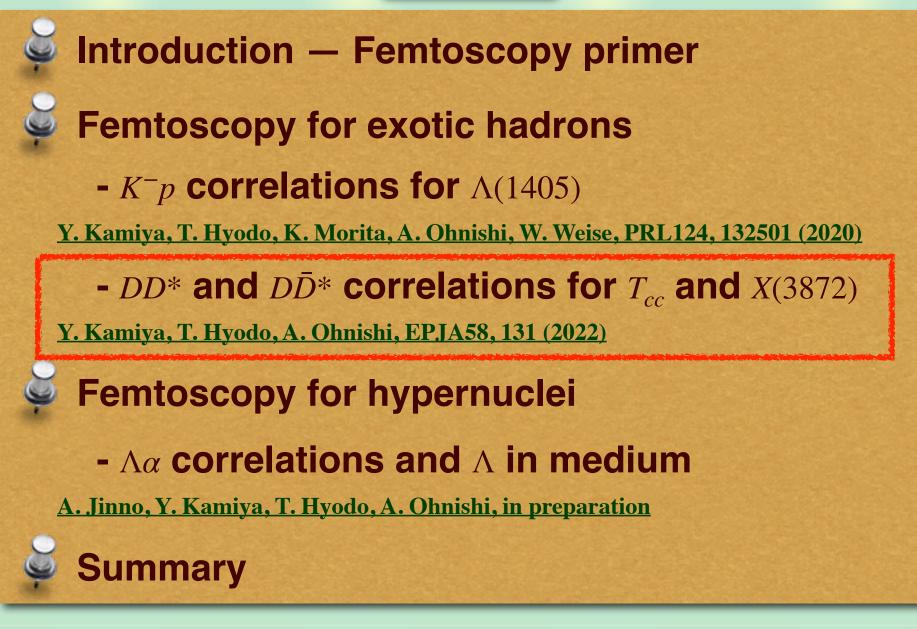
$$C_{K^{-p}}(\boldsymbol{q}) \simeq \int d^3 \boldsymbol{r} \, S_{K^{-p}}(\boldsymbol{r}) \, |\Psi_{K^{-p},\boldsymbol{q}}^{(-)}(\boldsymbol{r})|^2 + \sum_{i \neq K^{-p}} \omega_i \int d^3 \boldsymbol{r} \, S_i(\boldsymbol{r}) \, |\Psi_{i,\boldsymbol{q}}^{(-)}(\boldsymbol{r})|^2$$



### More strength is needed in the $\bar{K}^0 n$ channel



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## **Observation of** *T<sub>cc</sub>*

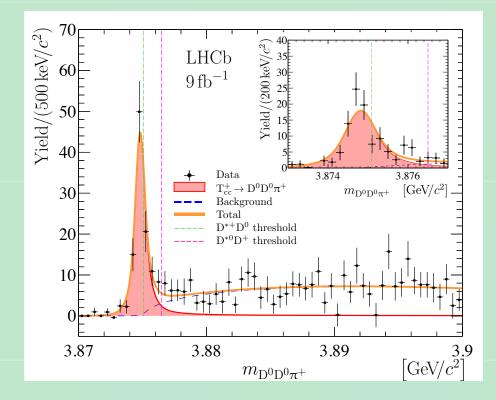
 $T_{cc}$  observed in  $D^0D^0\pi^+$  spectrum

LHCb collaboration, Nature Phys., 18, 751 (2022); Nature Comm., 13, 3351 (2022)

- Signal near DD\* threshold
- Charm  $C = +2 : \sim cc\bar{u}\bar{d}$
- Level structure

3870

 $3875 \begin{bmatrix} Energy (MeV) \\ ---- D^+ D^{*0} (3876.51) \\ ---- D^0 D^{*+} (3875.10) \\ \hline T_{cc} \end{bmatrix}$ 

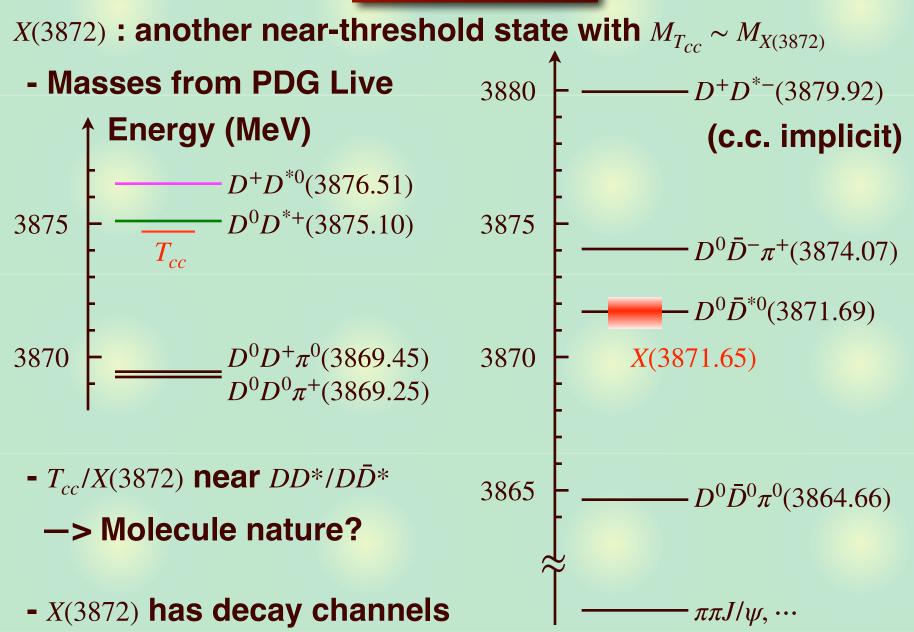


Very small (few MeV ~ keV) energy scale involved

 $D^0 D^+ \pi^0 (3869.45)$ 

 $D^0 D^0 \pi^+ (3869.25)$ 

 $T_{cc}$  and X(3872)



18

*DD*\*,*DD*\* **potentials** 

**Coupled-channel potentials** 

$$V_{DD^*/D\bar{D}^*} = \frac{1}{2} \begin{pmatrix} V_{I=1} + V_{I=0} & V_{I=1} - V_{I=0} \\ V_{I=1} - V_{I=0} & V_{I=1} + V_{I=0} + V_c \end{pmatrix} \frac{D^0 D^{*+} / \{D^0 \bar{D}^{*0}\}}{D^+ D^{*0} / \{D^+ D^{*-}\}}$$

 $\uparrow$  **Coulomb for**  $\{D^+D^{*-}\}$ 

- I = 0 : one-range gaussian potentials, I = 1 neglected  $V_{I=0} = V_0 \exp\{-m_{\pi}^2 r^2\}, \quad V_{I=1} = 0$  $\uparrow$  range by  $\pi$  exchange

 $V_0 \in \mathbb{C}$  <-- scattering lengths (molecule picture)

-  $T_{cc}$  :  $a_0^{D^0 D^{*+}} = -7.16 + i1.85$  fm (LHCb analysis)

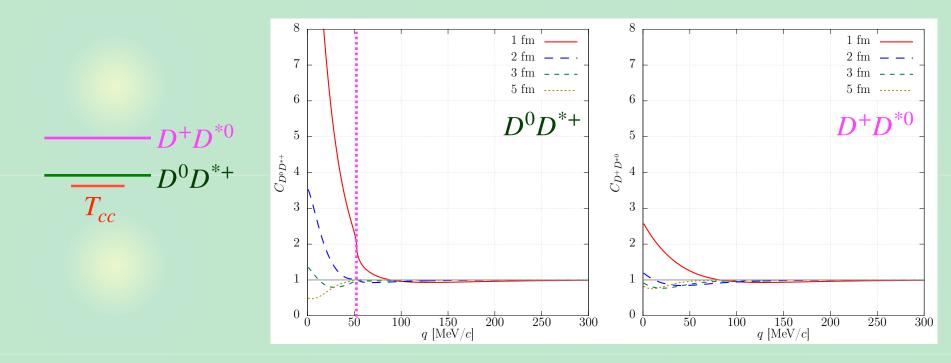
LHCb collaboration, Nature Comm., 13, 3351 (2022)

- X(3872) :  $a_0^{D^0 \bar{D}^{*0}} = -4.23 + i3.95 \text{ fm} (a_0 = -i/\sqrt{2\mu E_h} \text{ with PDG } E_h)$  19

 $DD^* \sim T_{cc}$  sector

### $D^0D^{*+}$ and $D^+D^{*0}$ correlation functions ( $cc\bar{u}d$ , exotic)

Y. Kamiya, T. Hyodo, A. Ohnishi, EPJA58, 131 (2022)



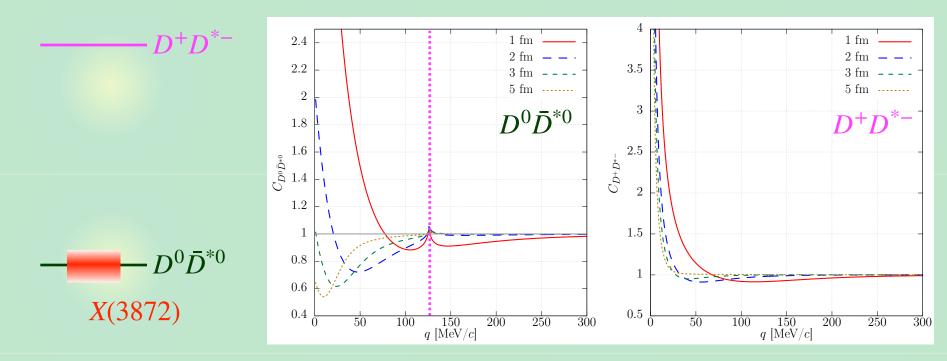
- Bound state feature (source size dep.) in both channels

- Strong signal in  $D^0D^{*+}$ , weaker one in  $D^+D^{*0}$
- $D^+D^{*0}$  cusp in  $D^0D^{*+}$  ( $q \sim 52$  MeV) is not very prominent

## $D\bar{D}^* \sim X(3872)$ sector

### $D^0 \overline{D}^{*0}$ and $D^+ \overline{D}^{*-}$ correlation functions ( $c \overline{c} q \overline{q}$ )

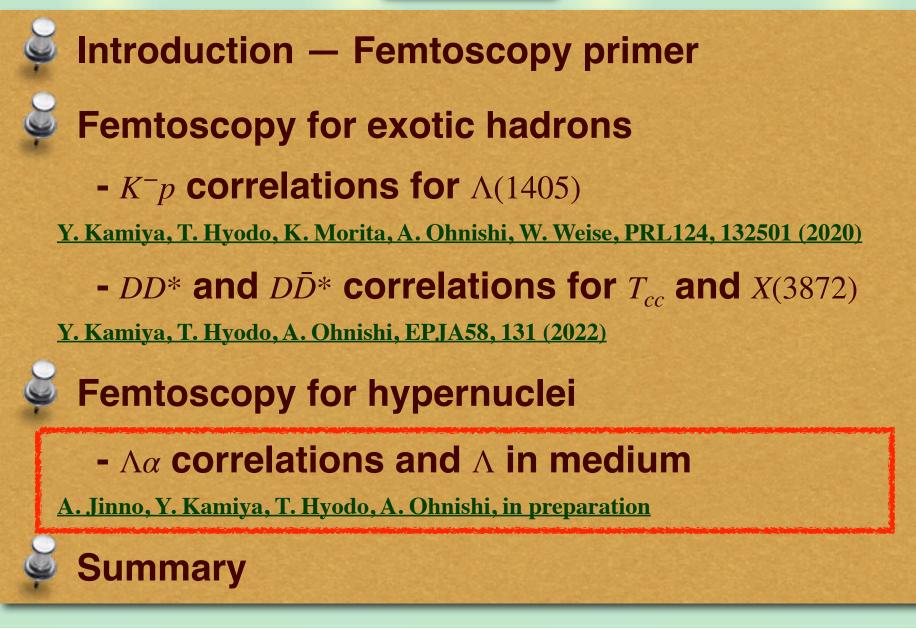
Y. Kamiya, T. Hyodo, A. Ohnishi, EPJA58, 131 (2022)



- Bound state feature in  $D^0 \overline{D}^{*0}$  correlation
- Sizable  $D^+D^{*-}$  cusp in  $D^0\overline{D}^{*0}$  ( $q \sim 126 \text{ MeV}$ )
- D+D\*- correlation : Coulomb attraction dominance



## Contents



Motivation

## Hyperon puzzle in neutron stars

- ANN three-body force for repulsion at high density

D. Gerstung, N. Kaiser, W. Weise, EPJA 55, 175 (2020)

## How to verify this in experiments?

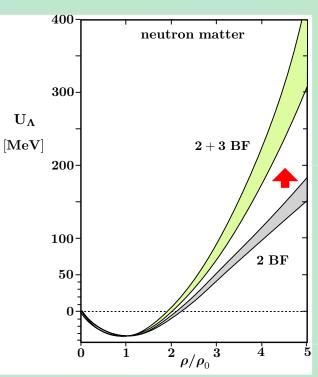
-  $\Lambda$  directed flow in heavy ion collisions

Y. Nara, A. Jinno, K. Murase, A. Ohnishi, PRC 106, 044902 (2022)

A-nucleus correlation function?

- Heavy nuclei are difficult to produce
- Strong binding of  $\alpha$  —> high central density  $\gtrsim 2\rho_0$

## Repulsion at high density by $\Lambda \alpha$ correlation function?



## $\Lambda \alpha$ potentials

## Skyrme-Hartree Fock methods for $\Lambda$ hypernuclei

## - LY4 : empirical potential

D.E. Lanskoy, Y. Yamamoto, PRC 55, 2330 (1997)

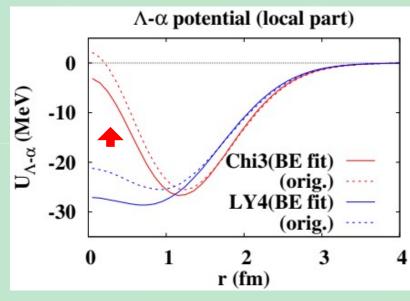
## - Chi3 : based on chiral EFT with $\Lambda NN$ force

A. Jinno, K. Murase, Y. Nara, A. Ohnishi, in preparation

## - Both potentials reproduce hypernuclear data from C to Pb

## $\Lambda \alpha$ potentials

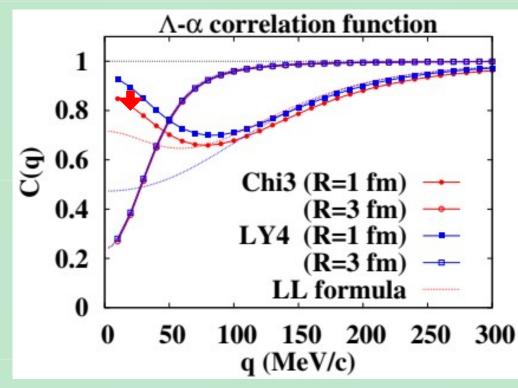
- overestimate <sup>5</sup><sub>A</sub>He binding energy
   adjustment of parameters
- LY4 : Woods-Saxon like
- Chi3 : central repulsion



## $\Lambda \alpha$ correlation functions

## **Results of correlation function**

A. Jinno, Y. Kamiya, T. Hyodo, A. Ohnishi, in preparation



- Bound state signature (dip at small q)
- Central repulsion: slightly stronger correlation for R = 1 fm
- Int. range ~ a few fm —> LL does not work for R = 1 fm

## Summary

