

$\Lambda(1405)$ as a hadronic molecule



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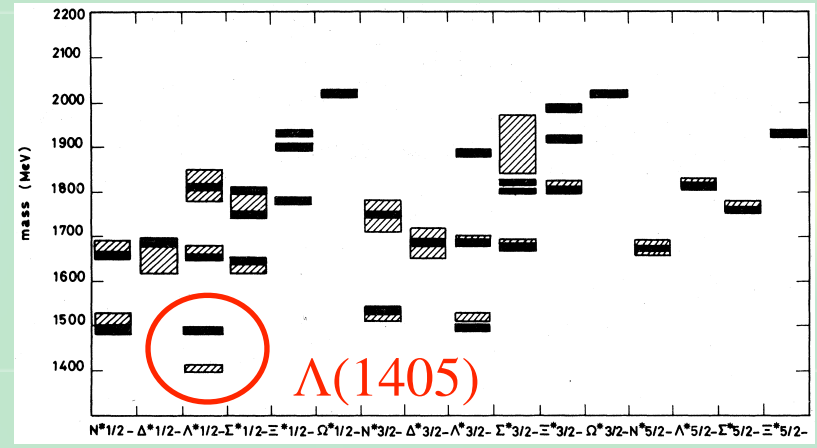
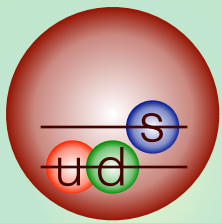


2022, Feb. 18th 1

$\Lambda(1405)$ in meson-baryon scattering

$\Lambda(1405)$ does not fit in standard picture \rightarrow exotic structure?

N. Isgur, G. Karl, PRD18, 4187 (1978)

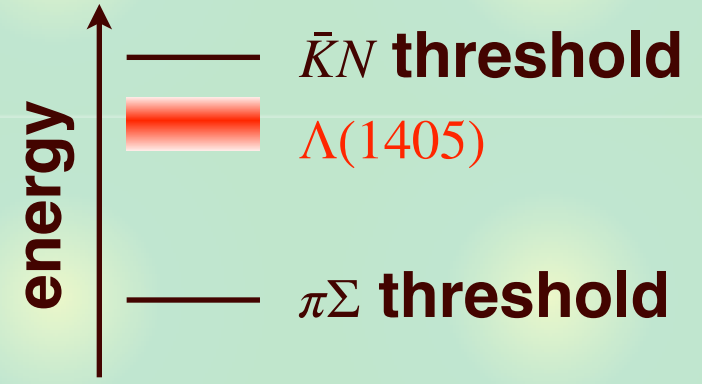
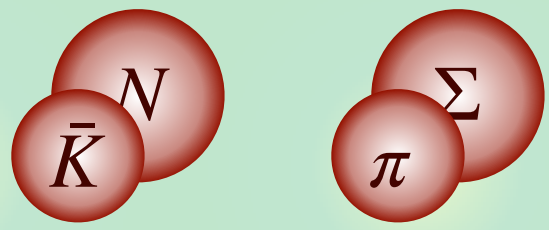


— : theory
 ▨ : experiment

Resonance in coupled-channel scattering

R.H. Dalitz, T.C. Wong, G. Rajasekaran, PR153, 1617 (1967)

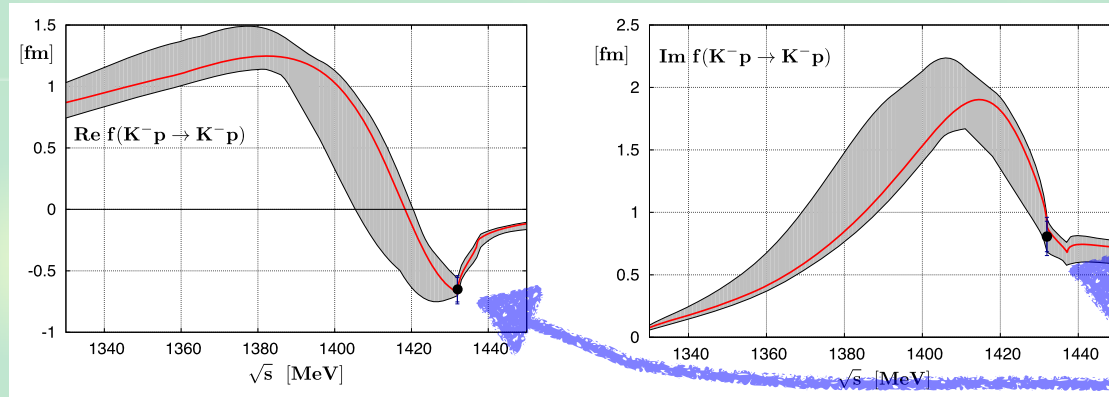
- coupling to MB states



Detailed analysis of $\bar{K}N-\pi\Sigma$ scattering is necessary.

Pole positions determined

Recent analyses with chiral SU(3) dynamics at NLO



**Kaonic hydrogen
by SIDDHARTA**

- [14,15] Y. Ikeda, T. Hyodo, W. Weise, *PLB* **706**, 63 (2011); *NPA* **881**, 98 (2012),
 [17] Z.H. Guo, J.A. Oller, *PRC* **87**, 035202 (2013),
 [18] M. Mai, U.G. Meißner, *EPJA* **51**, 30 (2015)

approach	pole 1 [MeV]	pole 2 [MeV]
Refs. [14, 15], NLO	$1424_{-23}^{+7} - i 26_{-14}^{+3}$	$1381_{-6}^{+18} - i 81_{-8}^{+19}$
Ref. [17], Fit II	$1421_{-2}^{+3} - i 19_{-5}^{+8}$	$1388_{-9}^{+9} - i 114_{-25}^{+24}$
Ref. [18], solution #2	$1434_{-2}^{+2} - i 10_{-1}^{+2}$	$1330_{-5}^{+4} - i 56_{-11}^{+17}$
Ref. [18], solution #4	$1429_{-7}^{+8} - i 12_{-3}^{+2}$	$1325_{-15}^{+15} - i 90_{-18}^{+12}$

- Two poles: superposition of two eigenstates

J.A. Oller, U.G. Meißner, *PLB* **500**, 263 (2001);

D. Jido, J.A. Oller, E. Oset, A. Ramos, U.G. Meißner, *NPA* **723**, 205 (2003)

$\Lambda(1405)$ in PDG

2020 update of PDG

P.A. Zyla, et al., PTEP 2020, 083C01 (2020); <http://pdg.lbl.gov/>

- Particle Listing section:

Citation: P.A. Zyla et al. (Particle Data Group), Prog. Theor. Exp. Phys. **2020**, 083C01 (2020)

$\Lambda(1405) 1/2^-$

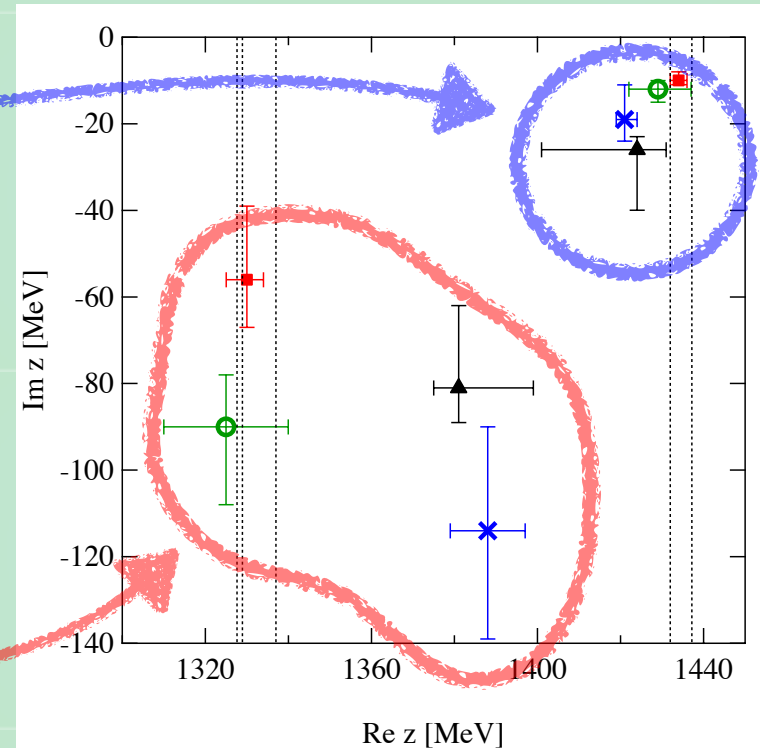
$I(J^P) = 0(\frac{1}{2}^-)$ Status: ****

Citation: P.A. Zyla et al. (Particle Data Group), Prog. Theor. Exp. Phys. **2020**, 083C01 (2020)

$\Lambda(1380) 1/2^-$

new!

$J^P = \frac{1}{2}^-$ Status: **



T. Hyodo, M. Niiyama, PPNP 120, 103868 (2021);

T. Hyodo, W. Weise, arXiv:2202.06181 [nucl-th]

- $\Lambda(1405)$ is no longer at 1405 MeV but ~ 1420 MeV.
- Lower pole: two-star resonance $\Lambda(1380)$

Next step: internal structure



Structure of $\Lambda(1405)$?



Weak binding relation for **stable** bound states

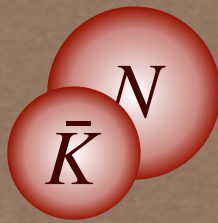
S. Weinberg, Phys. Rev. 137, B672 (1965)



Generalization to **unstable** resonances

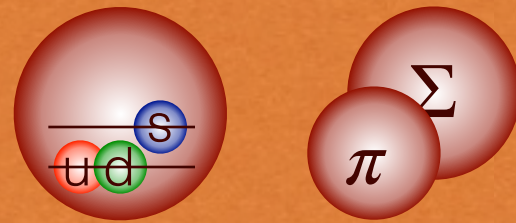
Y. Kamiya, T. Hyodo, PRC93, 035203 (2016); PTEP2017, 023D02 (2017)

Compositeness X
threshold channel



or

“Elementarity” Z
other contributions



observables (a_0, E_h)

Weak-binding relation for stable states

Compositeness X of s-wave **weakly bound state** ($R \gg R_{\text{typ}}$)

S. Weinberg, *Phys. Rev.* **137**, B672 (1965);

T. Hyodo, *Int. J. Mod. Phys. A* **28**, 1330045 (2013)

$$|d\rangle = \sqrt{X} |NN\rangle + \sqrt{1-X} |\text{others}\rangle$$

NN

continuum



deuteron

range of interaction

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\}, \quad R = \frac{1}{\sqrt{2\mu B}}$$

↑
scattering length

↑
radius of bound state

- Deuteron is NN composite: $a_0 \sim R \Rightarrow X \sim 1$
- Internal structure from **observables** (a_0, B)

Quantitative discussion?

Detour: deuteron in more detail

Weak-binding relation

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\}, \quad R = \frac{1}{\sqrt{2\mu B}}$$

- empirical values for deuteron case

$$a_0 \sim 5.42 \text{ fm}, \quad R \sim 4.32 \text{ fm}$$

- neglecting $\mathcal{O}(R_{\text{typ}}/R)$ term : contradiction with $0 \leq X \leq 1$?

Y. Kamiya, T. Hyodo, PoS INPC2016, 270 (2017),

Y. Li, F.K. Guo, J.Y. Pang, J.J. Wu, arXiv:2110.02766 [hep-ph],

J. Song, L.R. Dai, E. Oset arXiv:2201.04414 [hep-ph].

$$X \sim 1.68$$

If $0 \leq X \leq 1$, then $a_0 < R$

—> For systems with $a_0 > R$, $\mathcal{O}(R_{\text{typ}}/R)$ term is important.

T. Kinugawa, T. Hyodo, in preparation

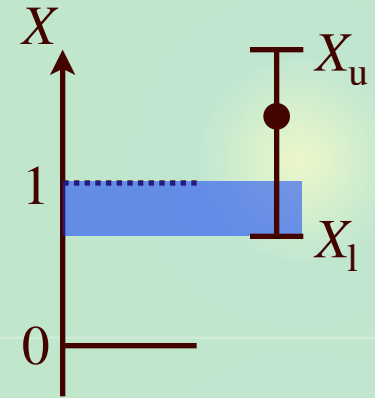
Detour: range correction

Uncertainty estimation with $\mathcal{O}(R_{\text{typ}}/R)$ term

Y. Kamiya, T. Hyodo, PTEP2017, 023D02 (2017)

$$X_u = \frac{a_0/R + \xi}{2 - a_0/R - \xi}, \quad X_l = \frac{a_0/R - \xi}{2 - a_0/R + \xi}, \quad \xi = \frac{R_{\text{typ}}}{R}$$

- exclude region outside $0 \leq X \leq 1$



Application and finite range correction

$$R_{\text{typ}} = \max\{R_{\text{int}}, R_{\text{eff}}\}$$

T. Kinugawa, T. Hyodo, arXiv:2111.06619; 2112.00249; 2201.04283 [hep-ph]

Bound state	$R_{\text{typ}} = R_{\text{eff}}$	$R_{\text{typ}} = R_{\text{int}}$	This work
d	$1.68^{+3.19}_{-0.943}$	$1.68^{+2.14}_{-0.823}$	$0.738 \leq X \leq 1$
$X(3872)$	$0.743^{+0.282}_{-0.213}$	$0.743^{+0.0675}_{-0.0627}$	$0.530 \leq X \leq 1$
$N\Omega$ dibaryon	$1.40^{+1.20}_{-0.600}$	$1.40^{+0.523}_{-0.364}$	$0.801 \leq X \leq 1$
$\Omega\Omega$ dibaryon	$1.56^{+1.95}_{-0.773}$	$1.56^{+1.22}_{-0.626}$	$0.791 \leq X \leq 1$
${}^3_{\Lambda}\text{H}$	$1.35^{+0.531}_{-0.366}$	$1.35^{+1.241}_{-0.603}$	$0.745 \leq X \leq 1$
${}^4\text{He}$ dimer	$1.08^{+0.179}_{-0.152}$	$1.08^{+0.129}_{-0.115}$	$0.926 \leq X \leq 1$

Effective field theory

Low-energy scattering with near-threshold bound state

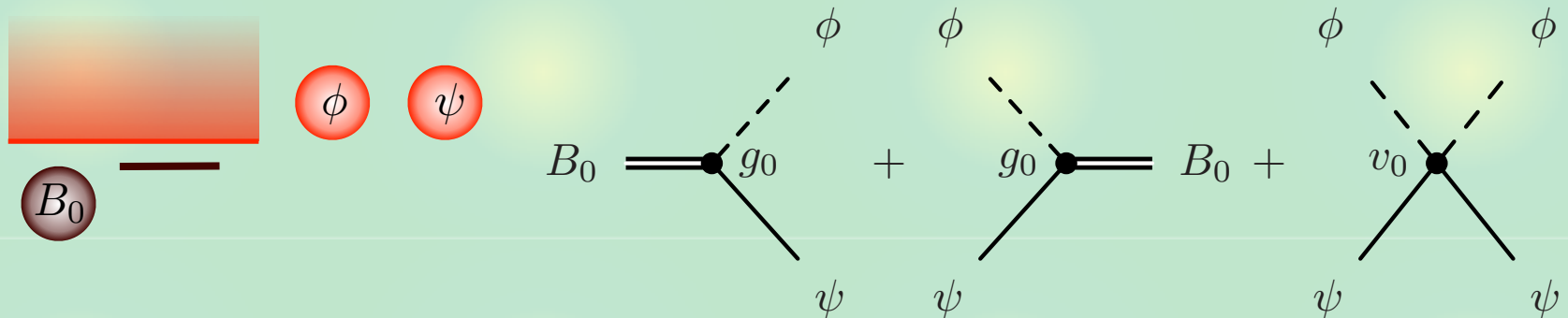
- Nonrelativistic EFT with contact interaction

D.B. Kaplan, Nucl. Phys. B494, 471 (1997)

E. Braaten, M. Kusunoki, D. Zhang, Annals Phys. 323, 1770 (2008)

$$H_{\text{free}} = \int d\mathbf{r} \left[\frac{1}{2M} \nabla \psi^\dagger \cdot \nabla \psi + \frac{1}{2m} \nabla \phi^\dagger \cdot \nabla \phi + \frac{1}{2M_0} \nabla B_0^\dagger \cdot \nabla B_0 + \omega_0 B_0^\dagger B_0 \right]$$

$$H_{\text{int}} = \int d\mathbf{r} \left[g_0 \left(B_0^\dagger \phi \psi + \psi^\dagger \phi^\dagger B_0 \right) + v_0 \psi^\dagger \phi^\dagger \phi \psi \right]$$



- **cutoff**: $\Lambda \sim 1/R_{\text{typ}}$ (interaction range of microscopic theory)

- At low momentum $p \ll \Lambda$, interaction \sim contact

Compositeness and “elementarity”

Eigenstates

$$H_{\text{free}} |B_0\rangle = \omega_0 |B_0\rangle, \quad H_{\text{free}} |\mathbf{p}\rangle = \frac{\mathbf{p}^2}{2\mu} |\mathbf{p}\rangle$$

free (discrete + continuum)

$$(H_{\text{free}} + H_{\text{int}}) |B\rangle = -B |B\rangle$$

full (bound state)

- normalization of $|B\rangle$ + completeness relation

$$\langle B|B\rangle = 1, \quad 1 = |B_0\rangle\langle B_0| + \int \frac{d\mathbf{p}}{(2\pi)^3} |\mathbf{p}\rangle\langle \mathbf{p}|$$

- projections onto free eigenstates

$$1 = Z + X, \quad Z \equiv |\langle B_0|B\rangle|^2, \quad X \equiv \int \frac{d\mathbf{p}}{(2\pi)^3} |\langle \mathbf{p}|B\rangle|^2$$

“elementarity”



compositeness

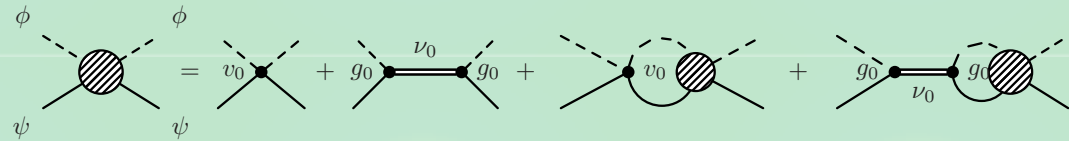


Z, X : real and nonnegative \rightarrow interpreted as **probability**

Weak binding relation

$\psi\phi$ scattering amplitude (exact result)

$$f(E) = -\frac{\mu}{2\pi} \frac{1}{[v(E)]^{-1} - G(E)}$$



$$v(E) = v_0 + \frac{g_0^2}{E - \omega_0}, \quad G(E) = \frac{1}{2\pi^2} \int_0^\Lambda dp \frac{p^2}{E - p^2/(2\mu) + i0^+}$$

Compositeness $X \leftarrow v(E), G(E)$

$$X = \frac{G'(-B)}{G'(-B) - [1/v(-B)]'}$$

$1/R = \sqrt{2\mu B}$ expansion of scattering length a_0

$$a_0 = -f(E=0) = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\} \text{renormalization dependent}$$

renormalization independent

If $R \gg R_{\text{typ}}$, correction terms neglected: $X \leftarrow (a_0, B)$

Inclusion of decay channel

Introduce decay channel

$$H'_{\text{free}} = \int d\mathbf{r} \left[\frac{1}{2M'} \nabla \psi'^{\dagger} \cdot \nabla \psi' - \nu_{\psi} \psi'^{\dagger} \psi' + \frac{1}{2m'} \nabla \phi'^{\dagger} \cdot \nabla \phi' - \nu_{\phi} \phi'^{\dagger} \phi' \right]$$

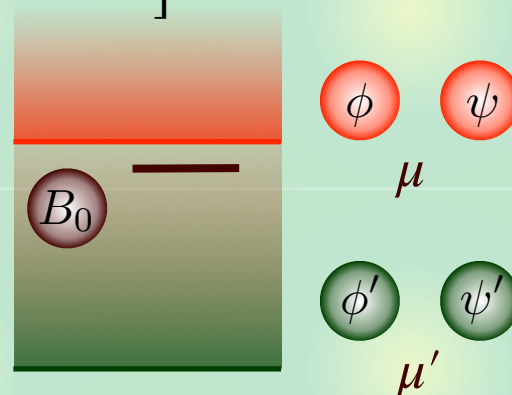
$$H'_{\text{int}} = \int d\mathbf{r} \left[g'_0 \left(B_0^{\dagger} \phi' \psi' + \psi'^{\dagger} \phi'^{\dagger} B_0 \right) + v'_0 \psi'^{\dagger} \phi'^{\dagger} \phi' \psi' + v_0^t (\psi'^{\dagger} \phi'^{\dagger} \phi' \psi' + \psi'^{\dagger} \phi'^{\dagger} \phi \psi) \right]$$

Quasi-bound state : complex eigenvalue

$$H = H_{\text{free}} + H'_{\text{free}} + H_{\text{int}} + H'_{\text{int}}$$

$$H |h\rangle = E_h |h\rangle, \quad E_h \in \mathbb{C}$$

$$\nu_{\psi} + \nu_{\phi} = \nu$$



Generalized relation : correction from threshold difference

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O} \left(\left| \frac{R_{\text{typ}}}{R} \right| \right) + \mathcal{O} \left(\left| \frac{\ell}{R} \right|^3 \right) \right\}, \quad R = \frac{1}{\sqrt{-2\mu E_h}}, \quad \ell \equiv \frac{1}{\sqrt{2\mu\nu}}$$

Y. Kamiya, T. Hyodo, PRC93, 035203 (2016); PTEP2017, 023D02 (2017)

If $|R| \gg (R_{\text{typ}}, \ell)$, **correction terms neglected:** $X \leftarrow (a_0, E_h)$

Evaluation of compositeness

Generalized weak-binding relation

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O} \left(\left| \frac{R_{\text{typ}}}{R} \right| \right) + \mathcal{O} \left(\left| \frac{\ell}{R} \right|^3 \right) \right\}, \quad R = \frac{1}{\sqrt{-2\mu E_h}}, \quad \ell \equiv \frac{1}{\sqrt{2\mu\nu}}$$

(a_0, E_h) determinations by several groups

- neglecting correction terms:

	E_h [MeV]	a_0 [fm]	$X_{\bar{K}N}$	$\tilde{X}_{\bar{K}N}$	$U/2$
Set 1 [35]	$-10 - i26$	$1.39 - i0.85$	$1.2 + i0.1$	1.0	0.3
Set 2 [36]	$-4 - i8$	$1.81 - i0.92$	$0.6 + i0.1$	0.6	0.0
Set 3 [37]	$-13 - i20$	$1.30 - i0.85$	$0.9 - i0.2$	0.9	0.1
Set 4 [38]	$2 - i10$	$1.21 - i1.47$	$0.6 + i0.0$	0.6	0.0
Set 5 [38]	$-3 - i12$	$1.52 - i1.85$	$1.0 + i0.5$	0.8	0.3

- In all cases, $X \sim 1$ with small $U/2$ (complex nature)

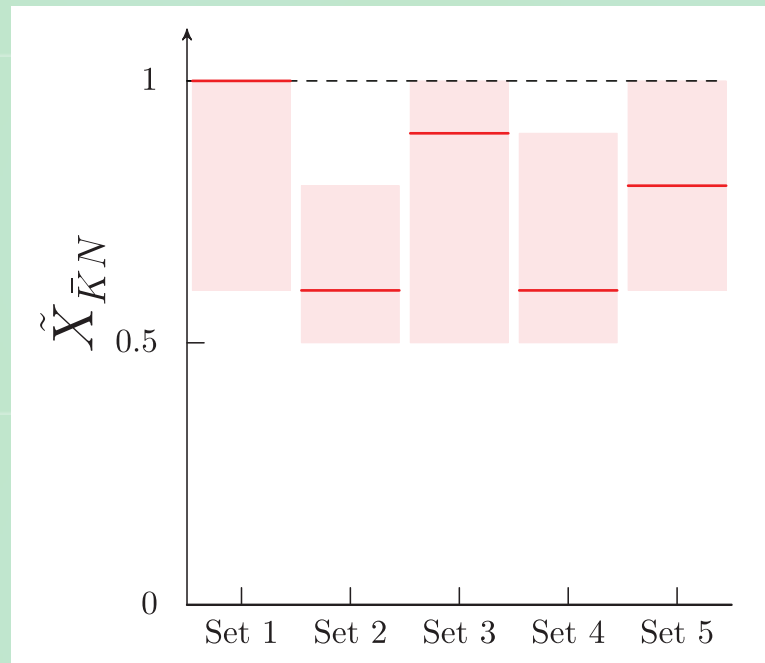
$\Lambda(1405)$: $\bar{K}N$ composite dominance \leftarrow observables

Uncertainty estimation

Estimation of correction terms: $|R| \sim 2$ fm

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\left|\frac{R_{\text{typ}}}{R}\right|\right) + \mathcal{O}\left(\left|\frac{\ell}{R}\right|^3\right) \right\}, \quad R = \frac{1}{\sqrt{-2\mu E_h}}, \quad \ell \equiv \frac{1}{\sqrt{2\mu\nu}}$$

- ρ meson exchange picture: $R_{\text{typ}} \sim 0.25$ fm
- energy difference from $\pi\Sigma$: $\ell \sim 1.08$ fm



$\bar{K}N$ composite dominance holds even **with correction terms.**

Summary

Pole structure of the $\Lambda(1405)$ region is now well constrained: “ $\Lambda(1405)$ ” \rightarrow $\Lambda(1405)$ **and $\Lambda(1380)$.**

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881, 98 (2012);

P.A. Zyla, *et al.* (Particle Data Group), PTEP 2020, 083C01 (2020)

T. Hyodo, M. Niiyama, PPNP 120, 103868 (2021);

T. Hyodo, W. Weise, arXiv:2202.06181 [nucl-th]

With the uncertainty estimation and range correction, deuteron is **quantitatively shown to be composite.**

T. Kinugawa, T. Hyodo, arXiv:2111.06619; 2112.00249; 2201.04283 [hep-ph]

Generalized weak-binding relation shows that $\Lambda(1405)$ is dominated by **molecular $\bar{K}N$ state.**

Y. Kamiya, T. Hyodo, PRC93, 035203 (2016); PTEP2017, 023D02 (2017)