

Hadron-hadron interaction from heavy-ion results





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2021, Nov. 23rd ₁

Contents

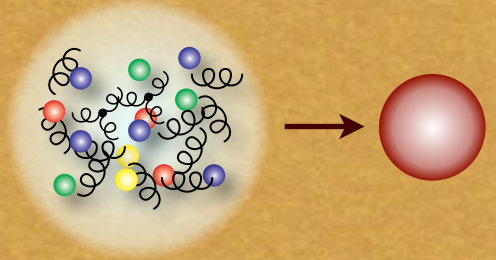
 Introduction — T_{cc} and $X(3872)$

 Production yields (1-hadron detection)

[S. Cho *et al.*, ExHIC collaboration, PRL 106, 212001 \(2011\);](#)

[S. Cho *et al.*, ExHIC collaboration, PRC 84, 064910 \(2011\);](#)

[S. Cho *et al.*, ExHIC collaboration, PPNP 95, 279 \(2017\)](#)



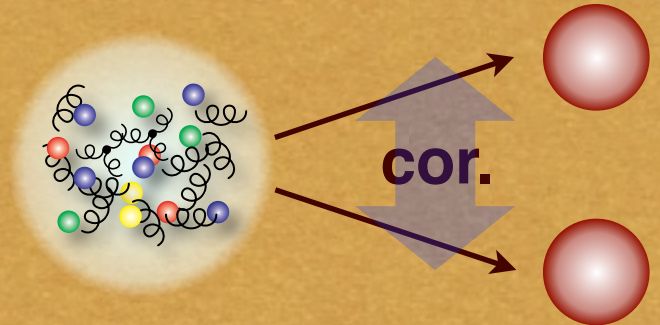
 Correlation functions (2-hadron detection)

- K^-p interactions

[Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise, PRL124, 132501 \(2020\)](#)

- D meson sector

[Y. Kamiya, T. Hyodo, A. Ohnishi, in preparation](#)



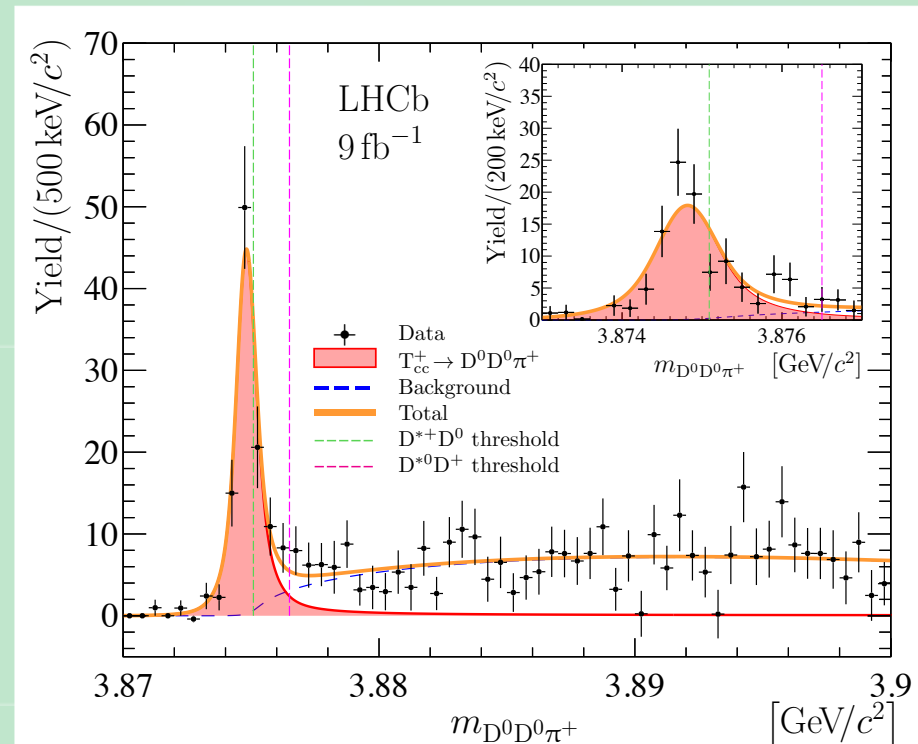
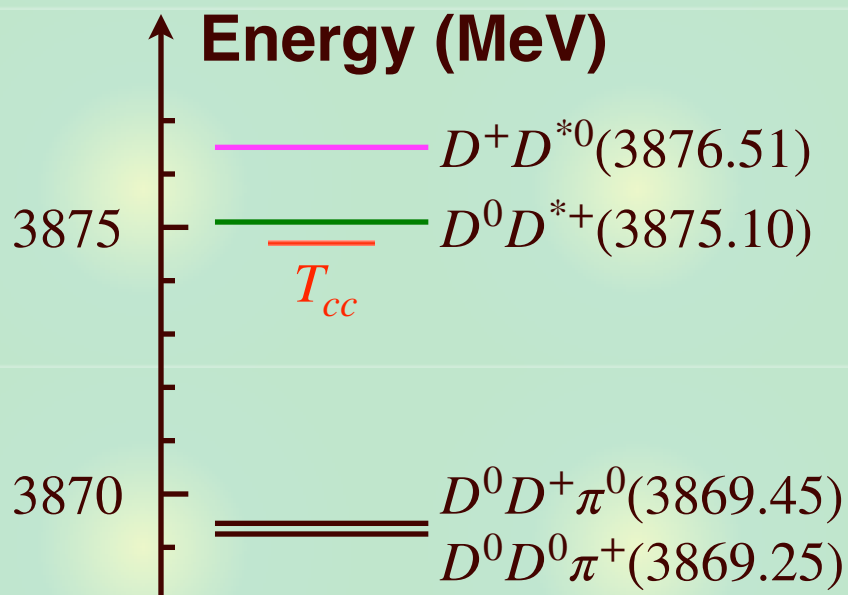
 Summary

Observation of T_{cc}

T_{cc} observed in $D^0 D^0 \pi^+$ spectrum

LHCb collaboration, arXiv 2109.01038 [hep-ex], 2109.01056 [hep-ex]

- Signal near DD^* threshold
- Charm $C = +2$: $\sim cc\bar{u}\bar{d}$
- Level structure

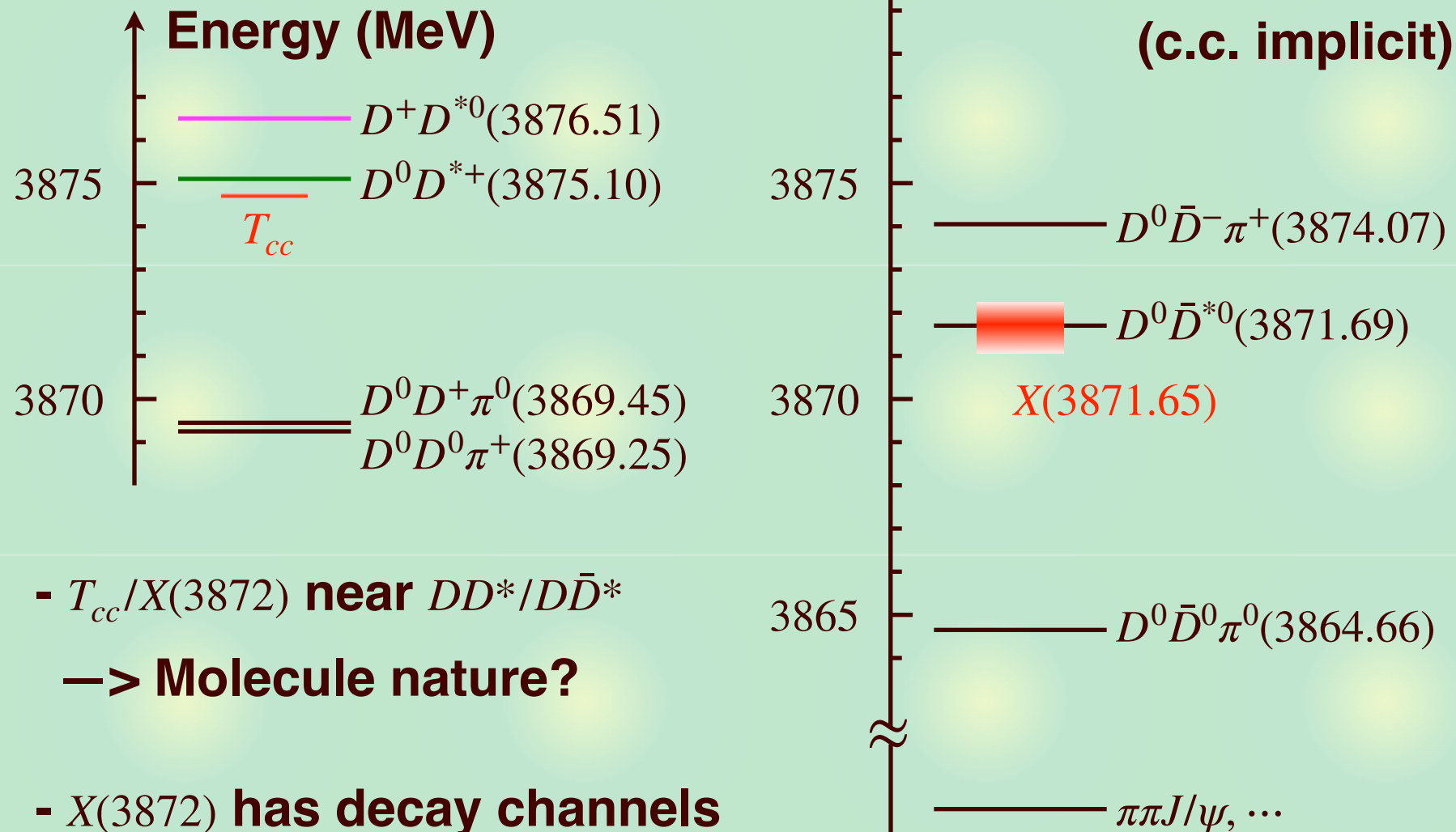


- Very small (**few MeV ~ keV**) energy scale involved

T_{cc} and $X(3872)$

$X(3872)$: another near-threshold state with $M_{T_{cc}} \sim M_{X(3872)}$

- Masses from PDG Live



ExHIC collaboration

Hadron production yields and internal structure

PRL **106**, 212001 (2011) PHYSICAL REVIEW LETTERS week ending 27 MAY 2011

Identifying Multiquark Hadrons from Heavy Ion Collisions

Sungtae Cho,¹ Takenori Furumoto,^{2,3} Tetsuo Hyodo,⁴ Daisuke Jido,² Che Ming Ko,⁵ Su Hong Lee,^{1,2} Marina Nielsen,⁶ Akira Ohnishi,² Takayasu Sekihara,^{2,7} Shigehiro Yasui,⁸ and Koichi Yazaki^{2,3}

(ExHIC Collaboration)

PHYSICAL REVIEW C **84**, 064910 (2011)

Exotic hadrons in heavy ion collisions

Sungtae Cho,¹ Takenori Furumoto,^{2,3} Tetsuo Hyodo,⁴ Daisuke Jido,² Che Ming Ko,⁵ Su Hong Lee,¹ Marina Nielsen,⁶ Akira Ohnishi,² Takayasu Sekihara,^{2,7} Shigehiro Yasui,⁸ and Koichi Yazaki^{2,9}

(ExHIC Collaboration)

Progress in Particle and Nuclear Physics 95 (2017) 279–322

Contents lists available at ScienceDirect

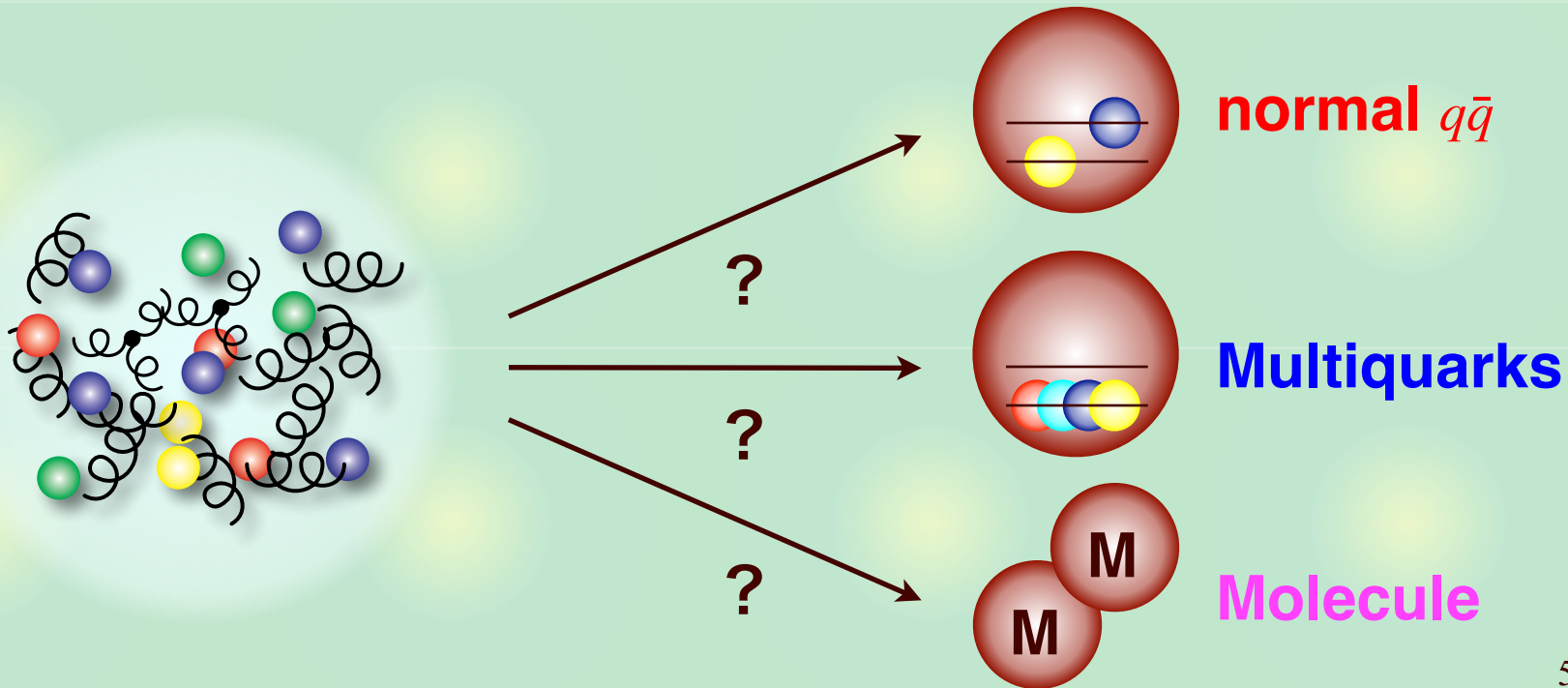
Progress in Particle and Nuclear Physics

journal homepage: www.elsevier.com/locate/ppnp

Review

Exotic hadrons from heavy ion collisions^{*}

Sungtae Cho^a, Tetsuo Hyodo^b, Daisuke Jido^c, Che Ming Ko^d, Su Hong Lee^{e,*}, Saori Maeda^f, Kenta Miyahara^g, Kenji Morita^b, Marina Nielsen^h, Akira Ohnishi^b, Takayasu Sekiharaⁱ, Taesoo Song^j, Shigehiro Yasui^f, Koichi Yazaki^k (ExHIC Collaboration)



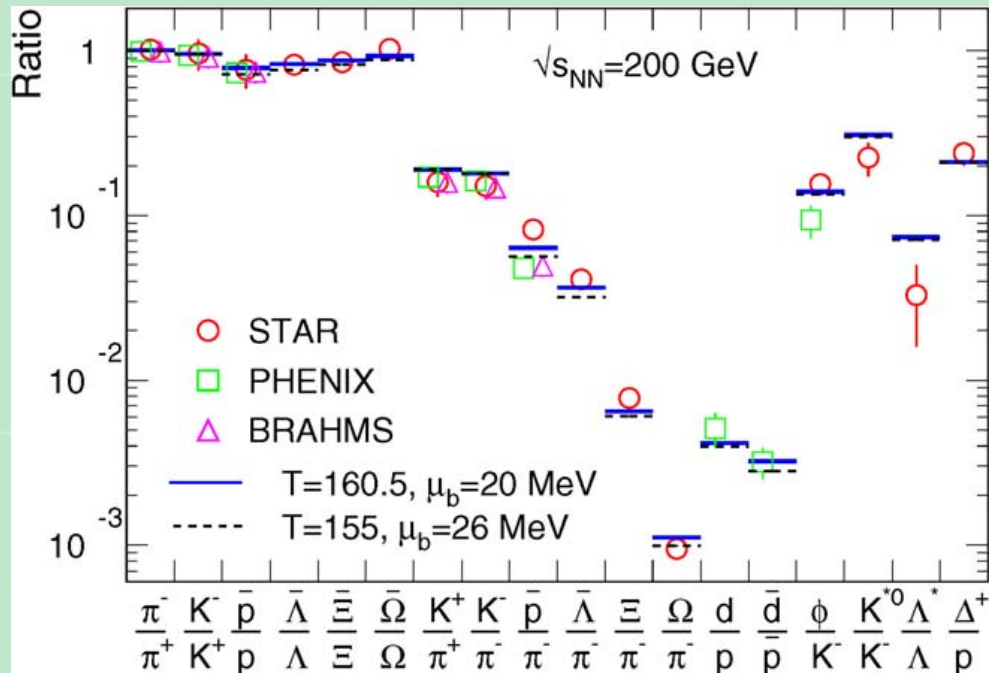
Statistical model

Statistical production yield for hadron h

A. Andronic, P. Braun-Munzinger, J. Stachel, NPA 772, 167 (2006)

$$N_h^{\text{stat}} = V_H \frac{g_h}{2\pi^2} \int_0^\infty \frac{p^2 dp}{\gamma_h^{-1} e^{E_h/T_H} \pm 1}$$

- Fugacity γ_h : chemical equilibrium for u, d, s , tuned for c, b



- works well for normal hadrons

Coalescence model

Coalescence (overlap of **constituents** and **hadron w.f.**)

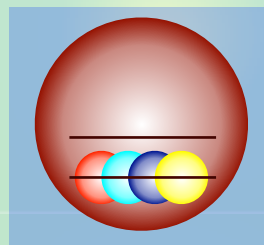
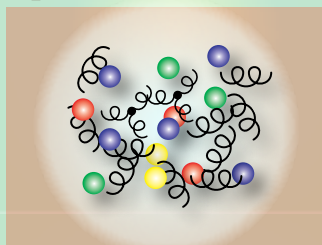
V. Greco, C. M. Ko, P. Levai, PRL 90, 202302 (2003)

$$N_h^{\text{coal}} = g_h \int \prod_{i=1}^n \left[\frac{1}{g_i} \frac{p_i \cdot d\sigma_i}{(2\pi)^3} \frac{d^3 p_i}{E_i} f(x_i, p_i) \right] f^W(x_1, \dots, x_n; p_1, \dots, p_n)$$

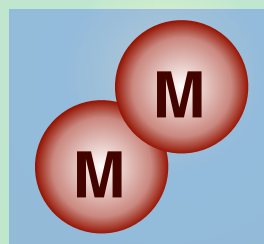
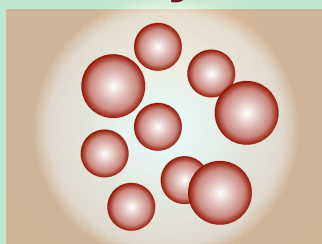
- Model parameters ← statistical yields of normal hadrons

Prediction for exotic hadrons

- **Multiquarks** by quark coalescence



- **Hadronic molecule** by hadron coalescence



Yield estimation

Coalescence-statistical ratio

S. Cho et al., ExHIC collaboration, PPNP 95, 279 (2017)

$$R_h^{CS} = \frac{N_h^{coal}}{N_h^{stat}}$$

- **Multiquarks suppressed**
- **Molecules enhanced**

Yields (Pb-Pb @ 5.02 TeV)

$$N_N = 32$$

$$N_{T_{cc}}^{stat} = 4.4 \times 10^{-3}$$

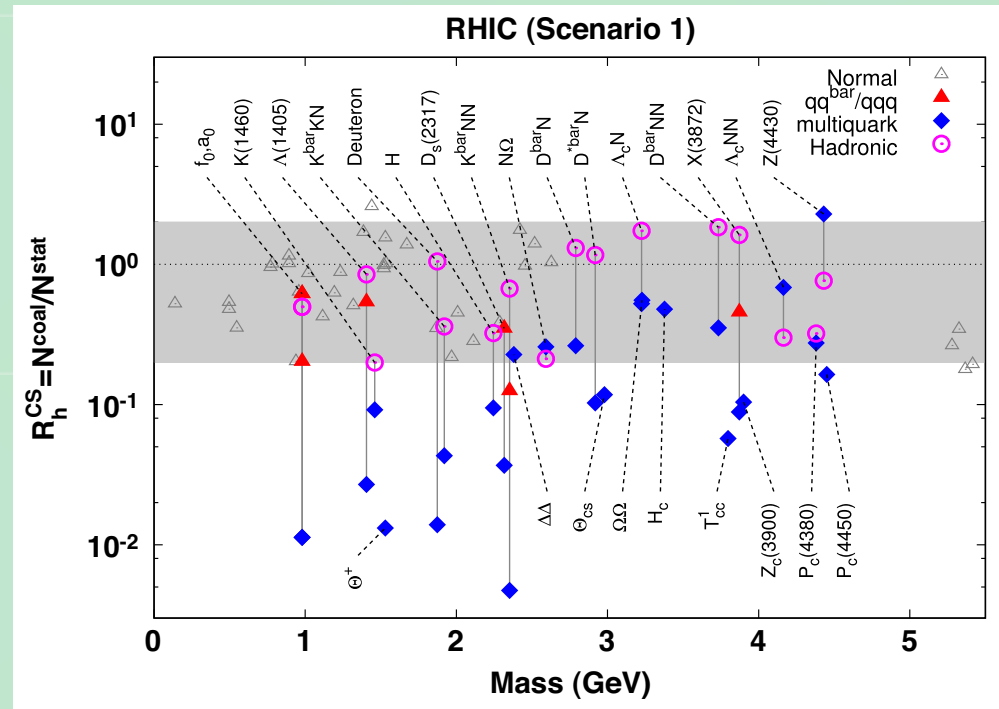
$$N_{T_{cc}}^{cc\bar{u}\bar{d}} = (1.8 - 2.1) \times 10^{-4}$$

$$N_{X(3872)}^{stat} = 2.8 \times 10^{-3}$$

$$N_{X(3872)}^{c\bar{c}q\bar{q}} = (1.8 - 2.1) \times 10^{-4}$$

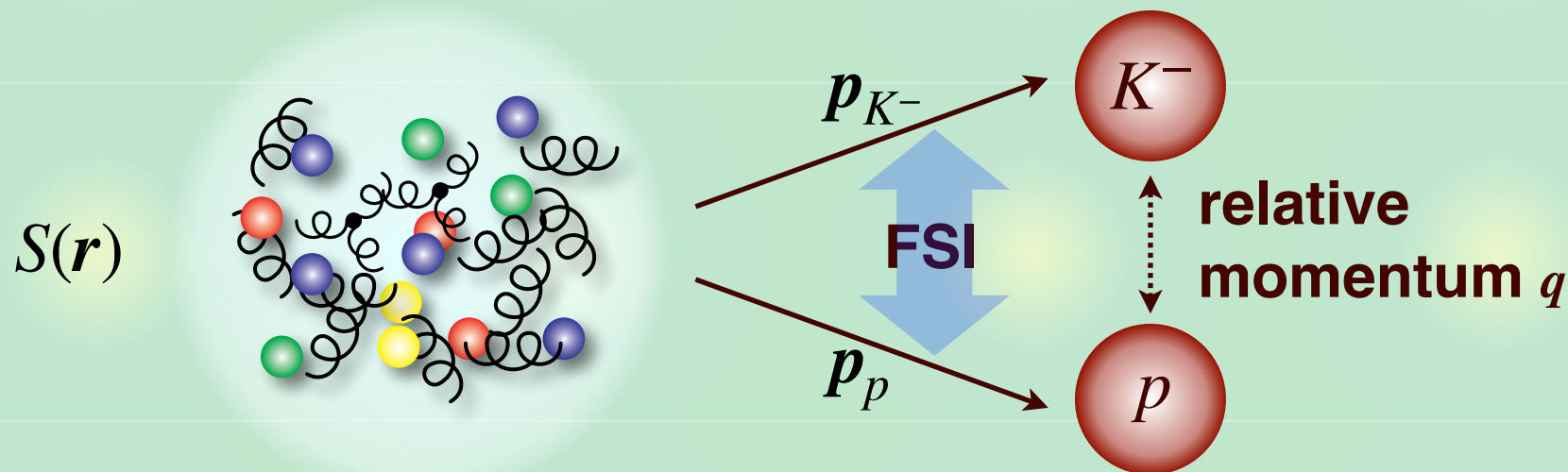
$$N_{X(3872)}^{molecule} = (4.5) \times 10^{-3}$$

If T_{cc} is a DD^* molecule, yield would be $\sim 10^{-3}$



Correlation function and hadron interaction

High-energy collision: chaotic source $S(\mathbf{r})$ of hadron emission



- Definition

$$C(\mathbf{q}) = \frac{N_{K^-p}(\mathbf{p}_{K^-}, \mathbf{p}_p)}{N_{K^-}(\mathbf{p}_{K^-})N_p(\mathbf{p}_p)} \quad (= 1 \text{ in the absence of FSI/QS})$$

- Theory (Koonin-Pratt formula)

$$C(\mathbf{q}) \simeq \int d^3\mathbf{r} S(\mathbf{r}) |\Psi_q^{(-)}(\mathbf{r})|^2$$

Source function \longleftrightarrow two-body wave function (FSI)

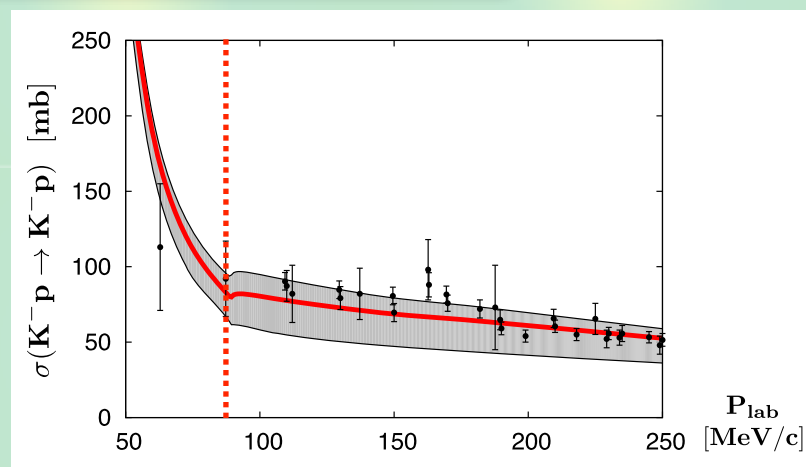
ALICE collaboration, *Nature* 588, 232 (2020); ...

Experimental data of K^-p correlation

K^-p total cross sections

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011)

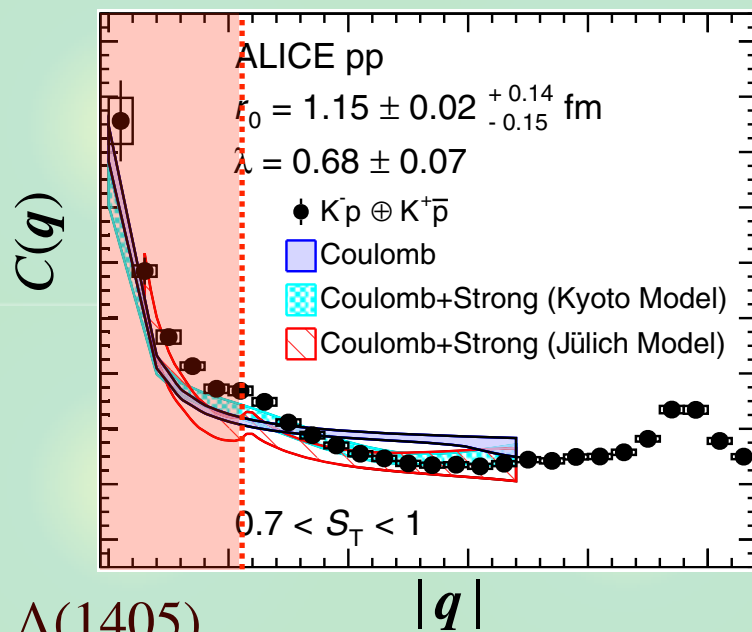
- Old bubble chamber data
- Resolution is not good
- Threshold cusp is not visible



K^-p correlation function

ALICE collaboration, PRL 124, 092301 (2020)

- Excellent **precision** (\bar{K}^0_n cusp)
- Low-energy data **below** \bar{K}^0_n



—> Important constraint on $\bar{K}N$ and $\Lambda(1405)$

Coupled-channel correlation function

Schrödinger equation (s-wave)

$$\begin{pmatrix} -\frac{\nabla^2}{2\mu_1} + V_{11}(r) + V_C(r) & V_{12}(r) & \dots \\ V_{21}(r) & -\frac{\nabla^2}{2\mu_2} + V_{22}(r) + \Delta_2 & \dots \\ \vdots & \vdots & \ddots \end{pmatrix} \begin{pmatrix} \psi_{K^-p}(r) \\ \psi_{\bar{K}^0n}(r) \\ \vdots \end{pmatrix} = E \begin{pmatrix} \psi_{K^-p}(r) \\ \psi_{\bar{K}^0n}(r) \\ \vdots \end{pmatrix}$$

Coulomb

threshold energy difference

Coupled-channel formulation

R. Lednicky, V.V. Lyuboshitz, V.L. Lyuboshitz, *Phys. Atom. Nucl.* **61**, 2050 (1997);
J. Haidenbauer, *NPA* **981**, 1 (2019)

$$C_{K^-p}(q) \simeq \int d^3\mathbf{r} S_{K^-p}(\mathbf{r}) |\Psi_{K^-p,q}^{(-)}(\mathbf{r})|^2 + \sum_{i \neq K^-p} \omega_i \int d^3\mathbf{r} S_i(\mathbf{r}) |\Psi_{i,q}^{(-)}(\mathbf{r})|^2$$

- **Transition from** $\bar{K}^0n, \pi^+\Sigma^-, \pi^0\Sigma^0, \pi^-\Sigma^+, \pi^0\Lambda$

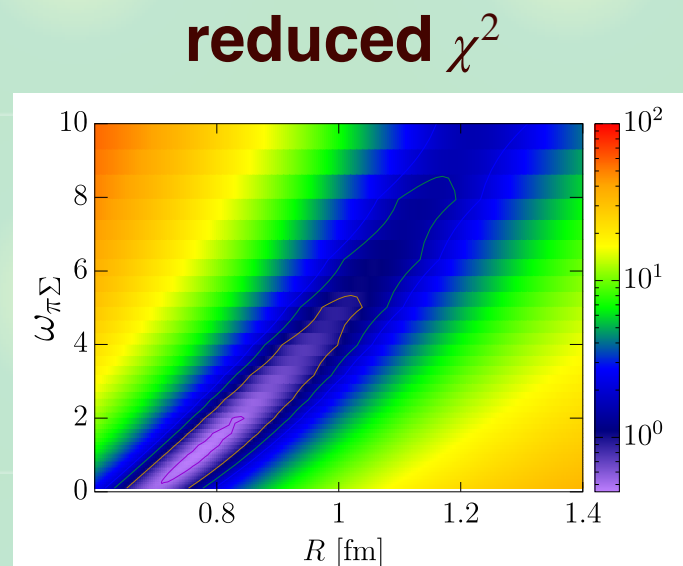
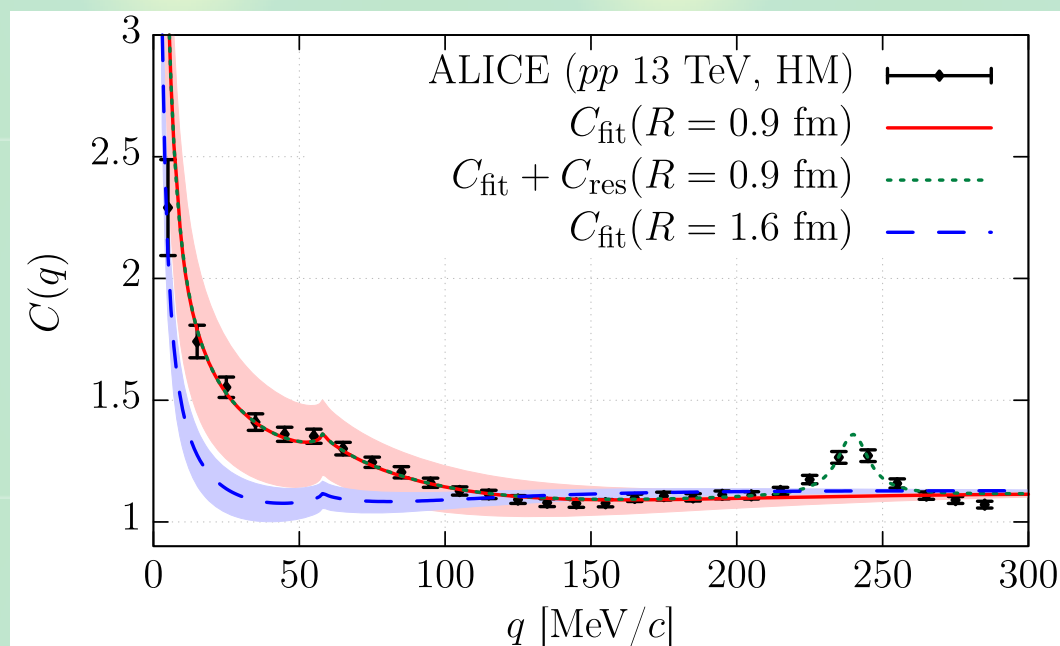
- ω_i : **weight of source channel** i relative to K^-p

Correlation from chiral SU(3) dynamics

Wave function $\Psi_q^{(-)}(r)$: coupled-channel $\bar{K}N-\pi\Sigma-\pi\Lambda$ potential

K. Miyahara, T. Hyodo, W. Weise, PRC98, 025201 (2018)

- Source function $S(r)$: Gaussian, $R \sim 1$ fm in K^+p data
- Source weight $\omega_{\pi\Sigma} \sim 2$ by simple statistical model estimate



Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise, PRL124, 132501 (2020)

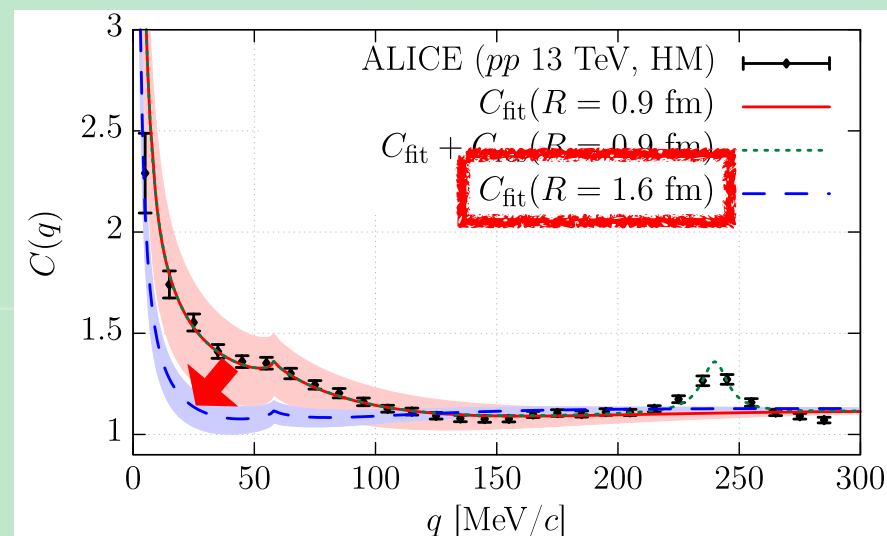
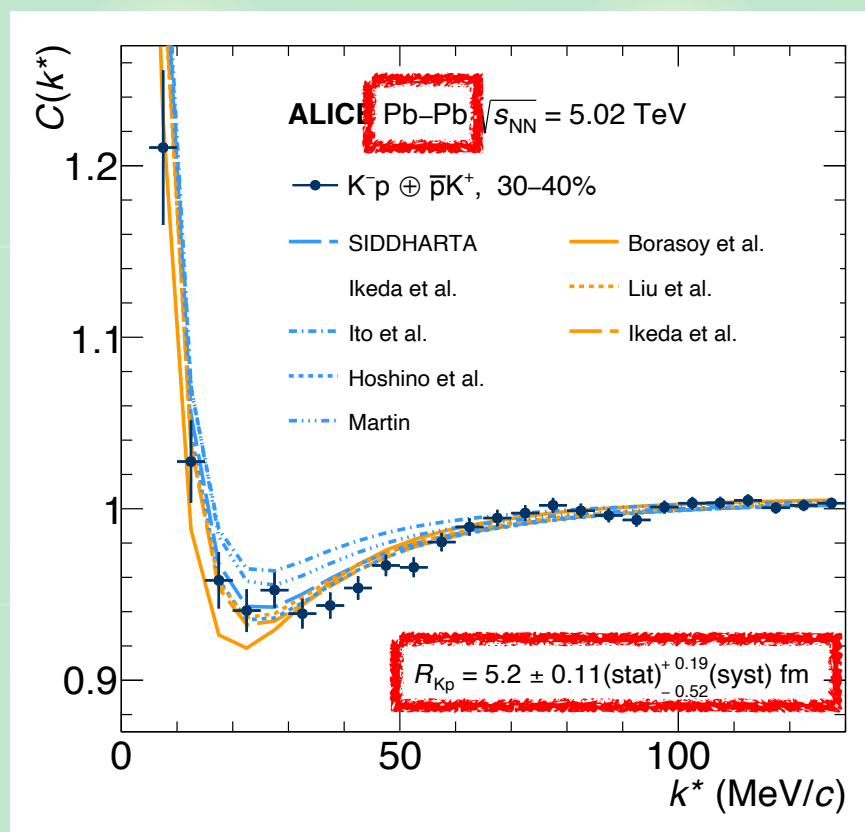
Correlation function by ALICE is well reproduced

Source size dependence

New data of Pb-Pb collisions at 5.02 TeV

ALICE collaboration, PLB 822, 136708 (2021)

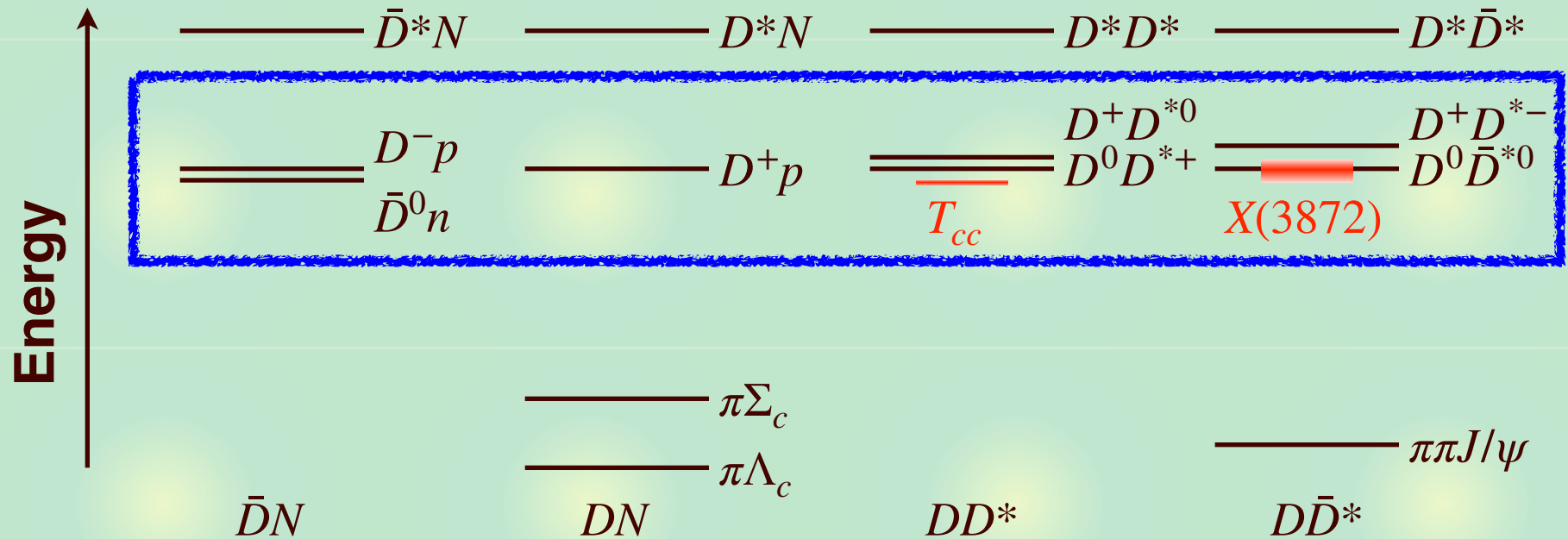
- Scattering length $a_{K^-p} = -0.91 + 0.92i$ fm



Correlation is suppressed at larger R , as predicted

Strategy

Schematic threshold structures in D meson sector



One-range gaussian potentials

$$V(r) = V_0 \exp\{-m^2 r^2\}$$

- $V_0 < -$ Scattering length in theoretical models ($\bar{D}N, DN$)
- $V_0 < -$ Binding energies of $T_{cc}, X(3872)$ ($DD^*, D\bar{D}^*$)

Exotic DN sector

D^-p correlation functions ($\bar{c}duud$, exotic)

- Coupled with \bar{D}^0n
- No decay channels below
- Theoretical models

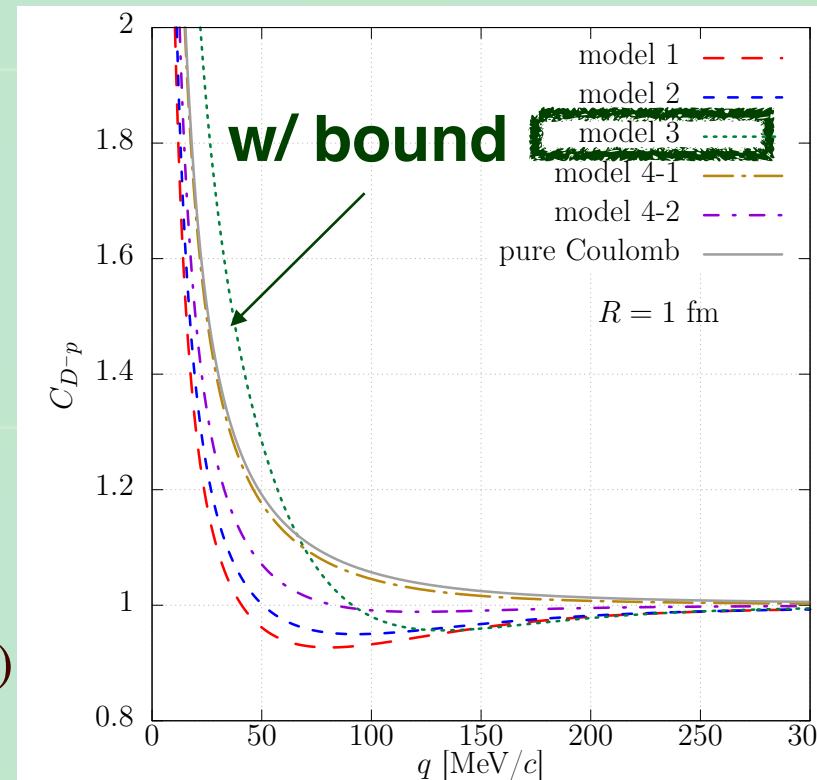
[1] J. Hofmann, M.F.M. Lutz, NPA763, 90 (2005);

[2] J. Haidenbauer *et al.*, EPJA33, 107 (2007);

[3] Y. Yamaguchi *et al.*, PRD84, 014032 (2011);

[4] C. Fontoura *et al.*, PRD87, 025206 (2013)

- Gaussian potentials with $a_0(I = 0,1)$



- Model 3 with a **bound state** : dip structure
- To be compared with experiments in future

Y. Kamiya, T. Hyodo, A. Ohnishi, in preparation

Non-exotic $\bar{D}N$ sector

D^+p correlation functions ($c\bar{d}uud$, non-exotic)

- No isospin partner in DN
- With decay channels ($\pi\Lambda_c, \pi\Sigma_c$)
- Theoretical models

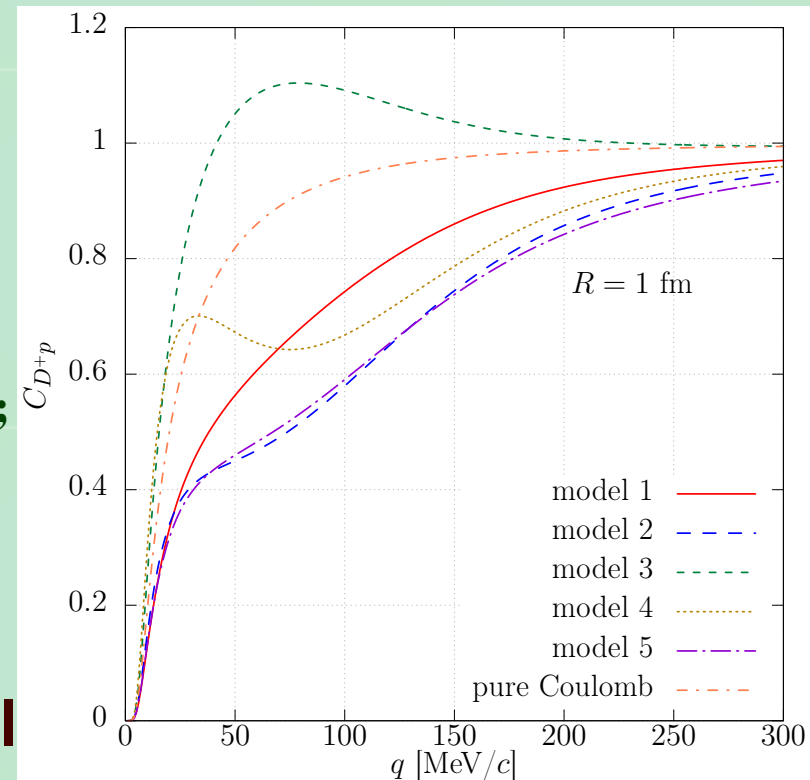
[1] J. Hofmann, M.F.M. Lutz, NPA763, 90 (2005);
 [2] T. Mizutani, A. Ramos, PRC74, 065201 (2006);
 [3] C. Garcia-Recio *et al.*, PRD79, 054004 (2009);
 [4] J. Haidenbauer *et al.*, EPJA47, 18 (2011);
 [5] U. Raha *et al.*, PRC98, 034002 (2018)

- Effective single-channel potential

$$\leftarrow a_0(I = 1)$$

- Sizable dependence on the scattering length

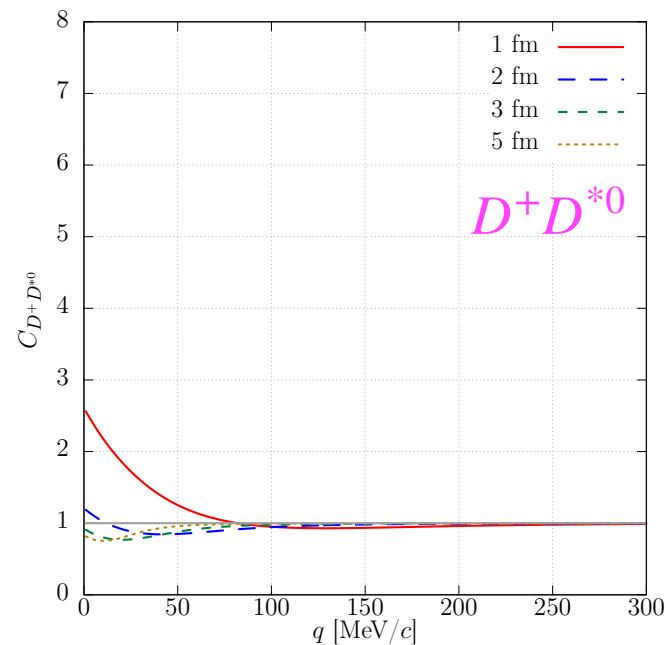
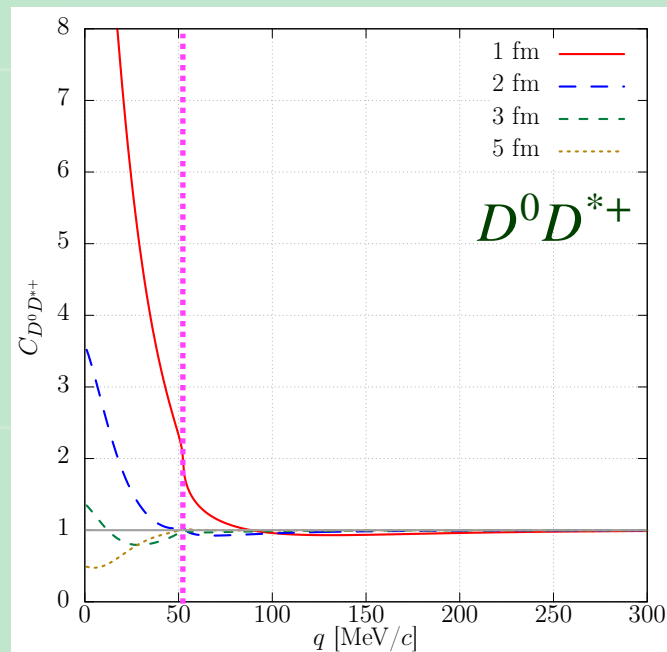
Y. Kamiya, T. Hyodo, A. Ohnishi, in preparation



$DD^* \sim T_{cc}$ sector

D^0D^{*+} and D^+D^{*0} correlation functions ($cc\bar{u}\bar{d}$, exotic)

— D^+D^{*0}
— D^0D^{*+}
— T_{cc}



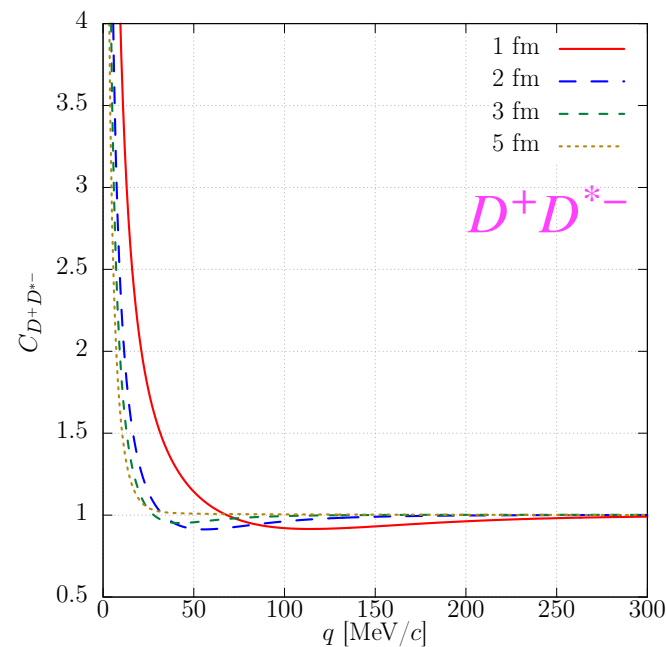
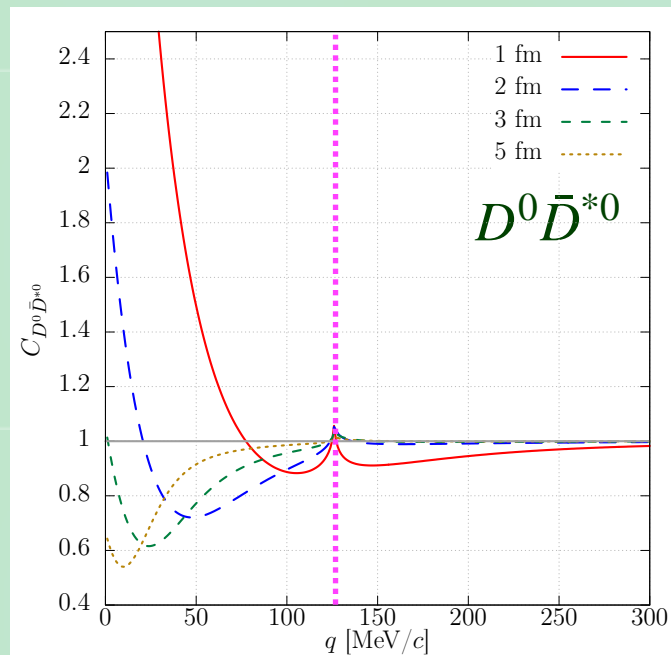
- Bound state feature (source size dep.) in both channels
- Strong signal in D^0D^{*+} , weaker one in D^+D^{*0}
- D^+D^{*0} cusp in D^0D^{*+} ($q \sim 52$ MeV) is not very prominent

$D\bar{D}^* \sim X(3872)$ sector

$D^0\bar{D}^{*0}$ and D^+D^{*-} correlation functions ($c\bar{c}q\bar{q}$)

D^+D^{*-}

$D^0\bar{D}^{*0}$
 $X(3872)$



- Bound state feature in $D^0\bar{D}^{*0}$ correlation
- Sizable D^+D^{*-} cusp in $D^0\bar{D}^{*0}$ ($q \sim 126$ MeV)
- D^+D^{*-} correlation : Coulomb attraction dominance

Summary

Hadron production yields ← internal structure

S. Cho *et al.*, ExHIC collaboration, PRL 106, 212001 (2011);

S. Cho *et al.*, ExHIC collaboration, PRC 84, 064910 (2011);

S. Cho *et al.*, ExHIC collaboration, PPNP 95, 279 (2017)

K^-p correlation in pp collisions

- well described by chiral SU(3) dynamics

Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise, PRL124, 132501 (2020)

D^-p and D^+p correlations

- sizable scattering length dependence

D^0D^{*+} and $D^0\bar{D}^{*0}$ correlations

- (quasi-)bound nature of T_{cc} and $X(3872)$

Y. Kamiya, T. Hyodo, A. Ohnishi, in preparation

ExHIC parameters

Table 3.1

Statistical and coalescence model parameters for Scenario 1 and 2 at RHIC (200 GeV), LHC (2.76 TeV) and LHC (5.02 TeV), and those given in Refs. [14,15]. Quark masses are taken to be $m_q = 350$ MeV, $m_s = 500$ MeV, $m_c = 1500$ MeV and $m_b = 4700$ MeV. In Refs. [14,15], light quark masses were taken to be $m_q = 300$ MeV.

	RHIC		LHC (2.76 TeV)		LHC (5.02 TeV)		RHIC	LHC (5 TeV)
	Sc. 1	Sc. 2	Sc. 1	Sc. 2	Sc. 1	Sc. 2	Refs [14,15]	
T_H (MeV)		162				156		175
V_H (fm ³)		2100				5380	1908	5152
μ_B (MeV)		24				0	20	0
μ_s (MeV)		10				0	10	0
γ_c		22		39		50	6.40	15.8
γ_b		4.0×10^7		8.6×10^8		1.4×10^9	2.2×10^6	3.3×10^7
T_C (MeV)	162	166	156	166	156	166		175
V_C (fm ³)	2100	1791	5380	3533	5380	3533	1000	2700
ω (MeV)	590	608	564	609	564	609		550
ω_s (MeV)	431	462	426	502	426	502		519
ω_c (MeV)	222	244	219	278	220	279		385
ω_b (MeV)	183	202	181	232	182	234		338
$N_u = N_d$	320	302	700	593	700	593	245	662
$N_s = N_{\bar{s}}$	183	176	386	347	386	347	150	405
$N_c = N_{\bar{c}}$		4.1		11		14	3	20
$N_b = N_{\bar{b}}$		0.03		0.44		0.71	0.02	0.8
T_F (MeV)		119				115		125
V_F (fm ³)		20355				50646	11322	30569
N_K		67.5				134	142 ^a	363 ^a
$N_{\bar{K}}$		59.6				134	127 ^a	363 ^a
N_N		20				32	62 ^a	150 ^a
N_{Δ}		18				28	–	–
N_{Λ}		3.8				6.5	–	–
N_{Ξ}		2.6				4.4	4.7	13
N_{Ω}		0.37				0.62	0.81	2.3
$N_D = N_{\bar{D}}$		1.5		4.0		5.2	1.0	6.9
$N_{D^*} = N_{\bar{D}^*}$		2.0		5.4		6.9	1.5	10
$N_{D_1} = N_{\bar{D}_1}$		0.20		0.49		0.63	0.19	1.3
$N_B = N_{\bar{B}}$		8.1×10^{-3}		0.12		0.20	5.3×10^{-3}	0.21
$N_{B^*} = N_{\bar{B}^*}$		1.9×10^{-2}		0.27		0.45	1.2×10^{-2}	0.49
N_{Λ_c}		0.17		0.36		0.46	–	–
N_{Σ_c}		0.2		0.41		0.52	–	–
$N_{\Sigma_c^*}$		0.28		0.56		0.71	–	–
N_{Ξ_c}		0.11		0.25		0.32	0.10	0.65

^a Values contain feed down contributions.

T_{cc} and $X(3872)$ yields

Particle	Scenario 1		Scenario 2		Mol.	Stat.
	$q\bar{q}/qqq$	Multiquark	$q\bar{q}/qqq$	Multiquark		
RHIC						
T_{cc}^1	–	5.0×10^{-5}	–	5.3×10^{-5}	–	8.9×10^{-4}
LHC (2.76 TeV)						
T_{cc}^1	–	1.1×10^{-4}	–	1.3×10^{-4}	–	2.7×10^{-3}
LHC (5.02 TeV)						
T_{cc}^1	–	1.8×10^{-4}	–	2.1×10^{-4}	–	4.4×10^{-3}

Particle	Scenario 1		Scenario 2		Mol.	Stat.
	$q\bar{q}/qqq$	Multiquark	$q\bar{q}/qqq$	Multiquark		
RHIC						
$D_s(2317)$	2.3×10^{-2}	2.4×10^{-3}	2.3×10^{-2}	2.5×10^{-3}	6.5×10^{-3}	6.6×10^{-2}
$X(3872)$	5.4×10^{-4}	5.0×10^{-5}	5.6×10^{-4}	5.3×10^{-5}	9.1×10^{-4}	5.7×10^{-4}
LHC (2.76 TeV)						
$D_s(2317)$	5.2×10^{-2}	4.3×10^{-3}	5.0×10^{-2}	4.5×10^{-3}	1.4×10^{-2}	1.5×10^{-1}
$X(3872)$	1.6×10^{-3}	1.1×10^{-4}	1.7×10^{-3}	1.3×10^{-4}	2.7×10^{-3}	1.7×10^{-3}
LHC (5.02 TeV)						
$D_s(2317)$	6.5×10^{-2}	5.4×10^{-3}	6.4×10^{-2}	5.7×10^{-3}	1.8×10^{-2}	1.9×10^{-1}
$X(3872)$	2.5×10^{-3}	1.8×10^{-4}	2.7×10^{-3}	2.1×10^{-4}	4.5×10^{-3}	2.8×10^{-3}

Boundary conditions

Asymptotic ($r \rightarrow \infty$) wave function

$$\begin{pmatrix} \psi_{K^{-p}}(r) \\ \psi_{\bar{K}^{0n}}(r) \\ \vdots \end{pmatrix} \propto \begin{pmatrix} \#e^{-iqr} + \#e^{iqr} \\ \#e^{-iq_2r} + \#e^{iq_2r} \\ \vdots \end{pmatrix} \quad \text{incoming + outgoing}$$

- Usual scattering: normalize incoming flux of beam

$$\begin{pmatrix} \psi_{K^{-p}}(r) \\ \psi_{\bar{K}^{0n}}(r) \\ \vdots \end{pmatrix} \propto \begin{pmatrix} e^{-iqr} + c_1^{(+)}e^{iqr} \\ + c_2^{(+)}e^{iq_2r} \\ \vdots \end{pmatrix} \quad \text{coefficient} \sim \text{S-matrix}$$

$$c_i^{(+)} \propto s_{1i}(q)$$

- Correlation function: normalize outgoing flux

$$\psi^{(-)} = \begin{pmatrix} \psi_{K^{-p}}(r) \\ \psi_{\bar{K}^{0n}}(r) \\ \vdots \end{pmatrix} \propto \begin{pmatrix} c_1^{(-)}e^{-iqr} + e^{iqr} \\ c_2^{(-)}e^{-iq_2r} \\ \vdots \end{pmatrix} \quad c_i^{(-)} \propto s_{1i}^{\dagger}(q)$$

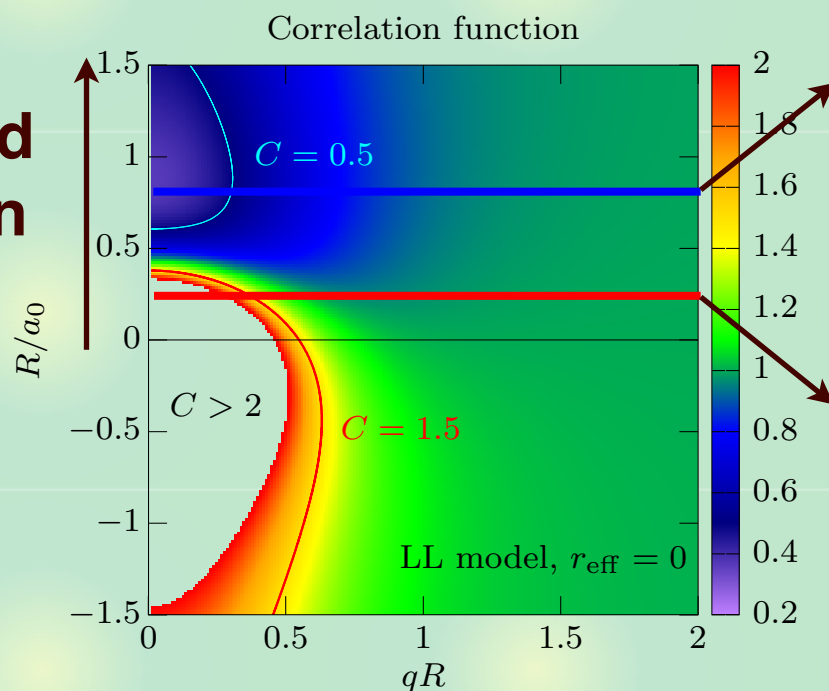
→ $\psi^{(-)}$ should be calculated with **full coupled channels**.

Bound state in correlation function

Y. Kamiya, K. Sasaki, T. Fukui, T. Hyodo, K. Morita, K. Ogata, A. Ohnishi, T. Hatsuda,
arXiv 2108.09644 [hep-ph]

Lednicky-Lyuboshitz model with $r_e = 0$

Bound region



- large source



- small source



Bound state \rightarrow small size enhancement & large size dip