

# $\Lambda(1405)$ as a hadronic molecule



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2020, Nov. 4th 1



## Introduction

- Structure of “unstable” resonance?



## Structure of $\Lambda(1405)$ resonance

- Pole positions

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881, 98 (2012);

P.A. Zyla, *et al.* (Particle Data Group), PTEP 2020, 083C01 (2020)

T. Hyodo, M. Niiyama, arXiv: 2010.07592 [hep-ph]

- $\bar{K}N$  compositeness

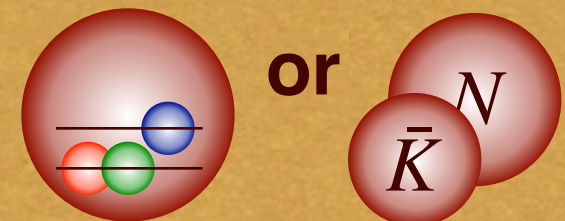
Y. Kamiya, T. Hyodo, PRC93, 035203 (2016); PTEP2017, 023D02 (2017)

T. Hyodo, JPS journal Vol. 75 No. 8, 478 (2020)

T. Kinugawa, T. Hyodo, in preparation



## Summary

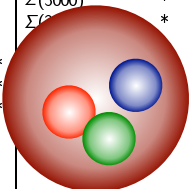


# Observed hadrons (2018)

PDG 2018 edition

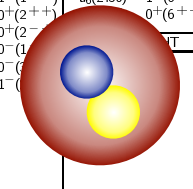
<http://pdg.lbl.gov/>

$p$	1/2 <sup>+</sup> ****	$\Delta(1232)$	3/2 <sup>+</sup> ****	$\Sigma^+$	1/2 <sup>+</sup> ****	$\Xi^0$	1/2 <sup>+</sup> ****	$\Lambda_c^+$	1/2 <sup>+</sup> ****
$n$	1/2 <sup>+</sup> ****	$\Delta(1600)$	3/2 <sup>+</sup> ***	$\Sigma^0$	1/2 <sup>+</sup> ****	$\Xi^-$	1/2 <sup>+</sup> ****	$\Lambda_c(2595)^+$	1/2 <sup>-</sup> ***
$N(1440)$	1/2 <sup>+</sup> ****	$\Delta(1620)$	1/2 <sup>-</sup> ****	$\Sigma^-$	1/2 <sup>+</sup> ****	$\Xi(1530)$	3/2 <sup>+</sup> ****	$\Lambda_c(2625)^+$	3/2 <sup>-</sup> ***
$N(1520)$	3/2 <sup>-</sup> ****	$\Delta(1700)$	3/2 <sup>-</sup> ****	$\Sigma(1385)$	3/2 <sup>+</sup> ****	$\Xi(1620)$	*	$\Lambda_c(2765)^+$	*
$N(1535)$	1/2 <sup>-</sup> ****	$\Delta(1750)$	1/2 <sup>+</sup> *	$\Sigma(1480)$	*	$\Xi(1690)$	***	$\Lambda_c(2880)^+$	5/2 <sup>+</sup> ***
$N(1650)$	1/2 <sup>-</sup> ****	$\Delta(1900)$	1/2 <sup>-</sup> **	$\Sigma(1560)$	*	$\Xi(1820)$	3/2 <sup>-</sup> ***	$\Lambda_c(2940)^+$	***
$N(1675)$	5/2 <sup>-</sup> ****	$\Delta(1905)$	5/2 <sup>+</sup> ****	$\Sigma(1580)$	3/2 <sup>-</sup> **	$\Xi(1950)$	***	$\Sigma_c(2455)$	1/2 <sup>+</sup> ****
$N(1680)$	5/2 <sup>+</sup> ****	$\Delta(1910)$	1/2 <sup>+</sup> ****	$\Sigma(1620)$	1/2 <sup>-</sup> *	$\Xi(2030)$	$\geq 5/2^?$ ***	$\Sigma_c(2520)$	3/2 <sup>+</sup> ****
$N(1685)$	*	$\Delta(1920)$	3/2 <sup>+</sup> ***	$\Sigma(1660)$	1/2 <sup>+</sup> ***	$\Xi(2120)$	**	$\Sigma_c(2800)$	***
$N(1700)$	3/2 <sup>-</sup> ***	$\Delta(1930)$	5/2 <sup>-</sup> ***	$\Sigma(1670)$	3/2 <sup>-</sup> ****	$\Xi(2250)$	**	$\Xi_c^+$	1/2 <sup>+</sup> ****
$N(1710)$	1/2 <sup>+</sup> ***	$\Delta(1940)$	3/2 <sup>-</sup> **	$\Sigma(1690)$	**	$\Xi(2370)$	**	$\Xi_c^0$	1/2 <sup>+</sup> ****
$N(1720)$	3/2 <sup>+</sup> ****	$\Delta(1950)$	7/2 <sup>+</sup> ****	$\Sigma(1730)$	3/2 <sup>+</sup> **	$\Xi(2500)$	*	$\Xi_c^-$	1/2 <sup>+</sup> ****
$N(1860)$	5/2 <sup>+</sup> **	$\Delta(2000)$	5/2 <sup>+</sup> **	$\Sigma(1750)$	1/2 <sup>-</sup> ***	$\Omega^-$	3/2 <sup>+</sup> ****	$\Xi_c^0$	1/2 <sup>+</sup> ****
$N(1875)$	3/2 <sup>-</sup> ***	$\Delta(2150)$	1/2 <sup>-</sup> *	$\Sigma(1770)$	1/2 <sup>+</sup> *	$\Omega(2250)^-$	***	$\Xi_c^+$	1/2 <sup>+</sup> ****
$N(1880)$	1/2 <sup>+</sup> **	$\Delta(2200)$	7/2 <sup>-</sup> *	$\Sigma(1775)$	5/2 <sup>-</sup> ****	$\Omega(2380)$	**	$\Xi_c^0$	1/2 <sup>+</sup> ****
$N(1895)$	1/2 <sup>-</sup> **	$\Delta(2300)$	9/2 <sup>+</sup> **	$\Sigma(1840)$	3/2 <sup>+</sup> **	$\Omega(2470)^-$	**	$\Xi_c^-$	1/2 <sup>+</sup> ****
$N(1900)$	3/2 <sup>+</sup> ***	$\Delta(2350)$	5/2 <sup>-</sup> *	$\Sigma(1880)$	1/2 <sup>+</sup> **			$\Xi_c^0$	1/2 <sup>+</sup> ****
$N(1990)$	7/2 <sup>+</sup> **	$\Delta(2390)$	7/2 <sup>+</sup> *	$\Sigma(1900)$	1/2 <sup>-</sup> *			$\Xi_c^-$	1/2 <sup>+</sup> ****
$N(2000)$	5/2 <sup>+</sup> **	$\Delta(2400)$	9/2 <sup>-</sup> **	$\Sigma(1915)$	5/2 <sup>+</sup> ****			$\Xi_c^0$	1/2 <sup>+</sup> ****
$N(2040)$	3/2 <sup>+</sup> **	$\Delta(2420)$	11/2 <sup>+</sup> ****	$\Sigma(1940)$	3/2 <sup>+</sup> **			$\Xi_c^-$	1/2 <sup>+</sup> ****
$N(2060)$	5/2 <sup>-</sup> **	$\Delta(2750)$	13/2 <sup>-</sup> **	$\Sigma(1940)$	3/2 <sup>-</sup> ***			$\Xi_c^0$	1/2 <sup>+</sup> ****
$N(2100)$	1/2 <sup>+</sup> *	$\Delta(2950)$	15/2 <sup>+</sup> **	$\Sigma(2000)$	1/2 <sup>-</sup> *			$\Xi_c^-$	1/2 <sup>+</sup> ****
$N(2120)$	3/2 <sup>-</sup> **			$\Sigma(2030)$	7/2 <sup>+</sup> ****			$\Omega_c^0$	1/2 <sup>+</sup> ****
$N(2190)$	7/2 <sup>-</sup> ****	$\Lambda$	1/2 <sup>+</sup> ****	$\Sigma(2070)$	5/2 <sup>+</sup> *			$\Omega_c(2770)^0$	3/2 <sup>+</sup> ****
$N(2220)$	9/2 <sup>+</sup> ****	$\Lambda(1405)$	1/2 <sup>-</sup> ****	$\Sigma(2080)$	3/2 <sup>+</sup> **				
$N(2250)$	9/2 <sup>-</sup> ****	$\Lambda(1520)$	3/2 <sup>-</sup> ****	$\Sigma(2100)$	7/2 <sup>-</sup> *			$\Xi_{cc}^+$	*
$N(2300)$	1/2 <sup>+</sup> **	$\Lambda(1600)$	1/2 <sup>+</sup> ***	$\Sigma(2250)$	***			$\Lambda_b^0$	1/2 <sup>+</sup> ***
$N(2570)$	5/2 <sup>-</sup> **	$\Lambda(1670)$	1/2 <sup>-</sup> ****	$\Sigma(2455)$	**			$\Lambda_b(5912)^0$	1/2 <sup>-</sup> ***
$N(2600)$	11/2 <sup>-</sup> ***	$\Lambda(1690)$	3/2 <sup>-</sup> ****	$\Sigma(2620)$	**			$\Lambda_b(5920)^0$	3/2 <sup>-</sup> ***
$N(2700)$	13/2 <sup>+</sup> **	$\Lambda(1710)$	1/2 <sup>+</sup> *	$\Sigma(3000)$	*			$\Lambda_b(5920)^0$	1/2 <sup>+</sup> ***
		$\Lambda(1800)$	1/2 <sup>-</sup> ***		*			$\Sigma_b$	1/2 <sup>+</sup> ***
		$\Lambda(1810)$	1/2 <sup>+</sup> ***		*			$\Sigma_b$	3/2 <sup>+</sup> ***
		$\Lambda(1820)$	5/2 <sup>+</sup> ****		*			$\Xi_b^0, \Xi_b^-$	1/2 <sup>+</sup> ***
		$\Lambda(1830)$	5/2 <sup>-</sup> ****		*			$\Xi_b(5935)^0$	1/2 <sup>+</sup> ***
		$\Lambda(1890)$	3/2 <sup>+</sup> ****		*			$\Xi_b(5945)^0$	3/2 <sup>+</sup> ***
		$\Lambda(2000)$	*		*			$\Xi_b(5955)^0$	3/2 <sup>+</sup> ***
		$\Lambda(2020)$	7/2 <sup>+</sup> *		*			$\Omega_b$	1/2 <sup>+</sup> ***
		$\Lambda(2050)$	3/2 <sup>-</sup> *		*				
		$\Lambda(2100)$	7/2 <sup>-</sup> ****		*				
		$\Lambda(2110)$	5/2 <sup>+</sup> ***		*				
		$\Lambda(2325)$	3/2 <sup>-</sup> *		*				
		$\Lambda(2350)$	9/2 <sup>+</sup> **		*				
		$\Lambda(2585)$	**		*				



155 baryons

LIGHT UNFLAVORED (S = C = B = 0)		STRANGE (S = ±1, C = B = 0)		CHARMED, STRANGE (C = S = ±1)		CC F <sub>c</sub> (F <sub>c</sub> )	
F(F <sub>c</sub> )		F(F <sub>c</sub> )		F(F <sub>c</sub> )		F(F <sub>c</sub> )	
• $\pi^\pm$	1 <sup>-</sup> (0 <sup>-</sup> )	• $\rho(1680)$	0 <sup>-</sup> (1 <sup>-</sup> )	• $K^\pm$	1/2(0 <sup>-</sup> )	• $D_s^\pm$	0(0 <sup>-</sup> )
• $\pi^0$	1 <sup>-</sup> (0 <sup>-</sup> )	• $\rho(1690)$	1 <sup>+</sup> (3 <sup>-</sup> )	• $K_S^0$	1/2(0 <sup>-</sup> )	• $D_s^\pm$	0(?)
• $\eta$	0 <sup>+</sup> (0 <sup>+</sup> )	• $\rho(1700)$	1 <sup>+</sup> (1 <sup>-</sup> )	• $K_L^0$	1/2(0 <sup>-</sup> )	• $D_{s1}(2317)^\pm$	0(0 <sup>+</sup> )
• $\eta(500)$	0 <sup>+</sup> (0 <sup>+</sup> )	• $a_2(1700)$	1 <sup>-</sup> (2 <sup>+</sup> )	• $K_1^0$	1/2(0 <sup>+</sup> )	• $D_{s1}(2460)^\pm$	0(1 <sup>+</sup> )
• $\rho(770)$	1 <sup>+</sup> (1 <sup>-</sup> )	• $\omega(1710)$	0 <sup>+</sup> (0 <sup>+</sup> )	• $K_0^*(800)$	1/2(0 <sup>+</sup> )	• $D_{s1}(2536)^\pm$	0(1 <sup>+</sup> )
• $\omega(782)$	0 <sup>+</sup> (1 <sup>-</sup> )	• $\eta(1760)$	0 <sup>+</sup> (0 <sup>+</sup> )	• $K^*(892)$	1/2(1 <sup>-</sup> )	• $D_{s2}(2573)$	0(?)
• $\eta(958)$	0 <sup>+</sup> (0 <sup>+</sup> )	• $\pi(1800)$	1 <sup>-</sup> (0 <sup>+</sup> )	• $K_1(1270)$	1/2(1 <sup>+</sup> )	• $D_{s1}^*(2700)^\pm$	0(1 <sup>-</sup> )
• $\eta(980)$	0 <sup>+</sup> (0 <sup>+</sup> )	• $f_2(1810)$	0 <sup>+</sup> (2 <sup>+</sup> )	• $K_1(1400)$	1/2(1 <sup>+</sup> )	• $D_{s1}^*(2860)^\pm$	0(?)
• $a_0(980)$	1 <sup>-</sup> (0 <sup>+</sup> )	• $X(1835)$	?(2 <sup>-</sup> )	• $K^*(1410)$	1/2(1 <sup>+</sup> )	• $D_{s1}(3040)^\pm$	0(?)
• $\phi(1020)$	0 <sup>-</sup> (1 <sup>-</sup> )	• $X(1840)$	?(2 <sup>?</sup> )	• $K_0^*(1430)$	1/2(0 <sup>+</sup> )		
• $h_1(1170)$	0 <sup>-</sup> (1 <sup>+</sup> )	• $\omega_3(1850)$	0 <sup>-</sup> (3 <sup>-</sup> )	• $K_2^*(1430)$	1/2(2 <sup>+</sup> )		
• $b_1(1235)$	1 <sup>+</sup> (1 <sup>+</sup> )	• $\eta_2(1870)$	0 <sup>+</sup> (2 <sup>-</sup> )	• $K(1460)$	1/2(0 <sup>-</sup> )		
• $a_1(1260)$	1 <sup>-</sup> (1 <sup>+</sup> )	• $\pi_2(1880)$	1 <sup>-</sup> (2 <sup>+</sup> )	• $K_2(1580)$	1/2(2 <sup>-</sup> )		
• $f_2(1270)$	0 <sup>+</sup> (2 <sup>+</sup> )	• $\rho(1900)$	1 <sup>+</sup> (1 <sup>-</sup> )	• $K_1(1630)$	1/2(1 <sup>+</sup> )		
• $f_1(1285)$	0 <sup>+</sup> (1 <sup>+</sup> )	• $f_2(1910)$	0 <sup>+</sup> (2 <sup>+</sup> )	• $K_1(1650)$	1/2(1 <sup>+</sup> )		
• $\eta(1295)$	0 <sup>+</sup> (0 <sup>+</sup> )	• $f_2(1950)$	0 <sup>+</sup> (2 <sup>+</sup> )	• $K^*(1680)$	1/2(1 <sup>-</sup> )		
• $\pi(1300)$	1 <sup>-</sup> (0 <sup>+</sup> )	• $\rho_3(1990)$	1 <sup>+</sup> (3 <sup>-</sup> )	• $K_2(1770)$	1/2(2 <sup>-</sup> )		
• $a_2(1320)$	1 <sup>-</sup> (2 <sup>+</sup> )	• $f_2(2010)$	0 <sup>+</sup> (2 <sup>+</sup> )	• $K_3^*(1780)$	1/2(3 <sup>-</sup> )		
• $f_0(1370)$	0 <sup>+</sup> (0 <sup>+</sup> )	• $f_0(2020)$	0 <sup>+</sup> (0 <sup>+</sup> )	• $K_0^*(1820)$	1/2(2 <sup>-</sup> )		
• $h_1(1380)$	?(1 <sup>+</sup> )	• $a_4(2040)$	1 <sup>-</sup> (4 <sup>+</sup> )	• $K_2(1830)$	1/2(0 <sup>-</sup> )		
• $\pi_1(1400)$	1 <sup>-</sup> (1 <sup>+</sup> )	• $f_4(2050)$	0 <sup>+</sup> (4 <sup>+</sup> )	• $K_1^*(1830)$	1/2(0 <sup>+</sup> )		
• $\eta(1405)$	0 <sup>+</sup> (0 <sup>+</sup> )	• $\pi_2(2100)$	1 <sup>-</sup> (2 <sup>+</sup> )	• $K_0^*(1850)$	1/2(0 <sup>+</sup> )		
• $f_1(1420)$	0 <sup>+</sup> (1 <sup>+</sup> )	• $f_0(2100)$	0 <sup>+</sup> (0 <sup>+</sup> )	• $K_1^*(1880)$	1/2(2 <sup>+</sup> )		
• $\omega(1420)$	0 <sup>-</sup> (1 <sup>-</sup> )	• $f_2(2150)$	0 <sup>+</sup> (2 <sup>+</sup> )	• $K_4^*(2045)$	1/2(4 <sup>+</sup> )		
• $f_2(1430)$	0 <sup>+</sup> (2 <sup>+</sup> )	• $\rho(2150)$	1 <sup>+</sup> (1 <sup>-</sup> )	• $K_2(2250)$	1/2(2 <sup>-</sup> )		
• $a_0(1450)$	1 <sup>-</sup> (0 <sup>+</sup> )	• $\phi(2170)$	0 <sup>-</sup> (1 <sup>-</sup> )	• $K_3(2320)$	1/2(3 <sup>+</sup> )		
• $\eta(1450)$	1 <sup>-</sup> (1 <sup>-</sup> )	• $f_0(2200)$	0 <sup>+</sup> (0 <sup>+</sup> )	• $K_3^*(2380)$	1/2(5 <sup>-</sup> )		
• $\eta(1475)$	0 <sup>+</sup> (0 <sup>+</sup> )	• $f_2(2220)$	0 <sup>+</sup> (2 <sup>+</sup> )	• $K_4(2500)$	1/2(4 <sup>-</sup> )		
• $f_0(1500)$	0 <sup>+</sup> (0 <sup>+</sup> )	• $\eta(2225)$	0 <sup>+</sup> (0 <sup>+</sup> )	• $K(3100)$	?(2 <sup>?</sup> )		
• $f_1(1510)$	0 <sup>+</sup> (1 <sup>+</sup> )	• $\rho_3(2250)$	1 <sup>+</sup> (3 <sup>-</sup> )				
• $f_2(1525)$	0 <sup>+</sup> (2 <sup>+</sup> )	• $f_2(2300)$	0 <sup>+</sup> (2 <sup>+</sup> )				
• $f_3(1565)$	0 <sup>+</sup> (2 <sup>+</sup> )	• $f_4(2300)$	0 <sup>+</sup> (4 <sup>+</sup> )				
• $\rho(1570)$	1 <sup>+</sup> (1 <sup>-</sup> )	• $f_0(2330)$	0 <sup>+</sup> (0 <sup>+</sup> )				
• $h_1(1595)$	0 <sup>-</sup> (1 <sup>+</sup> )	• $f_2(2340)$	0 <sup>+</sup> (2 <sup>+</sup> )				
• $\pi_1(1600)$	1 <sup>-</sup> (1 <sup>+</sup> )	• $\rho_5(2350)$	1 <sup>+</sup> (5 <sup>-</sup> )				
• $a_1(1640)$	1 <sup>-</sup> (1 <sup>+</sup> )	• $a_6(2450)$	1 <sup>-</sup> (6 <sup>+</sup> )				
• $f_2(1640)$	0 <sup>+</sup> (2 <sup>+</sup> )		0 <sup>+</sup> (6 <sup>+</sup> )				
• $\Sigma_b(1645)$	0 <sup>+</sup> (2 <sup>+</sup> )						
• $\omega(1650)$	0 <sup>-</sup> (1 <sup>-</sup> )						
• $\omega_3(1670)$	0 <sup>-</sup> (3 <sup>-</sup> )						
• $\pi_2(1670)$	1 <sup>-</sup> (2 <sup>+</sup> )						



206 mesons

All ~ 370 hadrons emerge from single QCD Lagrangian.

# Observed hadrons (2020)

PDG 2020 edition

<http://pdg.lbl.gov/>

Only **color singlet** states are observed.

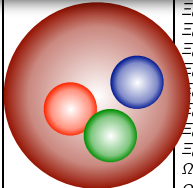
—> Color confinement problem

Flavor quantum numbers are described by  $qqq/q\bar{q}$ .

Why no  $qq\bar{q}\bar{q}$ ,  $qqqq\bar{q}$ , ... states (**exotic hadrons**)?

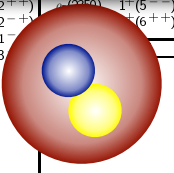
—> Exotic hadron problem, as nontrivial as confinement!

$\Lambda(1820)$	$5/2^+$	****
$\Lambda(1830)$	$5/2^-$	****
$\Lambda(1890)$	$3/2^+$	****
$\Lambda(2000)$	$1/2^-$	*
$\Lambda(2050)$	$3/2^-$	*
$\Lambda(2070)$	$3/2^+$	*
$\Lambda(2080)$	$5/2^-$	*
$\Lambda(2085)$	$7/2^+$	**
$\Lambda(2100)$	$7/2^-$	****
$\Lambda(2110)$	$5/2^+$	***
$\Lambda(2325)$	$3/2^-$	*
$\Lambda(2350)$	$9/2^+$	***
$\Lambda(2585)$		**



**162 baryons**

$D_s(2000)^0$	$1/2(0^-)$
$D_s(2420)^0$	$1/2(1^+)$
$D_s(2430)^0$	$1/2(1^?)$
$D_s(2460)^0$	$1/2(2^+)$
$D_s(2460)^+$	$1/2(2^+)$
$D_s(2550)^0$	$1/2(2^?)$
$D_s(2600)^0$	$1/2(2^?)$
$D_s(2640)^+$	$1/2(2^?)$
$D_s(2750)^0$	$1/2(3^-)$
$D_s(2740)^0$	$1/2(2^?)$
$D_s(3000)^0$	$1/2(2^?)$



**209 mesons**

All ~ 380 hadrons emerge from single QCD Lagrangian.

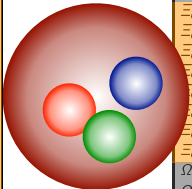


# Unstable states via strong interaction

## Stable/unstable hadrons

<http://pdg.lbl.gov/>

$p$	$1/2^+$ ****	$\Delta(1232)$	$3/2^+$ ****	$\Sigma^+$	$1/2^+$ ****	$\Xi^0$	$1/2^+$ ****	$\Xi^{++}$	***
$n$	$1/2^+$ ****	$\Delta(1600)$	$3/2^+$ ****	$\Sigma^0$	$1/2^+$ ****	$\Xi^-$	$1/2^+$ ****	$\Xi_{cc}^{++}$	***
$N(1440)$	$1/2^+$ ****	$\Delta(1620)$	$1/2^-$ ****	$\Sigma(1385)$	$3/2^+$ ****	$\Xi(1530)$	$3/2^+$ ****	$\Lambda_b^0$	$1/2^+$ ***
$N(1520)$	$3/2^-$ ****	$\Delta(1700)$	$3/2^-$ ****	$\Sigma(1580)$	$3/2^-$ ****	$\Xi(1620)$	*	$\Lambda_b(5912)^0$	$1/2^-$ ***
$N(1535)$	$1/2^-$ ****	$\Delta(1750)$	$1/2^+$ *	$\Sigma(1620)$	$1/2^-$ *	$\Xi(1690)$	***	$\Lambda_b(5920)^0$	$3/2^-$ ***
$N(1650)$	$1/2^-$ ****	$\Delta(1900)$	$1/2^-$ ***	$\Sigma(1660)$	$1/2^+$ ***	$\Xi(1820)$	$3/2^-$ ***	$\Lambda_b(6146)^0$	$3/2^+$ ***
$N(1675)$	$5/2^-$ ****	$\Delta(1905)$	$5/2^+$ ****	$\Sigma(1670)$	$3/2^-$ ****	$\Xi(1950)$	***	$\Lambda_b(6152)^0$	$5/2^+$ ****
$N(1680)$	$5/2^+$ ****	$\Delta(1910)$	$1/2^+$ ****	$\Sigma(1750)$	$1/2^-$ ***	$\Xi(2030)$	$\geq \frac{5}{2}^?$	$\Sigma_b$	$1/2^+$ ***
$N(1700)$	$3/2^-$ ***	$\Delta(1920)$	$3/2^+$ ***	$\Sigma(1775)$	$5/2^-$ ****	$\Xi(2120)$	*	$\Sigma_b^+$	$3/2^+$ ****
$N(1710)$	$1/2^+$ ****	$\Delta(1930)$	$5/2^-$ **	$\Sigma(1880)$	$1/2^+$ **	$\Xi(2250)$	**	$\Sigma_b(6097)^+$	***
$N(1720)$	$3/2^+$ ****	$\Delta(1940)$	$3/2^-$ **	$\Sigma(1900)$	$1/2^-$ **	$\Xi(2370)$	**	$\Sigma_b(6097)^-$	***
$N(1860)$	$5/2^+$ **	$\Delta(1950)$	$7/2^+$ ****	$\Sigma(2030)$	$7/2^+$ ****	$\Xi(2500)$	*	$\Xi_b^0, \Xi_b^-$	$1/2^+$ ****
$N(1875)$	$3/2^-$ ***	$\Delta(2000)$	$5/2^+$ **	$\Sigma(2070)$	$5/2^+$ **	$\Xi(2935)$	$1/2^+$ ***	$\Xi_b(5935)$	$1/2^+$ ****
$N(1880)$	$1/2^+$ ***	$\Delta(2150)$	$1/2^-$ *	$\Sigma(2250)$	***	$\Omega(2012)$	$3/2^+$ ****	$\Xi_b(5945)^0$	$3/2^+$ ****
$N(1895)$	$1/2^-$ ****	$\Delta(2200)$	$7/2^-$ ***	$\Sigma(2330)$	$3/2^+$ **	$\Omega(2250)$	***	$\Xi_b(5955)$	$3/2^+$ ****
$N(1890)$	$3/2^+$ ****	$\Delta(2300)$	$9/2^+$ **	$\Sigma(2400)$	$9/2^-$ **	$\Omega(2380)$	**	$\Xi_b(6227)$	***
$N(1990)$	$7/2^+$ **	$\Delta(2350)$	$5/2^-$ *	$\Sigma(2470)$	$7/2^+$ **	$\Omega(2470)$	**	$\Omega_b$	$1/2^+$ ****
$N(2000)$	$5/2^+$ **	$\Delta(2390)$	$7/2^+$ **	$\Lambda$	$1/2^+$ ****	$\Lambda_c(2595)^+$	$1/2^-$ ****	$P_c(4312)^+$	*
$N(2040)$	$3/2^+$ *	$\Delta(2400)$	$9/2^-$ **	$\Lambda(1405)$	$1/2^-$ ***	$\Lambda_c(2625)^+$	$3/2^-$ ***	$P_c(4380)^+$	*
$N(2060)$	$5/2^-$ ***	$\Delta(2420)$	$11/2^+$ ****	$\Lambda(1520)$	$3/2^-$ ****	$\Lambda_c(2765)^+$	*	$P_c(4440)^+$	*
$N(2100)$	$1/2^+$ ****	$\Delta(2750)$	$13/2^-$ **	$\Lambda(1670)$	$1/2^-$ ****	$\Lambda_c(2860)^+$	$3/2^+$ ****	$P_c(4457)^+$	*
$N(2120)$	$3/2^-$ ***	$\Delta(2950)$	$15/2^+$ **	$\Lambda(1810)$	$1/2^+$ ***	$\Lambda_c(2880)^+$	$5/2^+$ ****		
$N(2190)$	$7/2^-$ ****			$\Lambda(1820)$	$5/2^+$ ****	$\Lambda_c(2940)^+$	$3/2^-$ ****		
$N(2200)$	$9/2^+$ ****	$\Lambda$	$1/2^+$ ****	$\Lambda(1830)$	$5/2^-$ ****	$\Lambda_c(2980)^+$	***		
$N(2250)$	$9/2^-$ ****	$\Lambda$	$1/2^-$ ***	$\Lambda(1890)$	$3/2^+$ ****	$\Lambda_c(3080)$	***		
$N(2300)$	$1/2^+$ **	$\Lambda(1405)$	$1/2^-$ ***	$\Lambda(1890)$	$3/2^+$ ****	$\Lambda_c(3123)$	*		
$N(2570)$	$5/2^-$ **	$\Lambda(1520)$	$3/2^-$ ****	$\Lambda(2000)$	$1/2^-$ *	$\Omega_c^0$	$1/2^+$ ****		
$N(2600)$	$11/2^-$ ***	$\Lambda(1600)$	$1/2^+$ ****	$\Lambda(2050)$	$3/2^-$ *	$\Omega_c(2790)$	$1/2^-$ ***		
$N(2700)$	$13/2^+$ **	$\Lambda(1670)$	$1/2^-$ ****	$\Lambda(2070)$	$3/2^+$ *	$\Omega_c(2815)$	$3/2^-$ ****		
		$\Lambda(1690)$	$3/2^-$ ****	$\Lambda(2080)$	$5/2^-$ *	$\Omega_c(2890)$	**		
		$\Lambda(1710)$	$1/2^+$ *	$\Lambda(2085)$	$7/2^+$ **	$\Omega_c(2930)$	**		
		$\Lambda(1800)$	$1/2^-$ ***	$\Lambda(2100)$	$7/2^-$ ****	$\Omega_c(2970)$	***		
		$\Lambda(1810)$	$1/2^+$ ***	$\Lambda(2110)$	$5/2^+$ ****	$\Omega_c(3055)$	***		
		$\Lambda(1820)$	$5/2^+$ ****	$\Lambda(2325)$	$3/2^-$ *	$\Omega_c(3080)$	***		
		$\Lambda(1830)$	$5/2^-$ ****	$\Lambda(2350)$	$9/2^+$ **	$\Omega_c(3123)$	*		
		$\Lambda(1890)$	$3/2^+$ ****	$\Lambda(2585)$	**				
		$\Lambda(2000)$	$1/2^-$ *						
		$\Lambda(2050)$	$3/2^-$ *						
		$\Lambda(2070)$	$3/2^+$ *						
		$\Lambda(2080)$	$5/2^-$ *						
		$\Lambda(2085)$	$7/2^+$ **						
		$\Lambda(2100)$	$7/2^-$ ****						
		$\Lambda(2110)$	$5/2^+$ ****						
		$\Lambda(2325)$	$3/2^-$ *						
		$\Lambda(2350)$	$9/2^+$ **						
		$\Lambda(2585)$	**						



162 baryons

LIGHT UNFLAVORED ( $S=C=B=0$ )		STRANGE ( $S=\pm 1, C=B=0$ )		CHARMED, STRANGE ( $C=S=\pm 1$ )		$c\bar{c}$ continued $\bar{c}c$	
$f(\bar{f})$		$f(\bar{f})$		$f(\bar{f})$		$f(\bar{f})$	
$\pi^+$	$1(0^-)$	$\pi_2(1670)$	$1(2^-)$	$K^+$	$1/2(0^-)$	$D_s^+$	$0(1^-)$
$\pi^0$	$1(0^-)$	$\rho(1690)$	$0(1^-)$	$K^0$	$1/2(0^-)$	$D_s^0$	$0(2^-)$
$\eta$	$0(0^-)$	$\rho_2(1690)$	$1^+(3^-)$	$K_S^0$	$1/2(0^-)$	$D_{s1}^+(2317)$	$0(0^+)$
$\eta(500)$	$0^+(0^+)$	$\rho(1700)$	$1^+(1^-)$	$K_L^0$	$1/2(0^-)$	$D_{s1}(2460)^+$	$0(1^+)$
$\eta(770)$	$1^+(1^-)$	$\omega(1700)$	$1(2^+)$	$K_S^*(700)$	$1/2(0^+)$	$D_{s1}(2536)^+$	$0(1^+)$
$\eta(782)$	$0(1^-)$	$\phi(1700)$	$0^+(0^+)$	$K^*(892)$	$1/2(1^-)$	$D_{s2}^+(2573)$	$0(2^+)$
$\eta(1295)$	$0^+(0^+)$	$\phi(1710)$	$0^+(0^+)$	$K_1(1270)$	$1/2(1^+)$	$D_{s1}^*(2700)^+$	$0(1^-)$
$\eta(980)$	$0^+(0^+)$	$\phi(1760)$	$0^+(0^+)$	$K_1(1400)$	$1/2(1^+)$	$D_{s1}^*(2860)^+$	$0(1^-)$
$\omega(980)$	$0^+(0^+)$	$\pi(1800)$	$1(0^-)$	$K^*(1410)$	$1/2(1^-)$	$D_{s3}^*(2860)^+$	$0(3^-)$
$\omega(1020)$	$0(1^-)$	$\phi(1835)$	$?^?(0^-)$	$K_S^*(1430)$	$1/2(0^+)$	$D_{s1}(3040)^+$	$0(2^?)$
$h_1(1170)$	$0^+(1^+)$	$\phi_2(1850)$	$0(3^-)$	$K_S^*(1430)$	$1/2(2^+)$	BOTTOM ( $B=S=1$ )	
$h_1(1235)$	$1^+(1^+)$	$\eta(1870)$	$0^+(2^-)$	$K(1460)$	$1/2(0^-)$	$B^+$	$1/2(0^-)$
$a_1(1260)$	$1^+(1^+)$	$\pi_2(1880)$	$1(2^-)$	$K(1580)$	$1/2(2^-)$	$B^0$	$1/2(0^-)$
$f_1(1285)$	$0^+(1^+)$	$\pi(1900)$	$1^+(1^-)$	$K(1630)$	$1/2(2^+)$	$B^+/\bar{B}^0$ ADMIXTURE	
$\eta(1295)$	$0^+(0^+)$	$f_2(1910)$	$0^+(2^+)$	$K_1(1650)$	$1/2(1^+)$	$B^+/\bar{B}^0/B_s^0$ b-baryon ADMIXTURE	
$\pi(1300)$	$1(0^-)$	$a_1(1950)$	$1(0^+)$	$K_1(1650)$	$1/2(2^+)$	$V_{cb}$ and $V_{cb}$ CKM Matrix Elements	
$a_2(1320)$	$1(2^+)$	$f_2(1950)$	$0^+(2^+)$	$K_2(1770)$	$1/2(2^-)$	$B^+$	$1/2(2^+)$
$K_1(1370)$	$0^+(1^+)$	$a_1(1970)$	$1(4^+)$	$K_2(1780)$	$1/2(3^-)$	$B^0$	$1/2(2^+)$
$\pi_1(1400)$	$1(1^+)$	$\pi_2(1990)$	$1^+(3^-)$	$K_2^*(1820)$	$1/2(2^+)$	$B^+$	$1/2(2^+)$
$\eta(1405)$	$0^+(0^+)$	$\pi_2(2005)$	$1(2^-)$	$K(1830)$	$1/2(0^-)$	$B^+$	$1/2(1^-)$
$h_1(1415)$	$0^+(1^+)$	$f_2(2010)$	$0^+(2^+)$	$K_2(1950)$	$1/2(0^+)$	$B_1(5721)^+$	$1/2(1^+)$
$a_1(1420)$	$1(1^+)$	$f_2(2020)$	$0^+(0^+)$	$K_2(1980)$	$1/2(2^+)$	$B_1(5721)^0$	$1/2(1^+)$
$f_1(1420)$	$0^+(1^+)$	$f_4(2060)$	$0^+(4^+)$	$K_2^*(2045)$	$1/2(4^+)$	$B_1^*(5732)$	$?^?(2^?)$
$\omega(1420)$	$0(1^-)$	$\pi_2(2100)$	$1(2^-)$	$K_2^*(2045)$	$1/2(4^+)$	$B_1^*(5747)^+$	$1/2(2^+)$
$f_2(1430)$	$0^+(2^+)$	$\pi_2(2150)$	$1(2^-)$	$K_2(2250)$	$1/2(2^-)$	$B_2^*(5747)^0$	$1/2(2^+)$
$a_2(1450)$	$1(0^+)$	$\mu(2170)$	$1(1^-)$	$K_2^*(2300)$	$1/2(5^-)$	$B_1(5840)^+$	$1/2(2^?)$
$\rho(1450)$	$1^+(1^+)$	$\phi(2180)$	$0(1^-)$	$K_1(2500)$	$1/2(4^-)$	$B_1(5840)^0$	$1/2(2^?)$
$\eta(1475)$	$0^+(0^+)$	$f_2(2200)$	$0^+(0^+)$	$K(3100)$	$?^?(2^?)$	$B_1(5970)^+$	$1/2(2^?)$
$f_1(1500)$	$0^+(0^+)$	$\eta(2220)$	$0^+(2^+)$	CHARMED ( $C=\pm 1$ )		$B_1(5970)^0$	$1/2(2^?)$
$f_1(1510)$	$0^+(1^+)$	$\rho(2250)$	$1^+(3^-)$	$D^+$	$1/2(0^-)$	BOTTOM, STRANGE ( $B=S=\pm 1, S=\pm 1$ )	
$f_1(1525)$	$0^+(2^+)$	$\rho_2(2250)$	$1^+(3^-)$	$D^0$	$1/2(0^-)$	$B_s^+$	$0(0^-)$
$f_2(1565)$	$0^+(2^+)$	$\eta(2300)$	$0^+(2^+)$	$D^*$	$1/2(1^-)$	$B_s^0$	$0(1^-)$
$\rho(1570)$	$1^+(1^-)$	$f_4(2300)$	$0^+(4^+)$	$D^*(2007)^0$	$1/2(1^-)$	$B_s^*$	$?^?(2^?)$
$h(1595)$	$0^+(1^+)$	$f_2(2300)$	$0^+(0^+)$	$D^*(2010)^+$	$1/2(1^-)$	$\chi(5568)^+$	$?^?(2^?)$
$\pi_1(1600)$	$1^+(1^+)$	$f_2(2300)$	$0^+(0^+)$	$D_1^*(2300)^0$	$1/2(0^+)$	$B_{s1}(5830)^0$	$0(1^+)$
$a_1(1640)$	$1^+(1^+)$	$f_2(2340)$	$0^+(2^+)$	$D_1^*(2300)^+$	$1/2(0^+)$	$B_{s1}^*(5840)^0$	$0(2^+)$
$f_2(1640)$	$0^+(2^+)$	$\rho(2350)$	$1^+(5^-)$	$D_1(2420)^0$	$1/2(1^+)$	$B_{s1}^*(5880)$	$?^?(2^?)$
$\eta_2(1645)$	$0^+(2^+)$	$\omega(2350)$	$0(6^+)$	$D_1(2420)^+$	$1/2(2^?)$	BOTTOM, CHARMED ( $B=C=\pm 1$ )	
$\omega(1650)$	$0(1^-)$			$D_1(2430)^0$	$1/2(1^+)$	$B_c^+$	$0(0^-)$
$\omega_2(1670)$	$0(3^-)$			$D_1(2430)^+$	$1/2(1^+)$	$B_c(25)^+$	$0(0^-)$
				$D_2(2460)^0$	$1/2(2^+)$	BOTTOM, CHARMED ( $B=C=\pm 1$ )	
				$D_2(2460)^+$	$1/2(2^+)$	$B_c^+$	$0(0^-)$
				$D_2(2550)^0$	$1/2(2^?)$	$B_c(25)^+$	$0(0^-)$
				$D_2(2550)^+$	$1/2(2^?)$	$B_c(25)^+$	$0(0^-)$
				$D_2(2600)$	$1/2(2^?)$	BOTTOM, CHARMED ( $B=C=\pm 1$ )	
				$D_2(2640)^+$	$1/2(2^?)$	$B_c^+$	$0(0^-)$

# Nature of resonances

## Theoretical treatment for **unstable** hadrons

- **resonances** in hadron-hadron scattering
- **pole** of the scattering amplitude  $\longleftrightarrow$  “eigenstate”

T. Hyodo, M. Niiyama, arXiv: 2010.07592 [hep-ph]

- analytic continuation: unique

## Resonance as an “eigenstate” of Hamiltonian

- **complex energy**

G. Gamow, *Z. Phys.* **51**, 204 (1928)

Zur Quantentheorie des Atomkernes.

Von G. Gamow, z. Zt. in Göttingen.

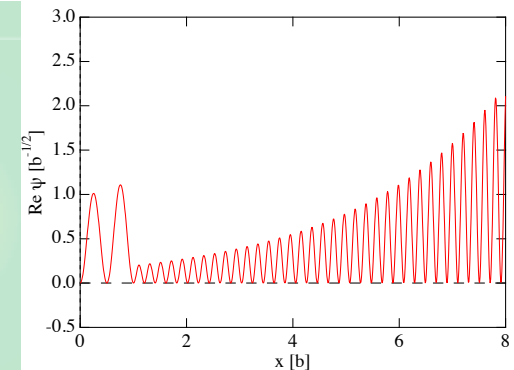
Mit 5 Abbildungen. (Eingegangen am 2. August 1928.)

Um diese Schwierigkeit zu überwinden, müssen wir annehmen, daß die Schwingungen gedämpft sind, und  $E$  komplex setzen:

$$E = E_0 + i \frac{\hbar \lambda}{4 \pi},$$

wo  $E_0$  die gewöhnliche Energie ist und  $\lambda$  das Dämpfungsdekrement (Zerfallskonstante). Dann sehen wir aber aus den Relationen (2 a) und (2 b),

- diverging wave function
- complex expectation value (norm,  $\langle r^2 \rangle$ )
- interpretation problem





## Introduction

- Structure of “unstable” resonance?



## Structure of $\Lambda(1405)$ resonance

- Pole positions

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881, 98 (2012);

P.A. Zyla, *et al.* (Particle Data Group), PTEP 2020, 083C01 (2020)

T. Hyodo, M. Niiyama, arXiv: 2010.07592 [hep-ph]

- $\bar{K}N$  compositeness

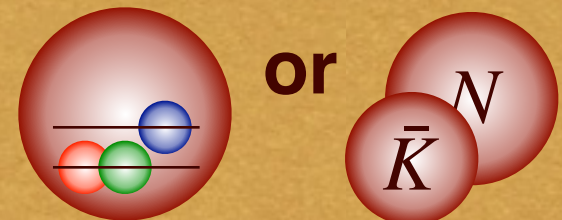
Y. Kamiya, T. Hyodo, PRC93, 035203 (2016); PTEP2017, 023D02 (2017)

T. Hyodo, JPS journal Vol. 75 No. 8, 478 (2020)

T. Kinugawa, T. Hyodo, in preparation



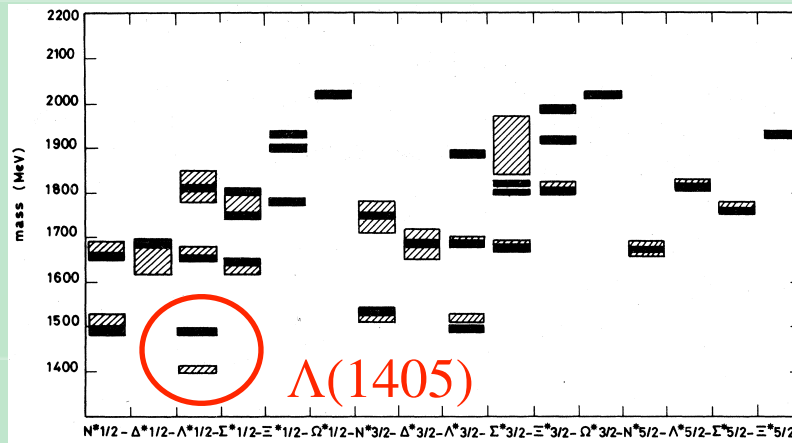
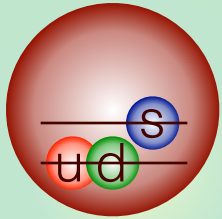
## Summary



# $\Lambda(1405)$ and $\bar{K}N$ scattering

$\Lambda(1405)$  does not fit in standard picture  $\rightarrow$  exotic candidate

N. Isgur and G. Karl, Phys. Rev. D18, 4187 (1978)

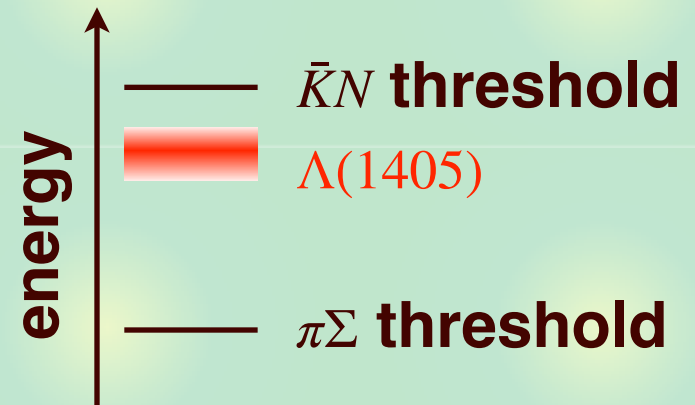
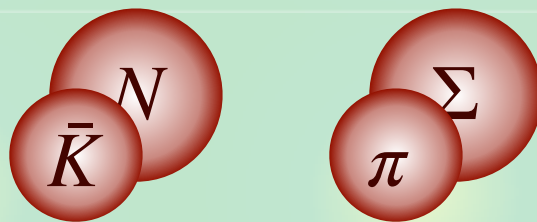


— : theory

▨ : experiment

## Resonance in coupled-channel scattering

- coupling to MB states



Detailed analysis of  $\bar{K}N$ - $\pi\Sigma$  scattering is necessary.



## Strategy for $\bar{K}N$ interaction

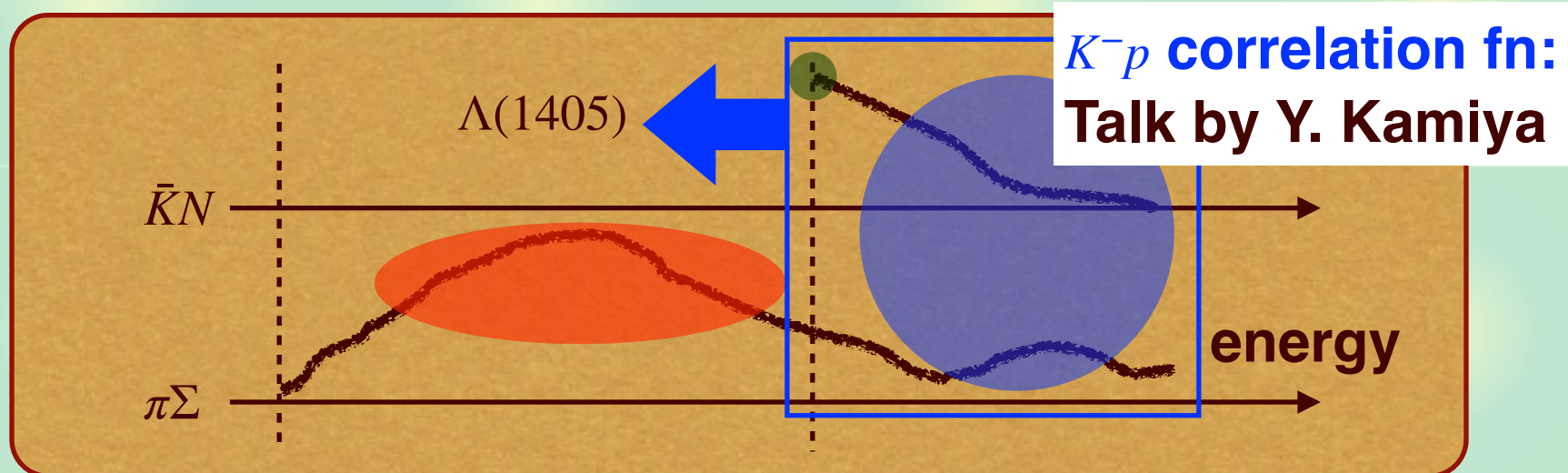
Above the  $\bar{K}N$  threshold : direct constraints

- $K^-p$  total cross sections (old data)
- $\bar{K}N$  threshold branching ratios (old data)
- $K^-p$  scattering length (new data : SIDDHARTA)

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881 98 (2012)

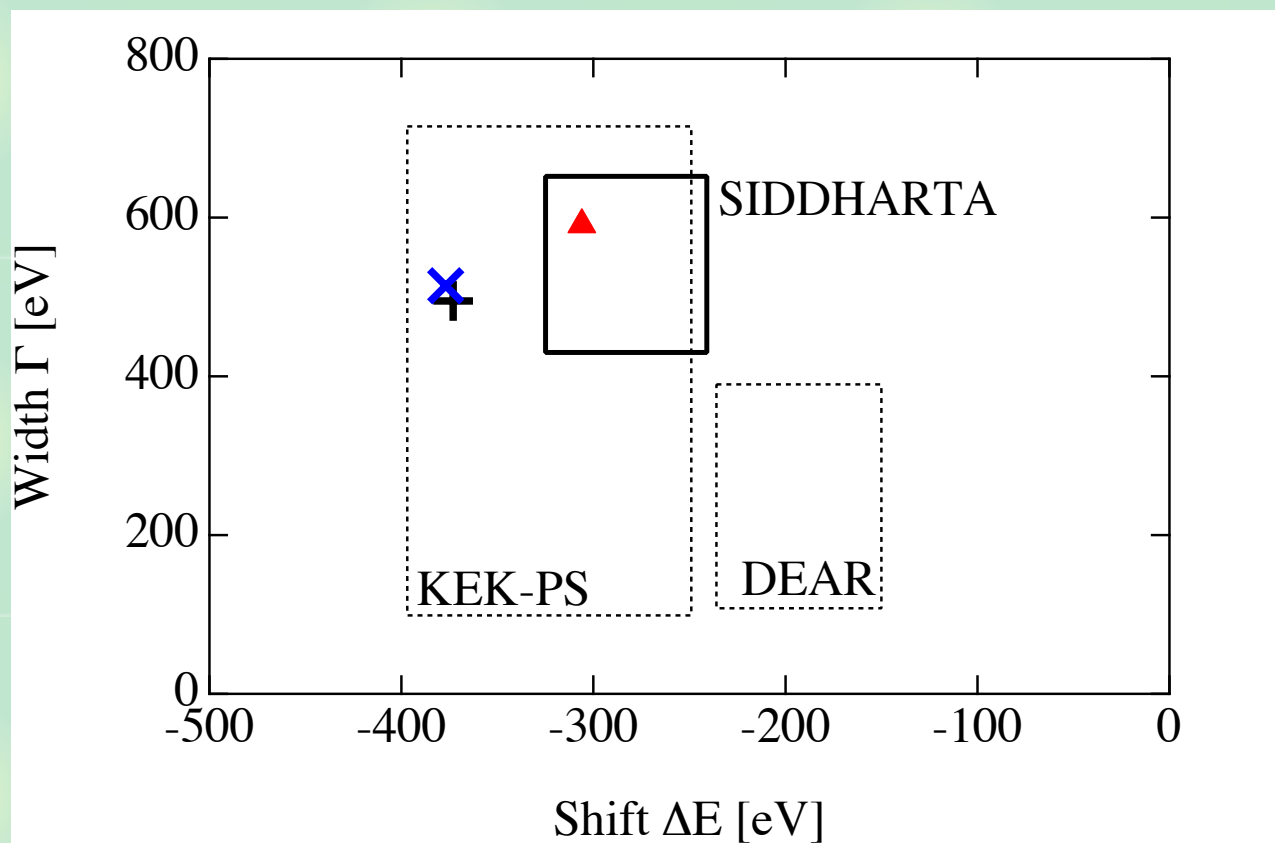
Below the  $\bar{K}N$  threshold: indirect constraints

- $\pi\Sigma$  mass spectra (new data : LEPS, CLAS, HADES, ...)



# Comparison with SIDDHARTA

	<b>TW</b>	<b>TWB</b>	<b>NLO</b>
$\chi^2/\text{d.o.f.}$	<b>1.12</b>	<b>1.15</b>	<b>0.957</b>



**TW** and **TWB** are reasonable, while best-fit requires **NLO**.

# Extrapolation to complex energy: two poles

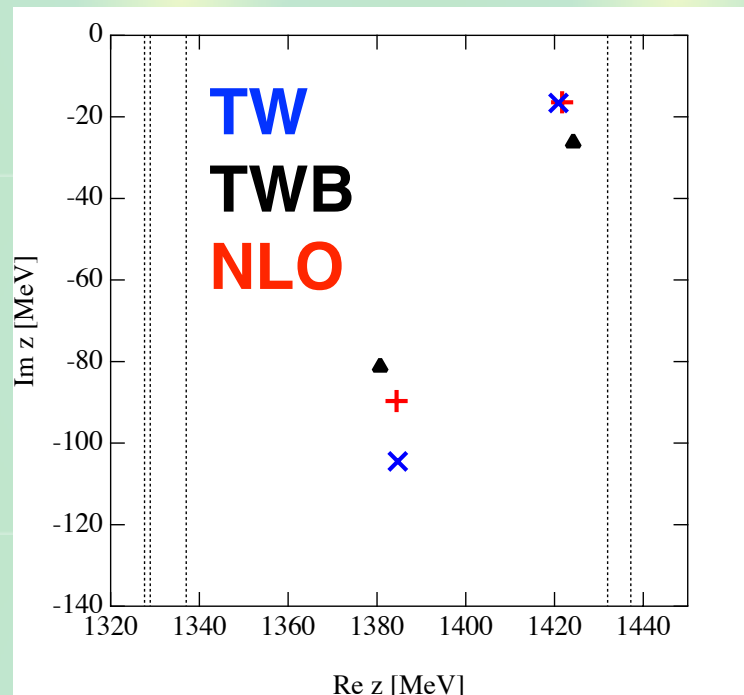
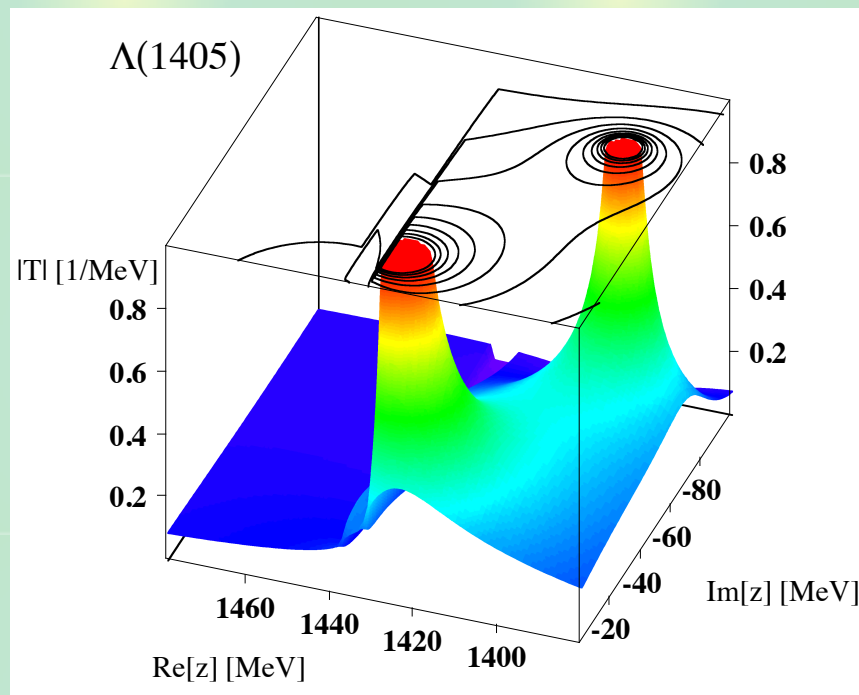
## Two poles : superposition of two eigenstates

J.A. Oller, U.G. Meißner, PLB 500, 263 (2001);

D. Jido, J.A. Oller, E. Oset, A. Ramos, U.G. Meißner, NPA 723, 205 (2003);

U.G. Meißner, Symmetry 12, 981 (2020); M. Mai, arXiv: 2010.00056 [nucl-th];

T. Hyodo, M. Niyama, arXiv: 2010.07592 [hep-ph]



T. Hyodo, D. Jido, Prog. Part. Nucl. Phys. 67, 55 (2012)

**NLO analysis confirms the two-pole structure.**

**PDG has changed**

## 2020 update of PDG

P.A. Zyla, et al., PTEP 2020, 083C01 (2020); <http://pdg.lbl.gov/>

### - Particle Listing section:

Citation: P.A. Zyla et al. (Particle Data Group), Prog. Theor. Exp. Phys. **2020**, 083C01 (2020)

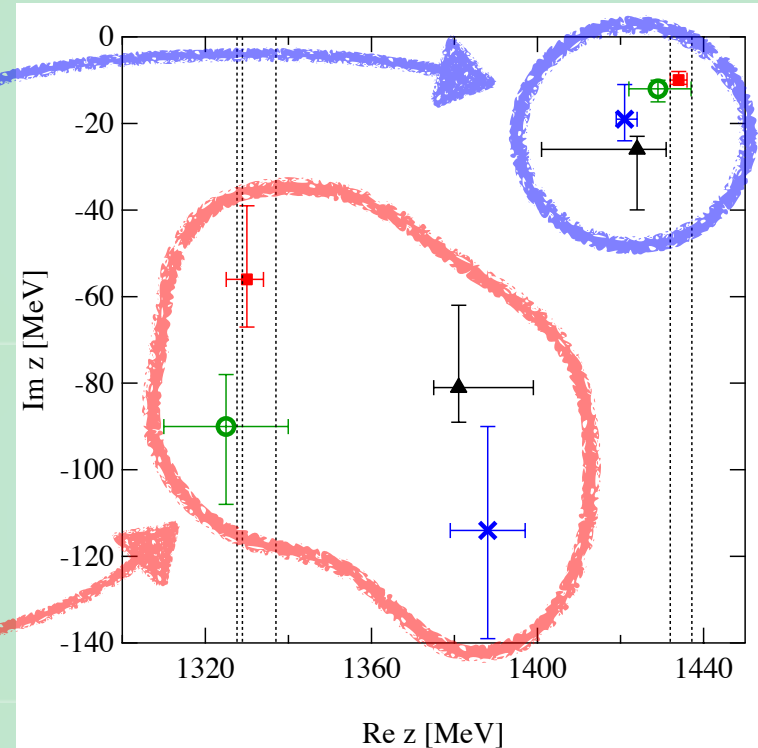
$\Lambda(1405) \ 1/2^-$

$I(J^P) = 0(\frac{1}{2}^-)$  Status: \* \* \* \*

Citation: P.A. Zyla et al. (Particle Data Group), Prog. Theor. Exp. Phys. **2020**, 083C01 (2020)

$\Lambda(1380) \ 1/2^-$

**new!**  $J^P = \frac{1}{2}^-$  Status: \* \*



T. Hyodo, M. Niiyama, arXiv: 2010.07592 [hep-ph]

- Lower pole: two-star resonance  $\Lambda(1380)$
- $\Lambda(1405)$  is no longer at 1405 MeV but  $\sim 1420$  MeV





## Introduction

- Structure of “unstable” resonance?



## Structure of $\Lambda(1405)$ resonance

- Pole positions

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881, 98 (2012);

P.A. Zyla, *et al.* (Particle Data Group), PTEP 2020, 083C01 (2020)

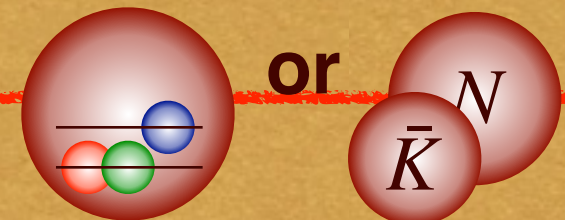
T. Hyodo, M. Niiyama, arXiv: 2010.07592 [hep-ph]

- $\bar{K}N$  compositeness

Y. Kamiya, T. Hyodo, PRC93, 035203 (2016); PTEP2017, 023D02 (2017)

T. Hyodo, JPS journal Vol. 75 No. 8, 478 (2020)

T. Kinugawa, T. Hyodo, in preparation



## Summary

# Weak-binding relation for stable states

Compositeness  $X$  of s-wave **weakly bound** state ( $R \gg R_{\text{typ}}$ )

S. Weinberg, *Phys. Rev.* **137**, B672 (1965);

T. Hyodo, *Int. J. Mod. Phys. A* **28**, 1330045 (2013)

$$|d\rangle = \sqrt{X} |NN\rangle + \sqrt{1-X} |\text{others}\rangle$$

$NN$

**continuum**



**deuteron**

range of interaction

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\}, \quad R = \frac{1}{\sqrt{2\mu B}}$$

↑
↓

↑
↑

scattering length
radius of state

- Deuteron is  $NN$  composite :  $a_0 \sim R \Rightarrow X \sim 1$
- Internal structure from **observable** ( $a_0, B$ )

**Problem: applicable only for stable states**

# Effective field theory

## Low-energy scattering with near-threshold bound state

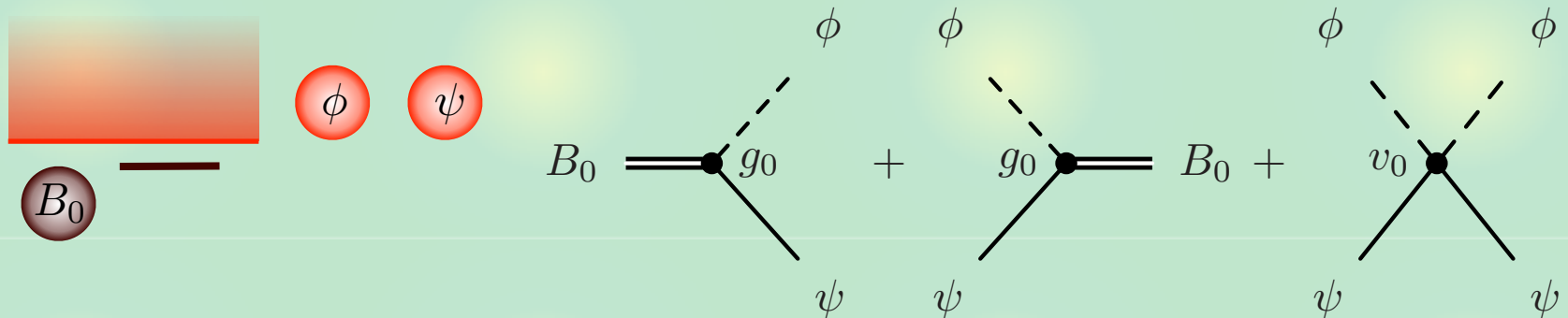
### - Nonrelativistic EFT with contact interaction

D.B. Kaplan, Nucl. Phys. B494, 471 (1997)

E. Braaten, M. Kusunoki, D. Zhang, Annals Phys. 323, 1770 (2008)

$$H_{\text{free}} = \int d\mathbf{r} \left[ \frac{1}{2M} \nabla \psi^\dagger \cdot \nabla \psi + \frac{1}{2m} \nabla \phi^\dagger \cdot \nabla \phi + \frac{1}{2M_0} \nabla B_0^\dagger \cdot \nabla B_0 + \omega_0 B_0^\dagger B_0 \right]$$

$$H_{\text{int}} = \int d\mathbf{r} \left[ g_0 \left( B_0^\dagger \phi \psi + \psi^\dagger \phi^\dagger B_0 \right) + v_0 \psi^\dagger \phi^\dagger \phi \psi \right]$$



- **cutoff** :  $\Lambda \sim 1/R_{\text{typ}}$  (interaction range of microscopic theory)

- At low momentum  $p \ll \Lambda$ , interaction  $\sim$  contact

# Compositeness and “elementariness”

## Eigenstates

$$H_{\text{free}} |B_0\rangle = \omega_0 |B_0\rangle, \quad H_{\text{free}} |\mathbf{p}\rangle = \frac{\mathbf{p}^2}{2\mu} |\mathbf{p}\rangle \quad \text{free (discrete + continuum)}$$

$$(H_{\text{free}} + H_{\text{int}}) |B\rangle = -B |B\rangle \quad \text{full (bound state)}$$

### - normalization of $|B\rangle$ + completeness relation

$$\langle B|B\rangle = 1, \quad 1 = |B_0\rangle\langle B_0| + \int \frac{d\mathbf{p}}{(2\pi)^3} |\mathbf{p}\rangle\langle \mathbf{p}|$$

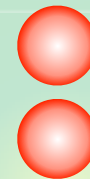
### - projections onto free eigenstates

$$1 = Z + X, \quad Z \equiv |\langle B_0|B\rangle|^2, \quad X \equiv \int \frac{d\mathbf{p}}{(2\pi)^3} |\langle \mathbf{p}|B\rangle|^2$$

“elementarity”



compositeness



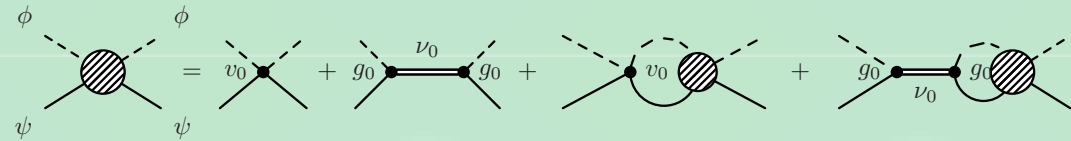
$Z, X$ : real and nonnegative  $\rightarrow$  interpreted as **probability**



# Weak binding relation

## $\psi\phi$ scattering amplitude (exact result)

$$f(E) = -\frac{\mu}{2\pi} \frac{1}{[v(E)]^{-1} - G(E)}$$



$$v(E) = v_0 + \frac{g_0^2}{E - \omega_0}, \quad G(E) = \frac{1}{2\pi^2} \int_0^\Lambda dp \frac{p^2}{E - p^2/(2\mu) + i0^+}$$

## Compositeness $X \leftarrow v(E), G(E)$

$$X = \frac{G'(-B)}{G'(-B) - [1/v(-B)]'}$$

$1/R = \sqrt{2\mu B}$  expansion of scattering length  $a_0$

$$a_0 = -f(E=0) = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\} \text{renormalization dependent}$$

renormalization independent

If  $R \gg R_{\text{typ}}$ , correction terms neglected:  $X \leftarrow (a_0, B)$

# Introduction of decay channel

## Introduce decay channel

$$H'_{\text{free}} = \int d\mathbf{r} \left[ \frac{1}{2M'} \nabla \psi'^{\dagger} \cdot \nabla \psi' - \nu_{\psi} \psi'^{\dagger} \psi' + \frac{1}{2m'} \nabla \phi'^{\dagger} \cdot \nabla \phi' - \nu_{\phi} \phi'^{\dagger} \phi' \right]$$

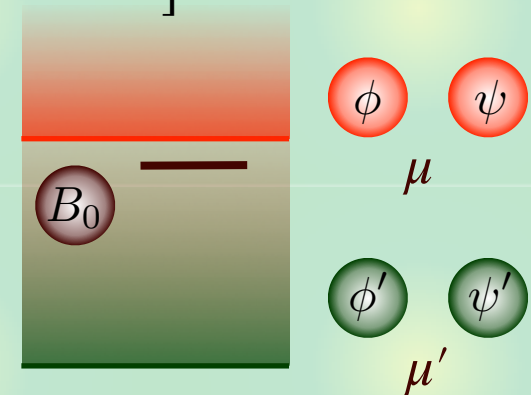
$$H'_{\text{int}} = \int d\mathbf{r} \left[ g'_0 \left( B_0^{\dagger} \phi' \psi' + \psi'^{\dagger} \phi'^{\dagger} B_0 \right) + \nu'_0 \psi'^{\dagger} \phi'^{\dagger} \phi' \psi' + \nu_0^t (\psi'^{\dagger} \phi'^{\dagger} \phi' \psi' + \psi'^{\dagger} \phi'^{\dagger} \phi \psi) \right]$$

## Quasi-bound state : complex eigenvalue

$$H = H_{\text{free}} + H'_{\text{free}} + H_{\text{int}} + H'_{\text{int}}$$

$$H |QB\rangle = E_{QB} |QB\rangle, \quad E_{QB} \in \mathbb{C}$$

$$\nu_{\psi} + \nu_{\phi} = \nu$$



## Generalized relation : **correction** from threshold difference

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O} \left( \left| \frac{R_{\text{typ}}}{R} \right| \right) + \mathcal{O} \left( \left| \frac{\ell}{R} \right|^3 \right) \right\}, \quad R = \frac{1}{\sqrt{-2\mu E_{QB}}}, \quad \ell \equiv \frac{1}{\sqrt{2\mu\nu}}$$

Y. Kamiya, T. Hyodo, PRC93, 035203 (2016); PTEP2017, 023D02 (2017)

c.f. V. Baru, *et al.*, Phys. Lett. B586, 53 (2004), ...

**If**  $|R| \gg (R_{\text{typ}}, \ell)$ , **correction terms neglected:**  $X \leftarrow (a_0, E_{QB})$

# Evaluation of compositeness

## Generalized weak-binding relation

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O} \left( \left| \frac{R_{\text{typ}}}{R} \right| \right) + \mathcal{O} \left( \left| \frac{\ell}{R} \right|^3 \right) \right\}, \quad R = \frac{1}{\sqrt{-2\mu E_{QB}}}, \quad \ell \equiv \frac{1}{\sqrt{2\mu\nu}}$$

$(a_0, E_{QB})$  determinations by several groups

- neglecting correction terms:

	$E_h$ [MeV]	$a_0$ [fm]	$X_{\bar{K}N}$	$\tilde{X}_{\bar{K}N}$	$U/2$
Set 1 [35]	$-10 - i26$	$1.39 - i0.85$	$1.2 + i0.1$	1.0	0.3
Set 2 [36]	$-4 - i8$	$1.81 - i0.92$	$0.6 + i0.1$	0.6	0.0
Set 3 [37]	$-13 - i20$	$1.30 - i0.85$	$0.9 - i0.2$	0.9	0.1
Set 4 [38]	$2 - i10$	$1.21 - i1.47$	$0.6 + i0.0$	0.6	0.0
Set 5 [38]	$-3 - i12$	$1.52 - i1.85$	$1.0 + i0.5$	0.8	0.3

- In all cases,  $X \sim 1$  with small  $U/2$  (complex nature)

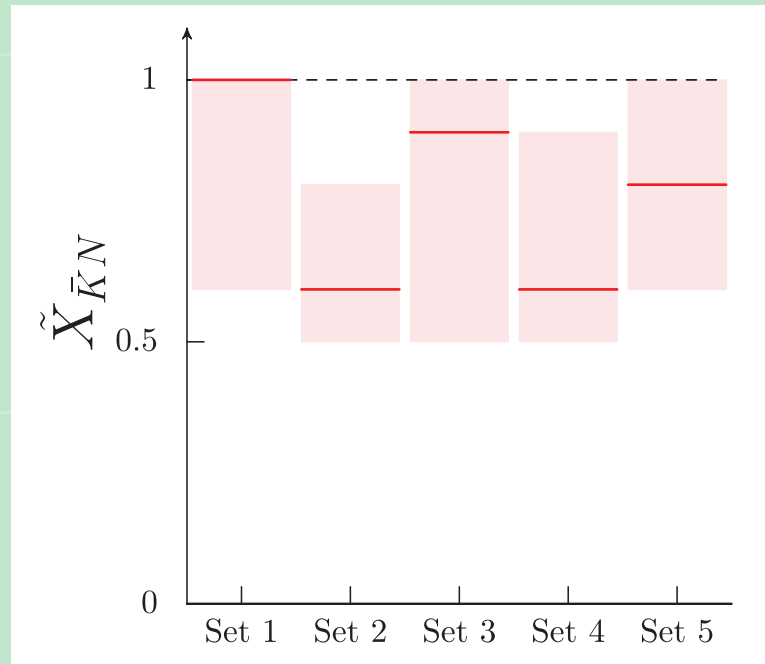
$\Lambda(1405)$ :  $\bar{K}N$  composite dominance  $\leftarrow$  observables

# Uncertainty estimation

Estimation of correction terms:  $|R| \sim 2$  fm

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\left|\frac{R_{\text{typ}}}{R}\right|\right) + \mathcal{O}\left(\left|\frac{\ell}{R}\right|^3\right) \right\}, \quad R = \frac{1}{\sqrt{-2\mu E_{QB}}}, \quad \ell \equiv \frac{1}{\sqrt{2\mu\nu}}$$

- $\rho$  meson exchange picture:  $R_{\text{typ}} \sim 0.25$  fm
- energy difference from  $\pi\Sigma$ :  $\ell \sim 1.08$  fm



$\bar{K}N$  composite dominance holds even **with correction terms.**



# Correction term and zero range limit

What happens if  $R_{\text{typ}} \rightarrow 0$  ?

$$a_0 = R \frac{2X}{1+X}$$

- Limit  $\Lambda \rightarrow \infty$  can be taken in renormalizable EFT
- EFT with only  $\psi, \phi$  fields should have  $X = 1$

$$\Rightarrow a_0 = R$$

- “effective range model” gives  $a_0 \neq R$  : contradiction?

**E. Braaten, M. Kusunoki, D. Zhang, *Annals Phys.* 323, 1770 (2008)**

$R_{\text{typ}}$  should be either  $R_{\text{int}}$  or length scale in the amplitude

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\}, \quad R_{\text{typ}} = \max(R_{\text{int}}, |r_e|, \dots)$$

- relevant to system with large  $|r_e|$

**T. Kinugawa, T. Hyodo, in preparation**

# Summary

- Structure of unstable resonance is **nontrivial**.
- Pole structure of the  $\Lambda(1405)$  region is now well constrained by the experimental data.  
“ $\Lambda(1405)$ ”  $\rightarrow$   $\Lambda(1405)$  **and**  $\Lambda(1380)$
- Generalized weak-binding relation shows that (higher-energy)  $\Lambda(1405)$  is dominated by **molecular  $\bar{K}N$**  component.

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881, 98 (2012);

P.A. Zyla, *et al.* (Particle Data Group), PTEP 2020, 083C01 (2020)

T. Hyodo, M. Niiyama, arXiv: 2010.07592 [hep-ph]

Y. Kamiya, T. Hyodo, PRC93, 035203 (2016); PTEP2017, 023D02 (2017)

T. Hyodo, JPS journal Vol. 75 No. 8, 478 (2020)

T. Kinugawa, T. Hyodo, in preparation