Size and structure of near-threshold states

Tetsuo Hyodo

Tokyo Metropolitan Univ.

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Classification of hadrons

PDG2018 : http://pdg.lbl.gov/

Usual hadrons <— quarks

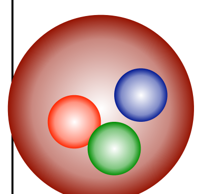
p	$1/2^+$	****
n	$1/2^+$	****
$N(1440)$	$1/2^+$	****
$N(1520)$	$3/2^-$	****
$N(1535)$	$1/2^-$	****
$N(1650)$	$1/2^-$	****
$N(1675)$	$5/2^-$	****
$N(1680)$	$5/2^+$	****
$N(1685)$		*
$N(1700)$	$3/2^-$	***
$N(1710)$	$1/2^+$	***
$N(1720)$	$3/2^+$	****
$N(1860)$	$5/2^+$	**
$N(1875)$	$3/2^-$	***
$N(1880)$	$1/2^+$	**
$N(1895)$	$1/2^-$	**
$N(1900)$	$3/2^+$	***
$N(1990)$	$7/2^+$	**
$N(2000)$	$5/2^+$	**
$N(2040)$	$3/2^+$	*
$N(2060)$	$5/2^-$	**
$N(2100)$	$1/2^+$	*
$N(2120)$	$3/2^-$	**
$N(2190)$	$7/2^-$	****
$N(2220)$	$9/2^+$	****
$N(2250)$	$9/2^-$	****
$N(2300)$	$1/2^+$	**
$N(2570)$	$5/2^-$	**
$N(2600)$	$11/2^-$	***
$N(2700)$	$13/2^+$	**

$\Delta(1232)$	$3/2^+$	****
$\Delta(1600)$	$3/2^+$	***
$\Delta(1620)$	$1/2^-$	****
$\Delta(1700)$	$3/2^-$	****
$\Delta(1750)$	$1/2^+$	*
$\Delta(1900)$	$1/2^-$	**
$\Delta(1905)$	$5/2^+$	****
$\Delta(1910)$	$1/2^+$	****
$\Delta(1920)$	$3/2^+$	***
$\Delta(1930)$	$5/2^-$	***
$\Delta(1940)$	$3/2^-$	**
$\Delta(1950)$	$7/2^+$	****
$\Delta(2000)$	$5/2^+$	**
$\Delta(2150)$	$1/2^-$	*
$\Delta(2200)$	$7/2^-$	*
$\Delta(2300)$	$9/2^+$	**
$\Delta(2350)$	$5/2^-$	*
$\Delta(2390)$	$7/2^+$	*
$\Delta(2400)$	$9/2^-$	**
$\Delta(2420)$	$11/2^+$	****
$\Delta(2750)$	$13/2^-$	**
$\Delta(2950)$	$15/2^+$	**
Λ	$1/2^+$	****
$\Lambda(1405)$	$1/2^-$	****
$\Lambda(1520)$	$3/2^-$	****
$\Lambda(1600)$	$1/2^+$	***
$\Lambda(1670)$	$1/2^-$	****
$\Lambda(1690)$	$3/2^-$	****
$\Lambda(1710)$	$1/2^+$	*
$\Lambda(1800)$	$1/2^-$	***
$\Lambda(1810)$	$1/2^+$	***
$\Lambda(1820)$	$5/2^+$	****
$\Lambda(1830)$	$5/2^-$	****
$\Lambda(1890)$	$3/2^+$	****
$\Lambda(2000)$		*
$\Lambda(2020)$	$7/2^+$	*
$\Lambda(2050)$	$3/2^-$	*
$\Lambda(2100)$	$7/2^-$	****
$\Lambda(2110)$	$5/2^+$	***
$\Lambda(2325)$	$3/2^-$	*
$\Lambda(2350)$	$9/2^+$	***
$\Lambda(2585)$		**

Σ^+	$1/2^+$	****
Σ^0	$1/2^+$	****
Σ^-	$1/2^+$	****
$\Sigma(1385)$	$3/2^+$	****
$\Sigma(1480)$		*
$\Sigma(1560)$		**
$\Sigma(1580)$	$3/2^-$	*
$\Sigma(1620)$	$1/2^-$	*
$\Sigma(1660)$	$1/2^+$	***
$\Sigma(1670)$	$3/2^-$	****
$\Sigma(1690)$		**
$\Sigma(1730)$	$3/2^+$	*
$\Sigma(1750)$	$1/2^-$	***
$\Sigma(1770)$	$1/2^+$	*
$\Sigma(1775)$	$5/2^-$	****
$\Sigma(1840)$	$3/2^+$	*
$\Sigma(1880)$	$1/2^+$	**
$\Sigma(1900)$	$1/2^-$	*
$\Sigma(1915)$	$5/2^+$	****
$\Sigma(1940)$	$3/2^+$	*
$\Sigma(1940)$	$3/2^-$	***
$\Sigma(2000)$	$1/2^-$	*
$\Sigma(2030)$	$7/2^+$	****
$\Sigma(2070)$	$5/2^+$	*
$\Sigma(2080)$	$3/2^+$	**
$\Sigma(2100)$	$7/2^-$	*
$\Sigma(2250)$		***
$\Sigma(2455)$		**
$\Sigma(2620)$		**
$\Sigma(3000)$		*
$\Sigma(3170)$		*

Ξ^0	$1/2^+$	****
Ξ^-	$1/2^+$	****
$\Xi(1530)$	$3/2^+$	****
$\Xi(1620)$		*
$\Xi(1690)$		***
$\Xi(1820)$	$3/2^-$	***
$\Xi(1950)$		***
$\Xi(2030)$	$\geq \frac{5}{2}?$	***
$\Xi(2120)$		*
$\Xi(2250)$		**
$\Xi(2370)$		**
$\Xi(2500)$		*
Ω^-	$3/2^+$	****
$\Omega(2250)^-$		***
$\Omega(2380)^-$		**
$\Omega(2470)^-$		**

Λ_c^+	$1/2^+$	****
$\Lambda_c(2595)^+$	$1/2^-$	***
$\Lambda_c(2625)^+$	$3/2^-$	***
$\Lambda_c(2765)^+$		*
$\Lambda_c(2880)^+$	$5/2^+$	***
$\Lambda_c(2940)^+$		***
$\Sigma_c(2455)$	$1/2^+$	****
$\Sigma_c(2520)$	$3/2^+$	***
$\Sigma_c(2800)$		***
Ξ_c^+	$1/2^+$	***
Ξ_c^0	$1/2^+$	***
$\Xi_c'^+$	$1/2^+$	***
$\Xi_c'^0$	$1/2^+$	***
$\Xi_c(2645)$	$3/2^+$	***
$\Xi_c(2790)$	$1/2^-$	***
$\Xi_c(2815)$	$3/2^-$	***
$\Xi_c(2930)$		*
$\Xi_c(2980)$		***
$\Xi_c(3055)$		***
$\Xi_c(3080)$		***
$\Xi_c(3123)$		*
Ω_c^0	$1/2^+$	***
$\Omega_c(2770)^0$	$3/2^+$	***
Ξ_{cc}^+		*
Λ_b^0	$1/2^+$	***
$\Lambda_b(5912)^0$	$1/2^-$	***
$\Lambda_b(5920)^0$	$3/2^-$	***
Σ_b	$1/2^+$	***
Σ_b^*	$3/2^+$	***
Ξ_b^0, Ξ_b^-	$1/2^+$	***
$\Xi_b'(5935)^-$	$1/2^+$	***
$\Xi_b(5945)^0$	$3/2^+$	***
$\Xi_b^*(5955)^-$	$3/2^+$	***
Ω_b^-	$1/2^+$	***



~ 150 baryons

LIGHT UNFLAVORED ($S = C = B = 0$)	$I^G(J^{PC})$
• $\pi^\pm$	$1^-(0^-)$
• π^0	$1^-(0^{-+})$
• η	$0^+(0^{-+})$
• $f_0(500)$	$0^+(0^{++})$
• $\rho(770)$	$1^+(1^{--})$
• $\omega(782)$	$0^-(1^{--})$
• $\eta'(958)$	$0^+(0^{-+})$
• $f_0(980)$	$0^+(0^{++})$
• $a_0(980)$	$1^-(0^{++})$
• $\phi(1020)$	$0^-(1^{--})$
• $h_1(1170)$	$0^-(1^{+-})$
• $b_1(1235)$	$1^+(1^{+-})$
• $a_1(1260)$	$1^-(1^{++})$
• $f_2(1270)$	$0^+(2^{++})$
• $f_1(1285)$	$0^+(1^{++})$
• $\eta(1295)$	$0^+(0^{-+})$
• $\pi(1300)$	$1^-(0^{-+})$
• $a_2(1320)$	$1^-(2^{++})$
• $f_0(1370)$	$0^+(0^{++})$
$h_1(1380)$	$?^-(1^{+-})$
• $\pi_1(1400)$	$1^-(1^{-+})$
• $\eta(1405)$	$0^+(0^{-+})$
• $f_1(1420)$	$0^+(1^{++})$
• $\omega(1420)$	$0^-(1^{--})$
$f_2(1430)$	$0^+(2^{++})$
• $a_0(1450)$	$1^-(0^{++})$
• $\rho(1450)$	$1^+(1^{--})$
• $\eta(1475)$	$0^+(0^{-+})$
• $f_0(1500)$	$0^+(0^{++})$
$f_1(1510)$	$0^+(1^{++})$
• $f_2'(1525)$	$0^+(2^{++})$
$f_2(1565)$	$0^+(2^{++})$
$\rho(1570)$	$1^+(1^{--})$
$h_1(1595)$	$0^-(1^{+-})$
• $\pi_1(1600)$	$1^-(1^{-+})$
$a_1(1640)$	$1^-(1^{++})$
$f_2(1640)$	$0^+(2^{++})$
• $\eta_2(1645)$	$0^+(2^{-+})$
• $\omega(1650)$	$0^-(1^{--})$
• $\omega_3(1670)$	$0^-(3^{--})$
• $\pi_2(1670)$	$1^-(2^{-+})$

LIGHT UNFLAVORED (continued)	$I^G(J^{PC})$
• $\phi(1680)$	$0^-(1^{--})$
• $\rho_3(1690)$	$1^+(3^{--})$
• $\rho(1700)$	$1^+(1^{--})$
$a_2(1700)$	$1^-(2^{++})$
• $f_0(1710)$	$0^+(0^{++})$
$\eta(1760)$	$0^+(0^{-+})$
• $\pi(1800)$	$1^-(0^{-+})$
$f_2(1810)$	$0^+(2^{++})$
$X(1835)$	$?^?(?^{-+})$
$X(1840)$	$?^?(?^{??})$
• $\phi_3(1850)$	$0^-(3^{--})$
$\eta_2(1870)$	$0^+(2^{-+})$
• $\pi_2(1880)$	$1^-(2^{-+})$
$\rho(1900)$	$1^+(1^{--})$
$f_2(1910)$	$0^+(2^{++})$
• $f_2(1950)$	$0^+(2^{++})$
$\rho_3(1990)$	$1^+(3^{--})$
• $f_2(2010)$	$0^+(2^{++})$
$f_0(2020)$	$0^+(0^{++})$
• $a_4(2040)$	$1^-(4^{++})$
• $f_4(2050)$	$0^+(4^{++})$
$\pi_2(2100)$	$1^-(2^{-+})$
$f_0(2100)$	$0^+(0^{++})$
$f_2(2150)$	$0^+(2^{++})$
$\rho(2150)$	$1^+(1^{--})$
• $\phi(2170)$	$0^-(1^{--})$
$f_0(2200)$	$0^+(0^{++})$
$f_J(2220)$	$0^+(2^{++}$ or $4^{++})$
$\eta(2225)$	$0^+(0^{-+})$
$\rho_3(2250)$	$1^+(3^{--})$
• $f_2(2300)$	$0^+(2^{++})$
$f_4(2300)$	$0^+(4^{++})$
$f_0(2330)$	$0^+(0^{++})$
• $f_2(2340)$	$0^+(2^{++})$
$\rho_5(2350)$	$1^+(5^{--})$
$a_6(2450)$	$1^-(6^{++})$
$f_6(2510)$	$0^+(6^{++})$

STRANGE ($S = \pm1$, $C = B = 0$)	$I(J^P)$
• $K^\pm$	$1/2(0^-)$
• K^0	$1/2(0^-)$
• K_S^0	$1/2(0^-)$
• K_L^0	$1/2(0^-)$
$K_0^*(800)$	$1/2(0^+)$
• $K^*(892)$	$1/2(1^-)$
• $K_1(1270)$	$1/2(1^+)$
• $K_1(1400)$	$1/2(1^+)$
• $K^*(1410)$	$1/2(1^-)$
• $K_0^*(1430)$	$1/2(0^+)$
• $K_2^*(1430)$	$1/2(2^+)$
$K(1460)$	$1/2(0^-)$
$K_2(1580)$	$1/2(2^-)$
$K(1630)$	$1/2(?^?)$
$K_1(1650)$	$1/2(1^+)$
• $K^*(1680)$	$1/2(1^-)$
• $K_2(1770)$	$1/2(2^-)$
• $K_3^*(1780)$	$1/2(3^-)$
• $K_2(1820)$	$1/2(2^-)$
$K(1830)$	$1/2(0^-)$
$K_0^*(1950)$	$1/2(0^+)$
$K_2^*(1980)$	$1/2(2^+)$
• $K_4^*(2045)$	$1/2(4^+)$
$K_2(2250)$	$1/2(2^-)$
$K_3(2320)$	$1/2(3^+)$
$K_5^*(2380)$	$1/2(5^-)$
$K_4(2500)$	$1/2(4^-)$
$K(3100)$	$?^?(?^{??})$

CHARMED ($C = \pm1$)	
• $D^\pm$	$1/2(0^-)$
• D^0	$1/2(0^-)$
• $D^*(2007)^0$	$1/2(1^-)$
• $D^*(2010)^\pm$	$1/2(1^-)$
• $D_0^*(2400)^0$	$1/2(0^+)$
$D_0^*(2400)^\pm$	$1/2(0^+)$
• $D_1(2420)^0$	$1/2(1^+)$
$D_1(2420)^\pm$	$1/2(?^?)$
$D_1(2430)^0$	$1/2(1^+)$
$D_2^*(2460)^0$	$1/2(2^+)$
$D_2^*(2460)^\pm$	$1/2(2^+)$
$D(2550)^0$	$1/2(0^-)$
$D(2600)$	$1/2(?^?)$
$D^*(2640)^\pm$	$1/2(?^?)$
$D(2750)$	$1/2(?^?)$

CHARMED, STRANGE ($C = S = \pm1$)	$I(J^P)$
• $D_s^\pm$	$0(0^-)$
• $D_s^{*\pm}$	$0(?^?)$
• $D_{s0}^*(2317)^\pm$	$0(0^+)$
• $D_{s1}(2460)^\pm$	$0(1^+)$
• $D_{s1}(2536)^\pm$	$0(1^+)$
• $D_{s2}(2573)$	$0(?^?)$
• $D_{s1}^*(2700)^\pm$	$0(1^-)$
$D_{sJ}^*(2860)^\pm$	$0(?^?)$
$D_{sJ}(3040)^\pm$	$0(?^?)$

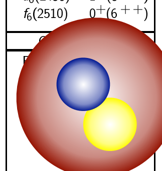
BOTTOM ($B = \pm1$)	
• $B^\pm$	$1/2(0^-)$
• B^0	$1/2(0^-)$
• $B^\pm/B^0$ ADMIXTURE	
• $B^\pm/B^0/B_s^0/b$-baryon ADMIXTURE	
V_{cb} and V_{ub} CKM Matrix Elements	
• B^*	$1/2(1^-)$
• $B_1(5721)^+$	$1/2(1^+)$
• $B_1(5721)^0$	$1/2(1^+)$
$B_J^*(5732)$	$?(?^?)$
• $B_2^*(5747)^+$	$1/2(2^+)$
• $B_2^*(5747)^0$	$1/2(2^+)$
• $B(5970)^+$	$?(?^?)$
• $B(5970)^0$	$?(?^?)$

BOTTOM, STRANGE ($B = \pm1$, $S = \mp1$)	
• B_s^0	$0(0^-)$
• B_s^*	$0(1^-)$
• $B_{s1}(5830)^0$	$0(1^+)$
• $B_{s2}^*(5840)^0$	$0(2^+)$
$B_{sJ}^*(5850)$	$?(?^?)$

BOTTOM, CHARMED ($B = C = \pm1$)	
• B_c^+	$0(0^-)$
$B_c(2S)^\pm$	$?^?(?^{??})$

$c\bar{c}$	$I^G(J^{PC})$
• $\eta_c(1S)$	$0^+(0^{-+})$
• $J/\psi(1S)$	$0^-(1^{--})$
• $\chi_{c0}(1P)$	$0^+(0^{++})$
• $\chi_{c1}(1P)$	$0^+(1^{++})$
• $h_c(1P)$	$?^?(1^{+-})$
• $\chi_{c2}(1P)$	$0^+(2^{++})$
• $\eta_c(2S)$	$0^+(0^{-+})$
• $\psi(2S)$	$0^-(1^{--})$
• $\psi(3770)$	$0^-(1^{--})$
$X(3823)$	$?^?(?^{?-})$
• $X(3872)$	$0^+(1^{++})$
• $X(3900)^\pm$	$?(1^+)$
$X(3900)^0$	$?(?^?)$
• $\chi_{c0}(3915)$	$0^+(0^{++})$
• $\chi_{c2}(2P)$	$0^+(2^{++})$
$X(3940)$	$?^?(?^{??})$
$X(4020)^\pm$	$?(?^?)$
• $\psi(4040)$	$0^-(1^{--})$
$X(4050)^\pm$	$?(?^?)$
$X(4140)$	$0^+(?^{?+})$
• $\psi(4160)$	$0^-(1^{--})$
$X(4160)$	$?^?(?^{??})$
$X(4230)$	$?^?(1^{--})$
$X(4240)^\pm$	$?^?(0^-)$
$X(4250)^\pm$	$?(?^?)$
• $X(4260)$	$?^?(1^{--})$
$X(4350)$	$0^+(?^{?+})$
• $X(4360)$	$?^?(1^{--})$
• $\psi(4415)$	$0^-(1^{--})$
• $X(4430)^\pm$	$?(1^+)$
• $X(4660)$	$?^?(1^{--})$

$b\bar{b}$	
$\eta_b(1S)$	$0^+(0^{-+})$
• $\Upsilon(1S)$	$0^-(1^{--})$
• $\chi_{b0}(1P)$	$0^+(0^{++})$
• $\chi_{b1}(1P)$	$0^+(1^{++})$
• $h_b(1P)$	$?^?(1^{+-})$
• $\chi_{b2}(1P)$	$0^+(2^{++})$
$\eta_b(2S)$	$0^+(0^{-+})$
• $\Upsilon(2S)$	$0^-(1^{--})$
• $\Upsilon(1D)$	$0^-(2^{--})$
• $\chi_{b0}(2P)$	$0^+(0^{++})$
• $\chi_{b1}(2P)$	$0^+(1^{++})$
$h_b(2P)$	$?^?(1^{+-})$
• $\chi_{b2}(2P)$	$0^+(2^{++})$
• $\Upsilon(3S)$	$0^-(1^{--})$
• $\chi_{b1}(3P)$	$0^+(1^{++})$
• $\Upsilon(4S)$	$0^-(1^{--})$
$X(10610)^\pm$	$1^+(1^+)$
$X(10610)^0$	$1^+(1^+)$
$X(10650)^\pm$	$?^+(1^+)$
• $\Upsilon(10860)$	$0^-(1^{--})$
• $\Upsilon(11020)$	$0^-(1^{--})$



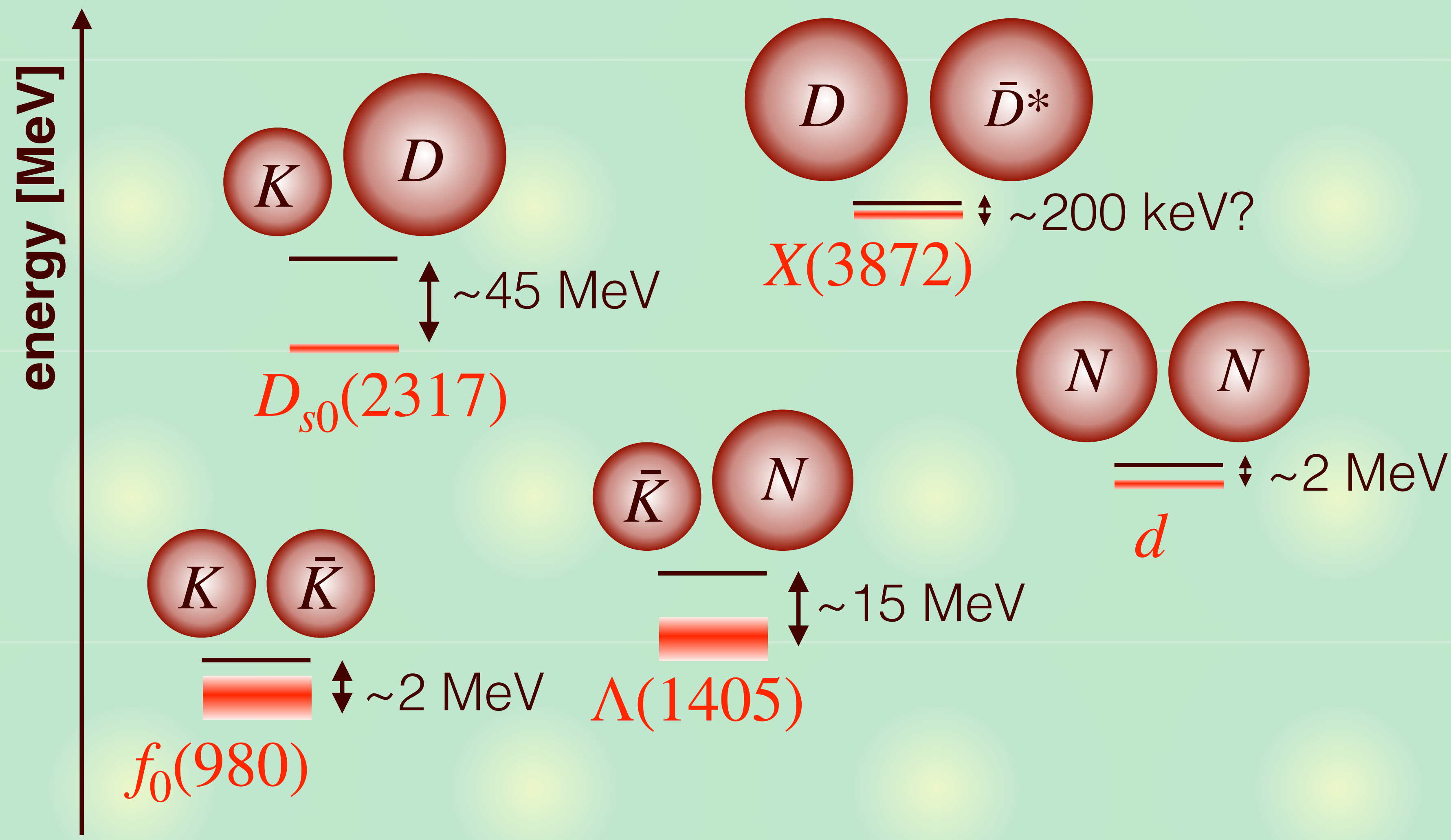
~ 210 mesons

Hadron clusters!

Can one construct a hadron from hadrons?

Hadron clusters

Hadrons near an s-wave two-body threshold



"hadronic molecules" (various flavors, baryon numbers, …)

Two-body universal physics

Near-threshold s-wave state: universal physics

E. Braaten, H.-W. Hammer, Phys. Rept. 428, 259 (2006);
P. Naidon, S. Endo, Rept. Prog. Phys. 80, 056001 (2017)

- scattering length $|a| \gg$ interaction range R_0
- size of (quasi-)bound state $\sim |a|$: loosely bound
- relation with eigenenergy E

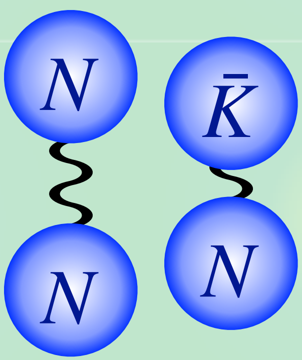
$$a(E) \sim \frac{i}{k} = \frac{i}{\sqrt{2\mu E}}$$

vdW

Examples: d, $\Lambda(1405)$, $^4\mathrm{He}$ dimer

	NN [fm]	$\bar{K}N$ [fm]	$^4\mathrm{He}$ [a_0]
$a(E)$	4.3	1.2-0.8i	178
a_{emp}	5.1	1.4-0.9i	189
R_0	1.4	0.4	10

strong



$^4\mathrm{He}$

Wave function

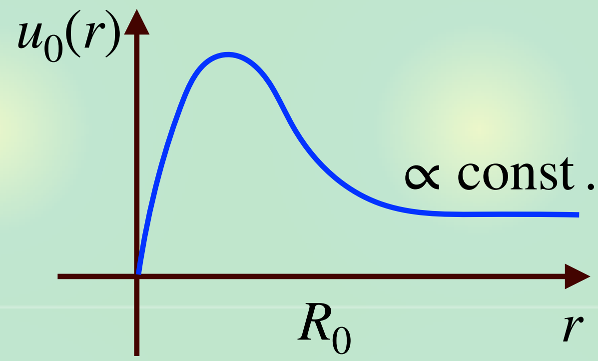
Why is a weakly bound s-wave state spatially large?

- Radial Schrödinger equation

$$-\frac{\hbar^2}{2\mu}\frac{d^2}{dr^2}u_0(r) + V(r)u_0(r) = Eu_0(r), \quad \psi_0(r) = Y_0^0\frac{u_0(r)}{r}$$

- at large distance ($V(r) = 0$) with zero energy ($E = 0$)

$$u_0(r) = C(r - a), \quad (r > R_0)$$



- scattering length a : intercept of $u_0(r)$
- bound state with $B = 0 \Rightarrow |a| = \infty$

Wave function is not normalizable

$$\int d^3r\,|\psi_0(r)|^2 = \int_0^\infty dr\,|u_0(r)|^2 = \infty$$

$B = 0$ state is not a bound state (zero-energy resonance)

Consequences

Mean squared radius

$$\langle r^2 \rangle = \int d^3r\ r^2 |\psi_0(r)|^2 = \int_0^\infty dr\ r^2 |u_0(r)|^2 = \infty$$

—> size of $B = 0$ state is infinitely large

Compositeness X (weight of the composite component)

T. Hyodo, Phys. Rev. C 90, 055208 (2014)

$$\Psi = \begin{pmatrix} \tilde{\psi}_0(p) \\ c \\ \vdots \end{pmatrix}$$

$$X = \int \frac{d^3p}{(2\pi)^3}\, |\tilde{\psi}_0(p)|^2 = \int d^3r\ |\psi_0(r)|^2$$ **infinite**

$$Z = |c|^2 + \cdots$$ **finite**

$$X + Z = 1$$

—> $B = 0$ state is completely composite ($X = 1,\ \ Z = 0$)

Weakly bound state ($B \neq 0$, except for fine tuning)

- large spatial size and composite dominance

p-wave state

What about p-wave states?

- Radial Schrödinger equation

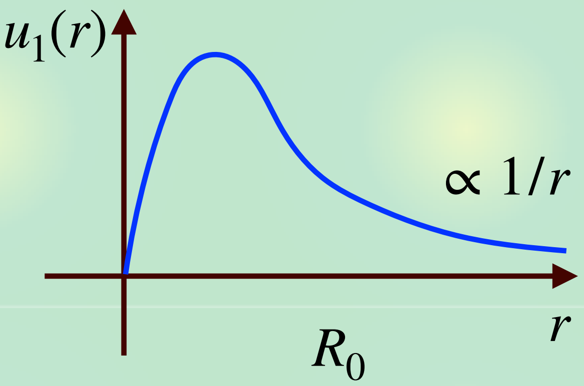
$$-\frac{\hbar^2}{2\mu}\frac{d^2}{dr^2}u_1(r) + V(r)u_1(r) + \frac{\hbar^2}{\mu r^2}u_1(r) = Eu_1(r), \quad \psi_1(\mathbf{r}) = Y_1^m(\Omega)\frac{u_1(r)}{r}$$

- at large distance ($V(r) = 0$) with zero energy ($E = 0$)

$$u_1(r) = Cr^s, \quad s = 2, -1$$

- bound state with $B = 0$

$$u_1(r) = \frac{C}{r}, \quad (r > R_0)$$



Wave function is normalizable

$$\int d^3r\,|\psi_1(\mathbf{r})|^2 = \int_0^\infty dr\,|u_1(r)|^2 < \infty$$

$B = 0$ state is a bound state

Consequences (p-wave)

Mean squared radius

$$\langle r^2\rangle = \int_0^\infty dr\ r^2 |u_1(r)|^2 = A_{\mathrm{in}} + |C|^2 \int_{R_0}^\infty dr\ 1 = \infty \qquad \leftarrow u_1(r) = \frac{C}{r}$$

—> size of $B = 0$ state is infinitely large

Compositeness X

T. Hyodo, Phys. Rev. C 90, 055208 (2014)

$$\Psi = \begin{pmatrix} \tilde{\psi}_1(\mathbf{p}) \\ c \\ \vdots \end{pmatrix}$$

$$X = \int \frac{d^3p}{(2\pi)^3}\ |\tilde{\psi}_1(\mathbf{p})|^2 = \int d^3r\ |\psi_1(\mathbf{r})|^2$$ **finite**

$$Z = |c|^2 + \cdots$$ **finite**

$$X + Z = 1$$

—> $B = 0$ can have any structure ($0 < X < 1$)

Weakly bound p-wave state ($B \neq 0$, except for fine tuning)

- large spatial size?

Higher partial waves

Wave function of a $B=0$ state with angular momentum ℓ

$$u_\ell(r) = \frac{C}{r^\ell} \quad (r > R_0)$$

- Mean squared radius

$$\langle r^2 \rangle = \int_0^\infty dr\ r^2 |u_\ell(r)|^2 = A_{\rm in} + |C|^2 \int_{R_0}^\infty dr\ r^{2-2\ell} = \infty \quad \text{if} \quad \ell \leq 1$$

—> diverges only for s- and p-waves

Generalized radius

$$\sqrt[n]{\langle r^n \rangle} = \sqrt[n]{\int d^3r\ r^n |\psi(r)|^2} = \infty \quad \text{if} \quad n - 2\ell \geq 0$$

- s-wave: radius diverges for all $n \geq 0$ <— universality

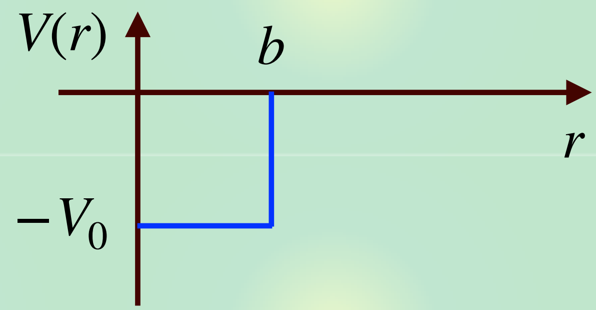
- p-wave: radius diverges for $n \geq 2$

- ℓ-th wave: radius diverges for $n \geq 2\ell$

example: square well potential

Square well potential

$$V(r) = \begin{cases} -V_0 & 0 \le r \le b \\ 0 & b < r \end{cases}$$



- s-wave

$$\langle r^2 \rangle = \left[\frac{b}{2} - \frac{\sin(2kb)}{4k} + \frac{\sin^2(kb)}{2\kappa} \right]^{-1} \left[\frac{1}{k^3} \left(\frac{(kb)^3}{6} + \frac{1-2(kb)^2}{8} \sin(2kb) - \frac{kb}{4} \cos(2kb) \right) + \frac{\sin^2(kb)}{2\kappa^3} \left[(\kappa b)^2 + \kappa b + \frac{1}{2} \right] \right]$$

$$k = \frac{\sqrt{2\mu(V_0 - B)}}{\hbar}, \quad \kappa = \frac{\sqrt{2\mu B}}{\hbar}$$

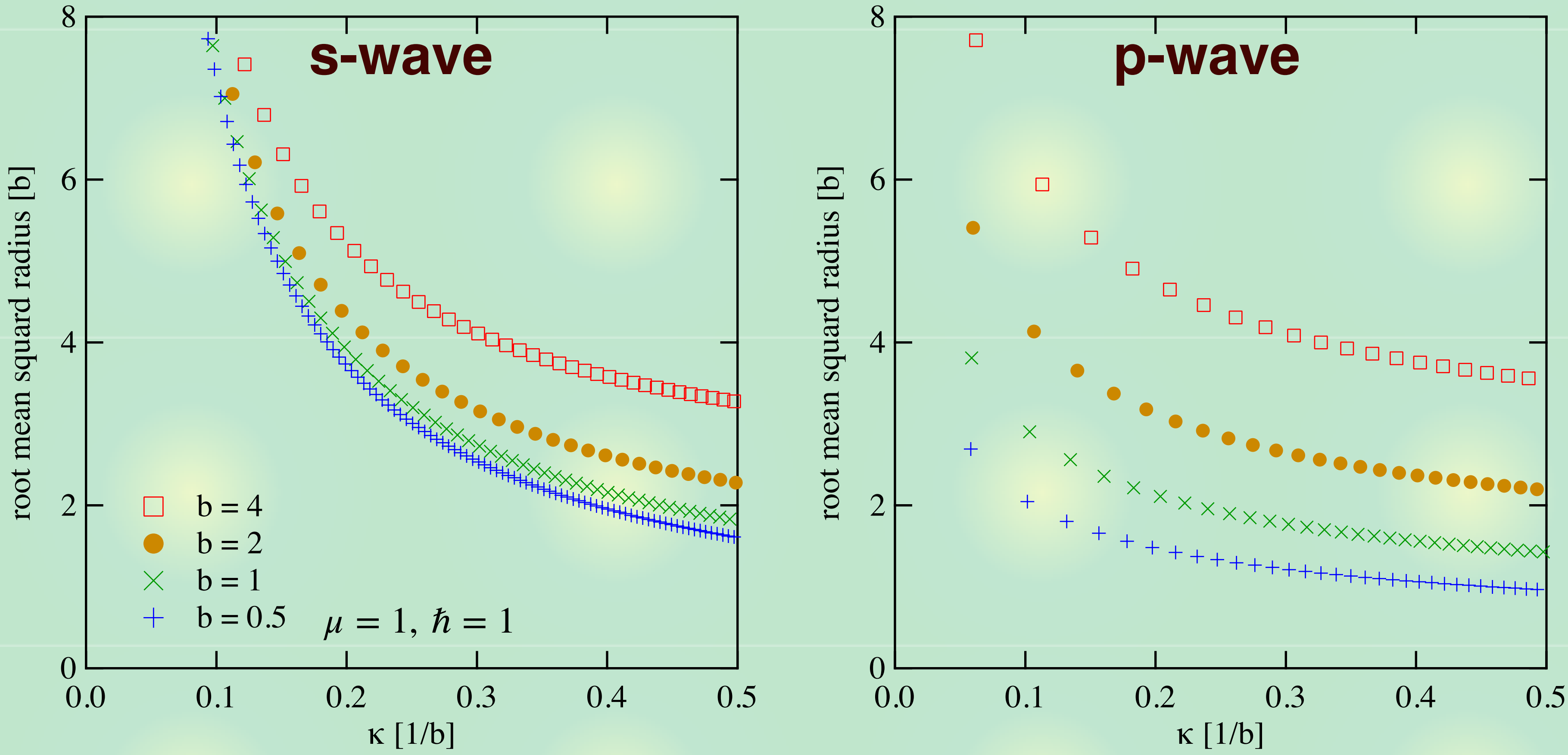
- p-wave

$$\langle r^2 \rangle = \left[\frac{b}{2} + \frac{\sin(2kb)}{4k} - \frac{\sin^2(kb)}{k^2 b} + \frac{\kappa \left(\sin(kb) - kb\cos(kb) \right)^2}{2k^2 \left(1 + \kappa b \right)^2} \left(1 + \frac{2}{\kappa b} \right) \right]^{-1}$$

$$\times \left[\frac{1}{k^3} \left(\frac{3kb + (kb)^3}{6} + \frac{-5 + 2(kb)^2}{8} \sin(2kb) + \frac{3kb}{4} \cos(2kb) \right) + \frac{\left(\sin(kb) - kb\cos(kb) \right)^2}{2\kappa k^2 \left(1 + \kappa b \right)^2} \left[(\kappa b)^2 + 3\kappa b + \frac{5}{2} \right] \right]$$

Demonstration

Root mean squared radius $\sqrt{\langle r^2 \rangle}$ v.s. $\kappa = \sqrt{2\mu B}/\hbar$



Both radii diverge in the weak binding limit ($\kappa \to 0$).

Dependence on b (well width) seems different.

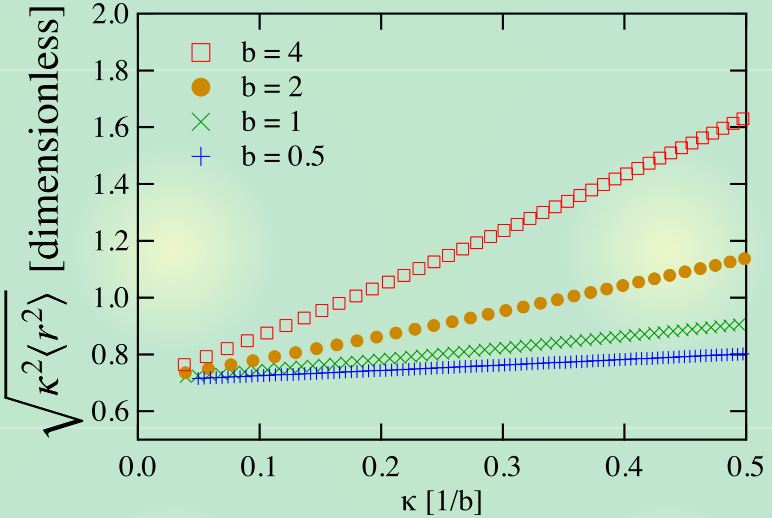
Weak-binding behavior

Behavior near the weak-binding limit

- s-wave

$$\langle r^2 \rangle = \frac{1}{2\kappa^2} + \mathcal{O}(\kappa^{-1})$$

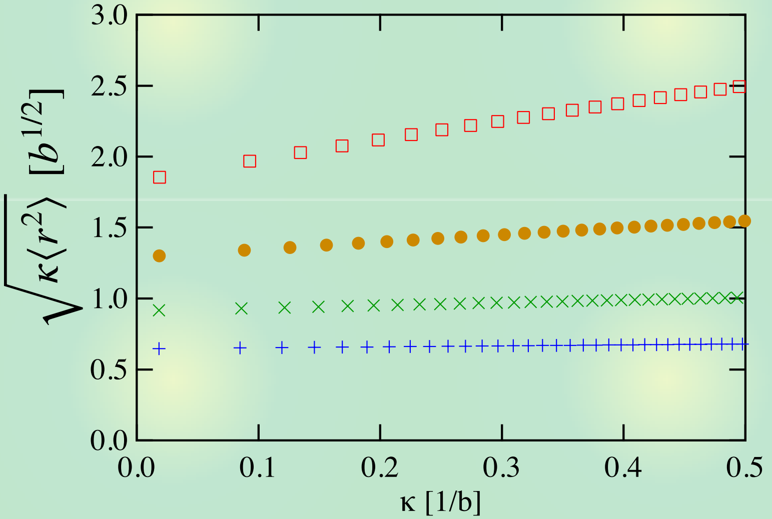
independent of b (well width)
<— universality



- p-wave

$$\langle r^2 \rangle = \frac{5b}{6}\frac{1}{\kappa} + \mathcal{O}(\kappa^0)$$

depends on b (well width)
non-universal divergence



Summary

- **Size and structure of $B = 0$ states**

	$\langle r^2 \rangle$	compositeness	$\int d^3r \lvert\psi\rvert^2$
s-wave	∞ (universal)	$X = 1$	∞
p-wave	∞ (non-universal)	$0 < X < 1$	finite

- **Implication: large mean squared radius of near-threshold ($B \neq 0$) p-wave bound states**

- **What about resonances?**

T. Hyodo, in preparation