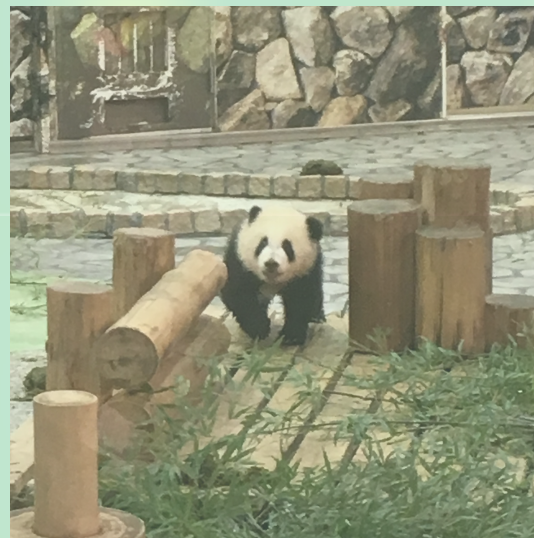


Status of $\Lambda(1405)$ in chiral dynamics




Tetsuo Hyodo

Tokyo Metropolitan Univ.

2019, Sep. 23rd 1



Introduction



Pole structure of the $\Lambda(1405)$ region

- Chiral SU(3) dynamics

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881, 98 (2012);
M. Tanabashi, *et al.* (Particle Data Group), PRD 98, 030001 (2018)



Some comments from recent studies

- K^-p correlation function —> Talk by Y. Kamiya (Friday)

ALICE collaboration, arXiv:1905.13470 [nucl-ex];
Y. Kamiya, T. Hyodo, K Morita, A. Ohnishi, in preparation

- Pole and finite volume spectrum

Y. Tsuchida, T. Hyodo, Phys. Rev. C97, 0552113 (2018)

\bar{K} meson and $\bar{K}N$ interaction

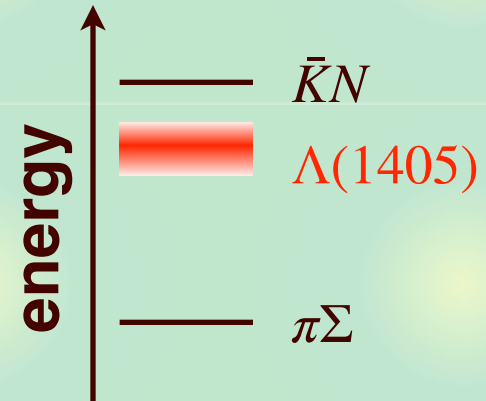
Two aspects of K/\bar{K} meson

- **NG boson** of chiral $SU(3)_R \otimes SU(3)_L \rightarrow SU(3)_V$
- **Massive** due to strange quark: $m_K \sim 496$ MeV
- > **Spontaneous/explicit** symmetry breaking

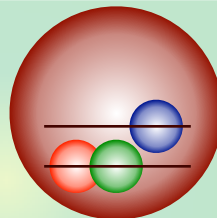
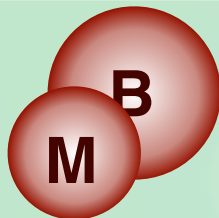
$\bar{K}N$ interaction ...

T. Hyodo, D. Jido, PPNP 67, 55 (2012)

- is coupled with $\pi\Sigma$ channel
- generates $\Lambda(1405)$ below threshold



molecule



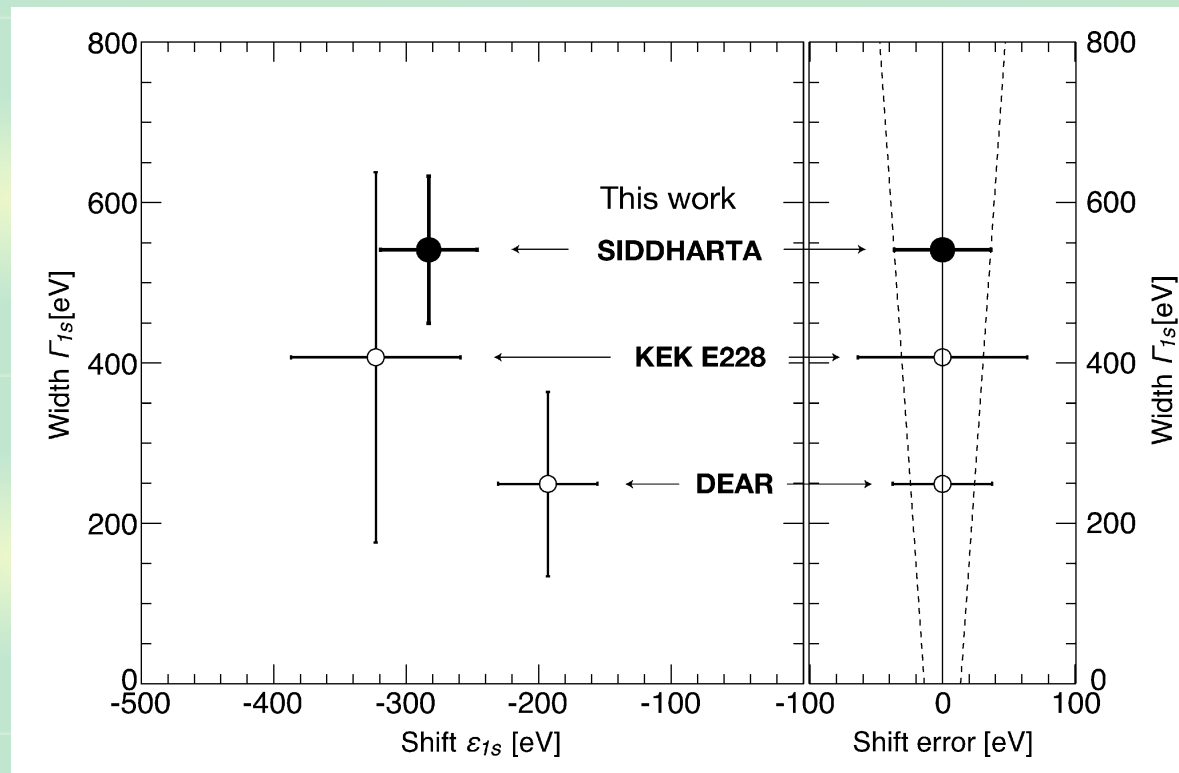
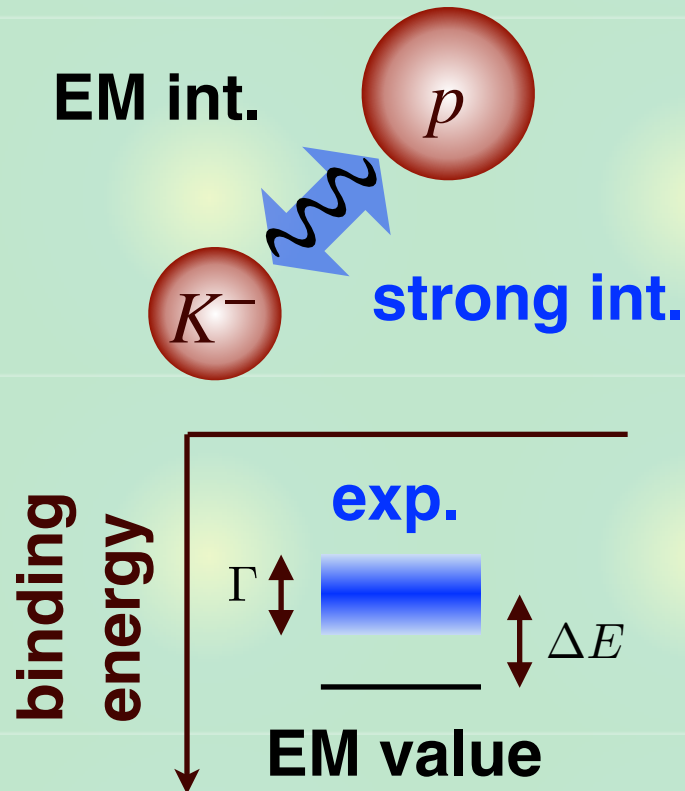
three-quark

- is fundamental building block for \bar{K} -nuclei, \bar{K} in medium, ...

SIDDHARTA measurement

Precise measurement of the kaonic hydrogen X-rays

M. Bazzi, *et al.*, PLB 704, 113 (2011); NPA 881, 88 (2012)



- Shift and width of atomic state \longleftrightarrow K^-p scattering length

U.G. Meissner, U. Raha, A. Rusetsky, Eur. Phys. J. C35, 349 (2004)

Quantitative constraint on the $\bar{K}N$ interaction at fixed energy

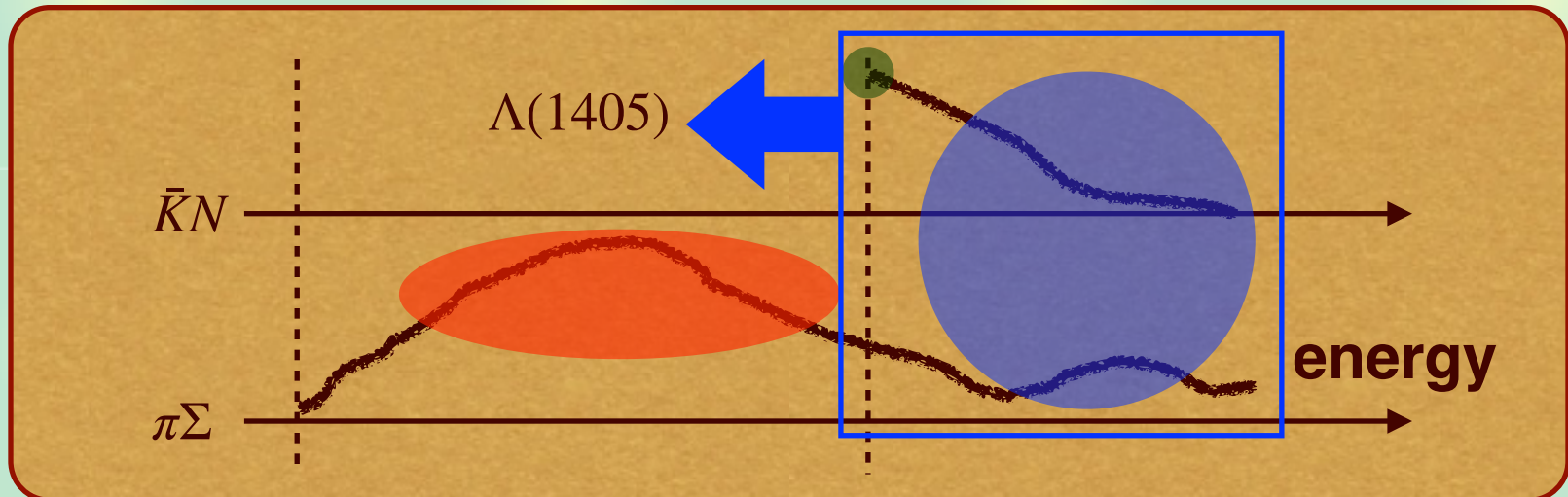
Strategy for $\bar{K}N$ interaction

Above the $\bar{K}N$ threshold: direct constraints

- K^-p total cross sections (old data)
- $\bar{K}N$ threshold branching ratios (old data)
- K^-p scattering length (new data: SIDDHARTA)

Below the $\bar{K}N$ threshold: indirect constraints

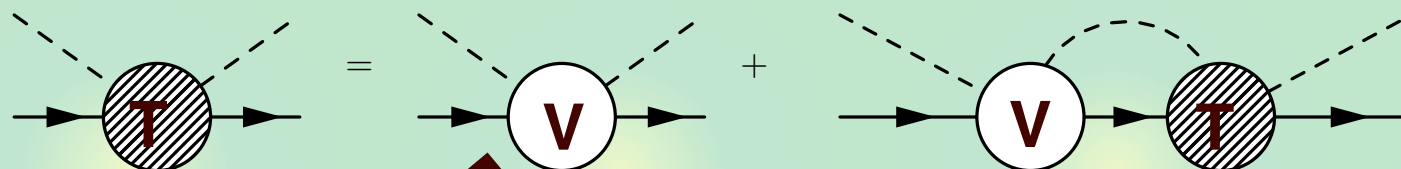
- $\pi\Sigma$ mass spectra (new data: LEPS, CLAS, HADES,...)



Construction of the realistic amplitude

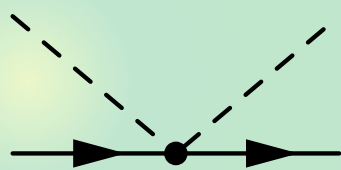
Chiral SU(3) coupled-channels ($\bar{K}N, \pi\Sigma, \pi\Lambda, \eta\Lambda, \eta\Sigma, K\Xi$) approach

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881 98 (2012)



Chiral perturbation theory

1) TW term

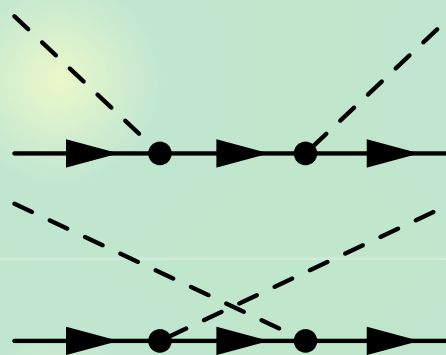


$\mathcal{O}(p)$

6 cutoffs

TW model

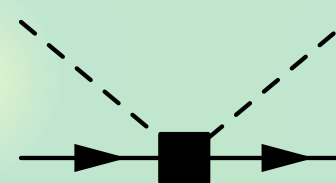
2) Born terms



$\mathcal{O}(p)$

TWB model

3) NLO terms



$\mathcal{O}(p^2)$

7 LECs

NLO model

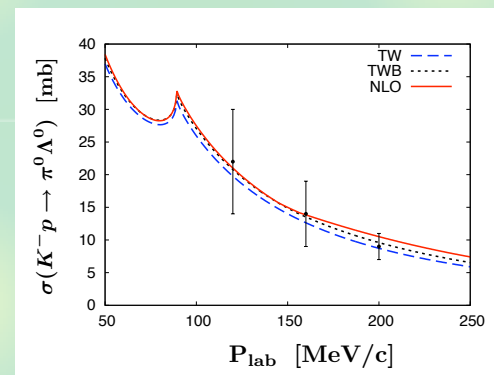
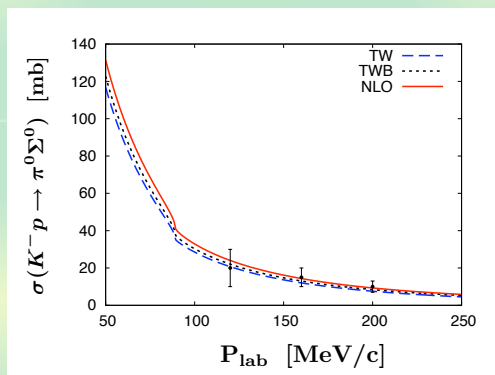
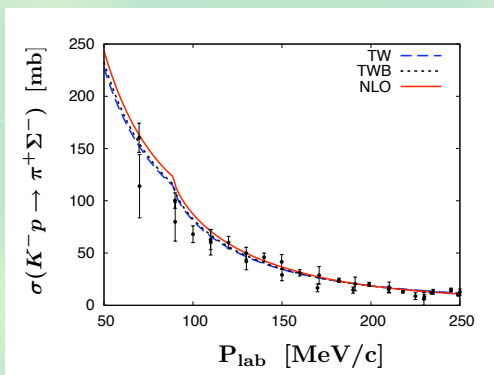
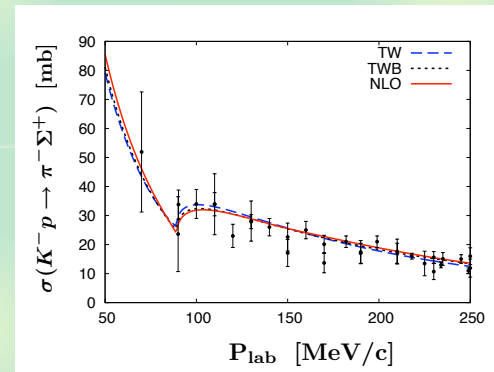
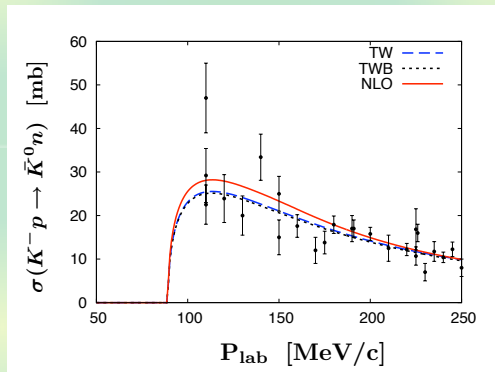
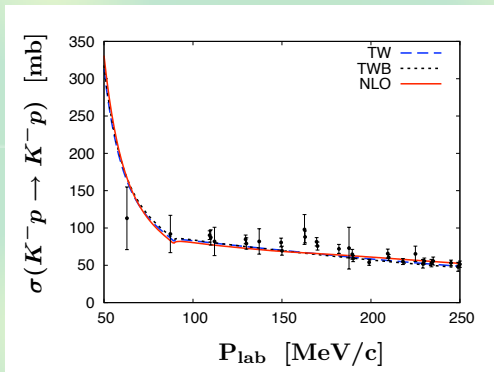
Best-fit results

K at rest

	TW	TWB	NLO	Experiment
ΔE [eV]	373	377	306	$283 \pm 36 \pm 6$ [10]
Γ [eV]	495	514	591	$541 \pm 89 \pm 22$ [10]
γ	2.36	2.36	2.37	2.36 ± 0.04 [11]
R_n	0.20	0.19	0.19	0.189 ± 0.015 [11]
R_c	0.66	0.66	0.66	0.664 ± 0.011 [11]
$\chi^2/\text{d.o.f}$	1.12	1.15	0.96	

} **SIDDHARTA**
 } **Branching ratios**

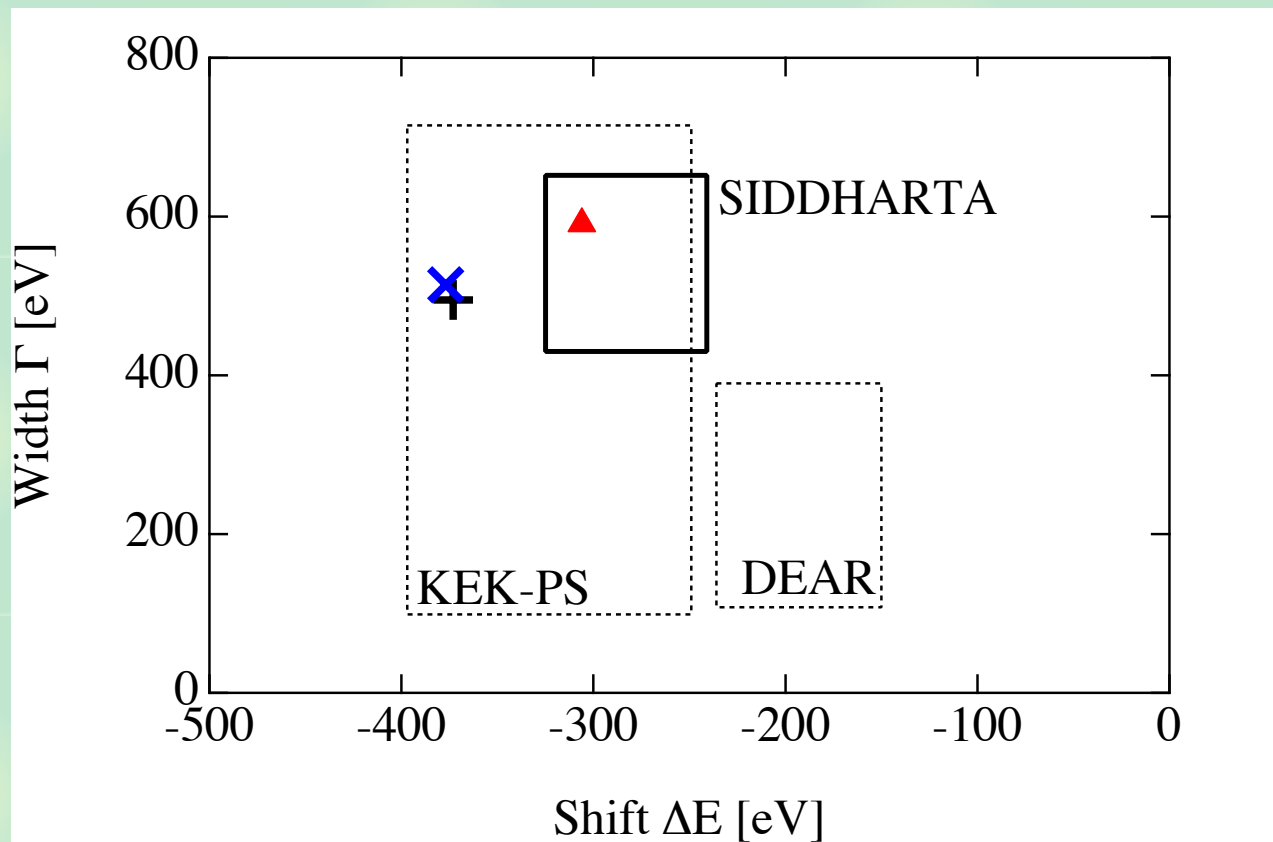
K^-p cross sections



Accurate description of all existing data ($\chi^2/\text{d.o.f} \sim 1$)

Comparison with SIDDHARTA

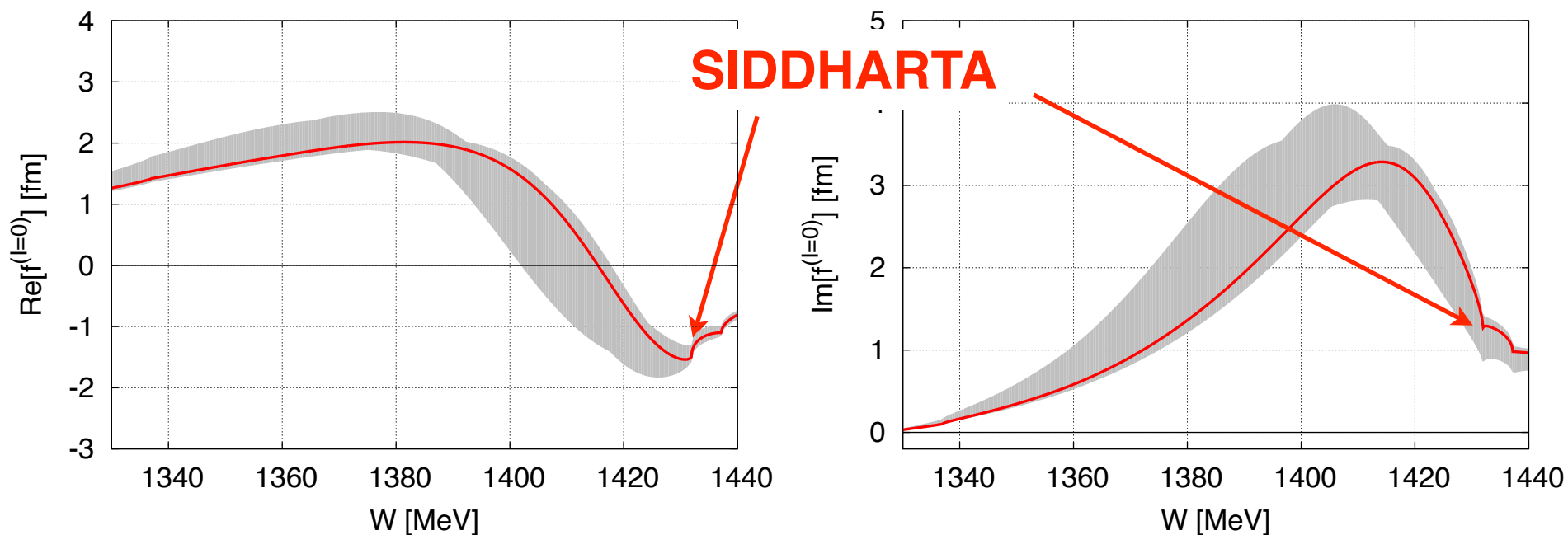
	TW	TWB	NLO
$\chi^2/\text{d.o.f.}$	1.12	1.15	0.957



TW and **TWB** are reasonable, while best-fit requires **NLO**.

Subthreshold extrapolation

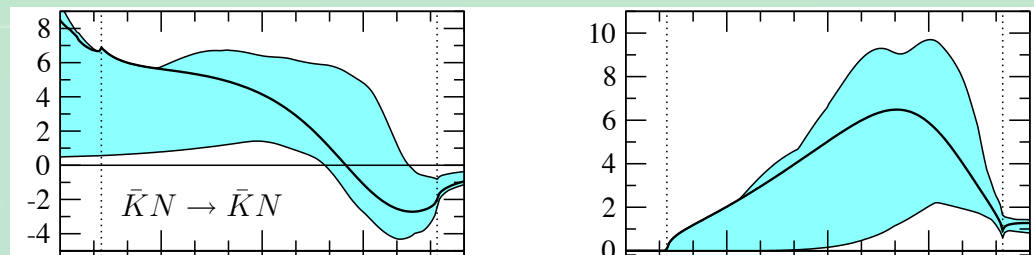
Uncertainty of $\bar{K}N \rightarrow \bar{K}N(I=0)$ amplitude below threshold



Y. Kamiya, K. Miyahara, S. Ohnishi, Y. Ikeda, T. Hyodo, E. Oset, W. Weise, NPA 954, 41 (2016)

- c.f. without SIDDHARTA

R. Nissler, Doctoral Thesis (2007)



SIDDHARTA is essential for subthreshold extrapolation.

Extrapolation to complex energy: two poles

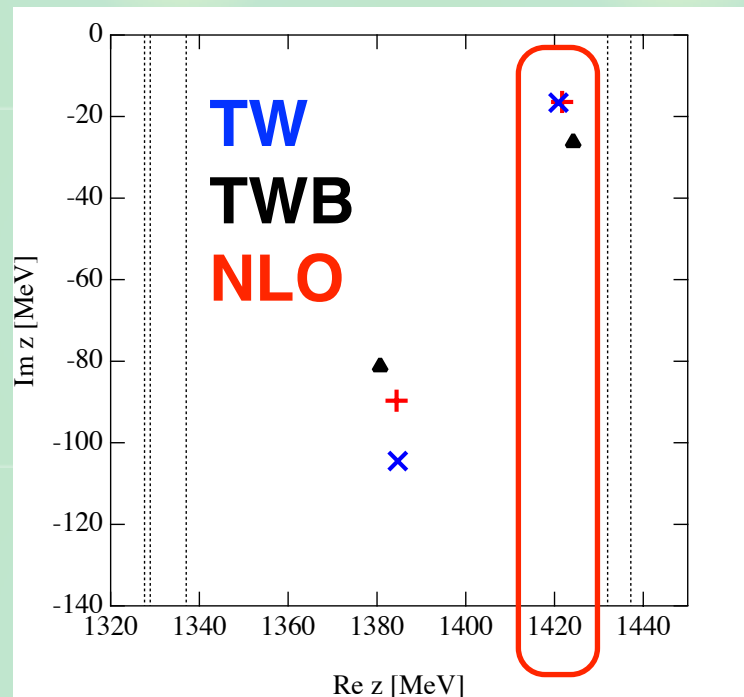
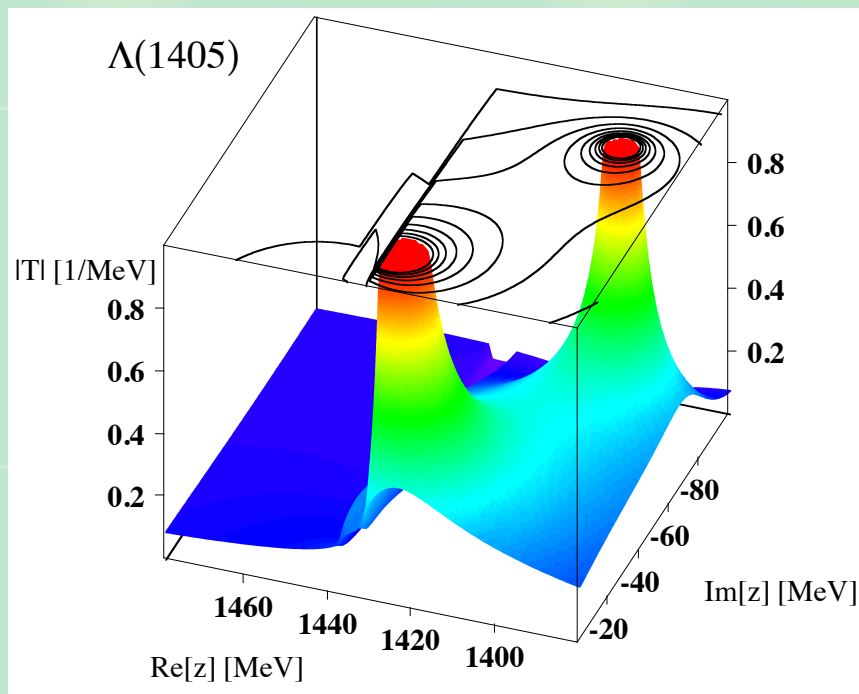
Two poles: superposition of two states

J.A. Oller, U.G. Meissner, PLB 500, 263 (2001);

D. Jido, J.A. Oller, E. Oset, A. Ramos, U.G. Meissner, NPA 723, 205 (2003);

T. Hyodo, W. Weise, PRC 77, 035204 (2008)

- Attractions of TW in 1 and 8 ($\bar{K}N$ and $\pi\Sigma$) channels



NLO analysis confirms the two-pole structure.

PDG changes

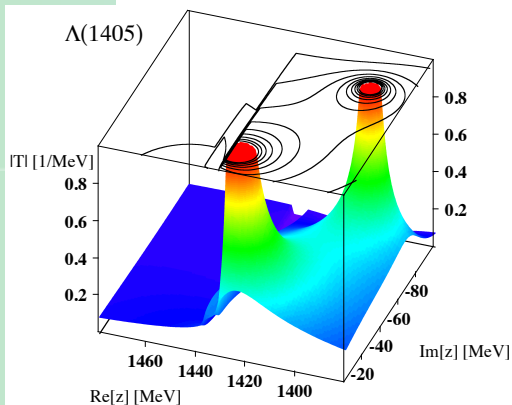
PDG particle listing of $\Lambda(1405)$

M. Tanabashi, et al., PRD 98, 030001 (2018), <http://pdg.lbl.gov/>

$\Lambda(1405) 1/2^-$

$I(J^P) = 0(\frac{1}{2}^-)$ Status **2014**

The nature of the $\Lambda(1405)$ has been a puzzle for decades: t... quark state or hybrid; two poles or one. We cannot here survey the rather extensive literature. See, for example, CIEPLY 10, KISSLINGER 11, SEKIHARA 11, and SHEVCHENKO 12A for discussions and earlier references.



of the chiral-unitary community 400-MeV region. ZYCHOR 08 inst the two-pole model, but this REVAI 09, which finds little basis io-pole models; and IKEDA 12,

1405) fits nicely into a $J^P = 1/2^-$ member members are the $\Lambda_c(2595)^+$, Fig. 1 of our note on "Charmed

MASS

TECN COMMENT

$\Lambda(1405) 1/2^-$

$I(J^P) = 0(\frac{1}{2}^-)$ Status **2018**

In the 1998 Note on the $\Lambda(1405)$ in PDG 98, R.H. Dalitz discussed the S-shaped cusp behavior of the intensity at the $N\bar{K}$ threshold observed in THOMAS 73 and HEMINGWAY 85. He commented that this behavior "is characteristic of S-wave coupling; the other below threshold hyperon, the $\Sigma(1385)$, has no such threshold distortion because its $N\bar{K}$ coupling is P-wave. For $\Lambda(1405)$ this asymmetry is the sole direct evidence that $J^P = 1/2^-$."

A recent measurement by the CLAS collaboration, MORIYA 14, definitely established the long-assumed $J^P = 1/2^-$ spin-parity assignment of the $\Lambda(1405)$. The experiment produced the $\Lambda(1405)$ spin-polarized in the photoproduction process $\gamma p \rightarrow K^+ \Lambda(1405)$ and measured the decay of the $\Lambda(1405)$ (polarized) $\rightarrow \Sigma^+$ (polarized) π^- . The observed isotropic decay of $\Lambda(1405)$ is consistent with spin $J = 1/2$. The polarization transfer to the Σ^+ (polarized) direction revealed negative parity, and thus established $J^P = 1/2^-$.

See the related review(s):

[Pole Structure of the \$\Lambda\(1405\)\$ Region](#)

$\Lambda(1405)$ REGION POLE POSITIONS

REAL PART

VALUE (MeV)	DOCUMENT ID	TECN
• • • We do not use the following data for averages, fits, limits, etc. • • •		
1429^{+8}_{-7}	1 MAI	15 DPWA
1325^{+15}_{-15}	2 MAI	15 DPWA
1434^{+2}_{-2}	3 MAI	15 DPWA
1330^{+4}_{-5}	4 MAI	15 DPWA
1421^{+3}_{-2}	5 GUO	13 DPWA
1388^{+9}_{-9}	6 GUO	13 DPWA
1424^{+7}_{-23}	7 IKEDA	12 DPWA
1381^{+18}_{-6}	8 IKEDA	12 DPWA

105. Pole Structure of the $\Lambda(1405)$ Region

Written November 2015 by Ulf-G. Meißner (Bonn Univ. / FZ Jülich) and Tetsuo Hyodo (YITP, Kyoto Univ.).

The $\Lambda(1405)$ resonance emerges in the meson-baryon scattering amplitude with the strangeness $S = -1$ and isospin $I = 0$. It is the archetype of what is called a dynamically generated resonance, as pioneered by Dalitz and Tuan [1]. The most powerful and

- **Our analysis (+ 2 other groups)**

- **Pole positions are now tabulated, prior to mass/width.**

Correlation function

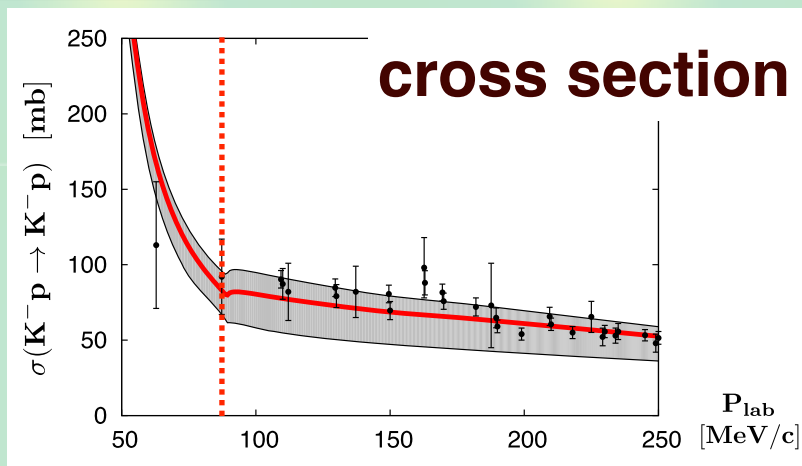
K^-p total cross sections

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011)

- Old bubble chamber data

K^-p correlation function

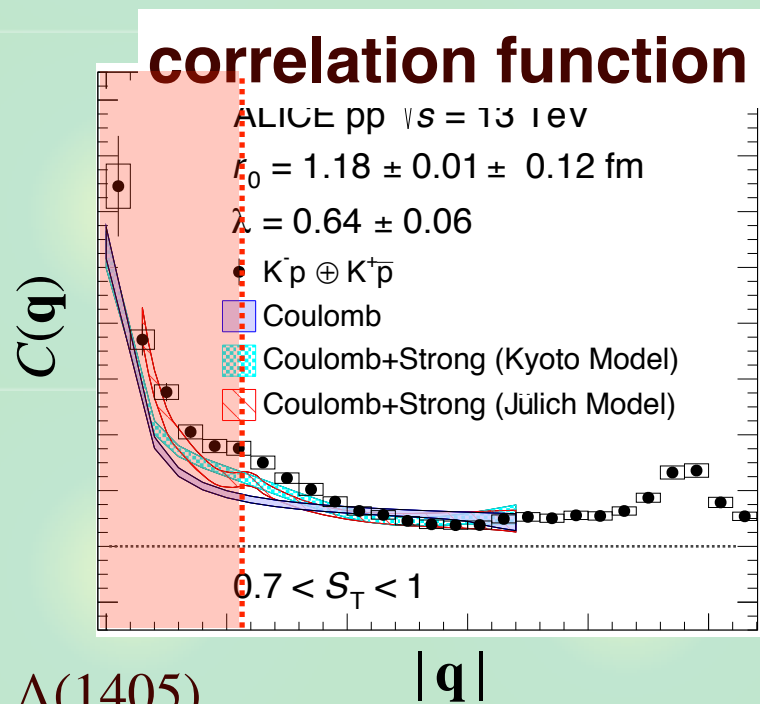
ALICE collaboration, arXiv:1905.13470 [nucl-ex]



$$C(\mathbf{q}) = \frac{N_{K^-p}(\mathbf{p}_{K^-}, \mathbf{p}_p)}{N_{K^-}(\mathbf{p}_{K^-})N_p(\mathbf{p}_p)}$$

- Excellent **precision** (\bar{K}^0n cusp)
- Low-energy data **below** \bar{K}^0n

$$C(\mathbf{q}) \simeq \int d^3\mathbf{r} S(\mathbf{r}) |\Psi_{\mathbf{q}}^{(-)}(\mathbf{r})|^2$$



—> Important constraint on $\bar{K}N$ and $\Lambda(1405)$

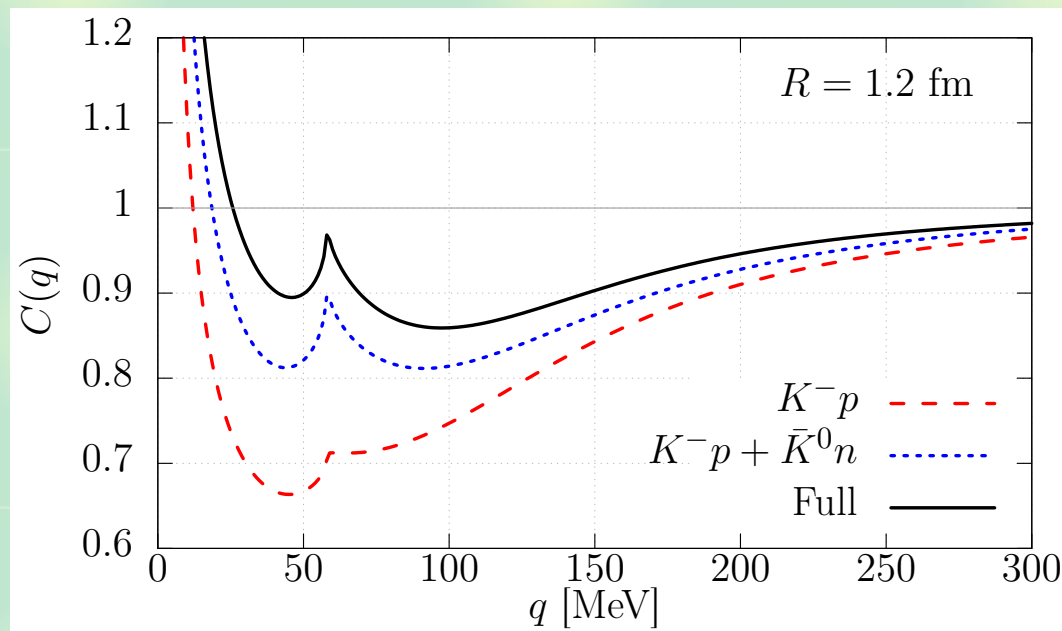
Results

Developing theoretical framework to calculate $C(q)$ with

- interaction: coupled-channel $\bar{K}N-\pi\Sigma-\pi\Lambda$ potential

K. Miyahara, T. Hyodo, W. Weise, PRC98, 025201 (2018)

- static spherical source $S(r)$, weight $\omega_i = 1$



$$\begin{pmatrix} \psi_{K^-p} \\ \psi_{\bar{K}^0n} \\ \psi_{\pi^+\Sigma^-} \\ \psi_{\pi^0\Sigma^0} \\ \psi_{\pi^-\Sigma^+} \\ \psi_{\pi^0\Lambda} \end{pmatrix}$$

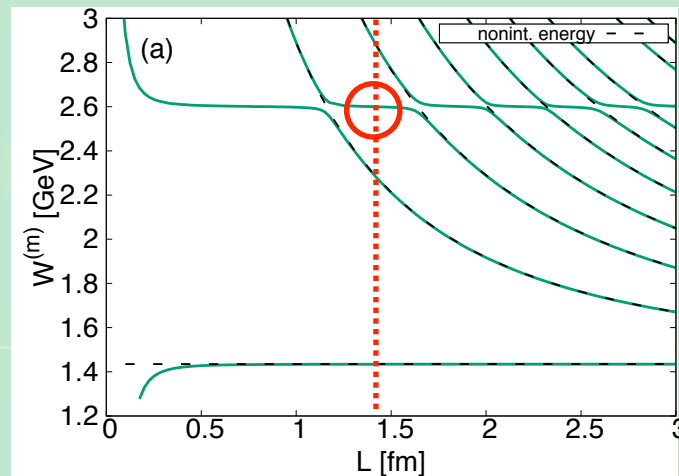
- \bar{K}^0n cusp is prominent with inclusion of $\psi_{\bar{K}^0n}$
- Coupled channels **enhance** K^-p correlation

Resonance and finite volume spectrum

Sharp resonance and finite volume spectrum (toy model)

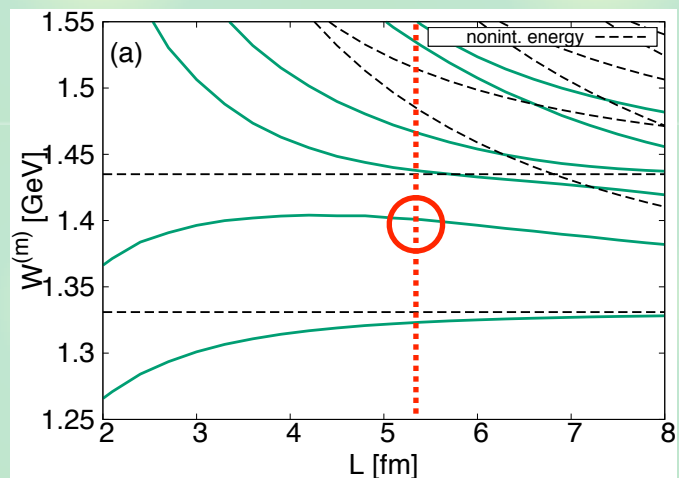
Y. Tsuchida, T. Hyodo, Phys. Rev. C97, 0552113 (2018)

- dashed: free eigenenergy
- solid: with interaction
- **additional** energy level




$\Lambda(1405)$ case (amplitude with two poles)

- Only **one** additional level
- # of additional energy levels = # of $\pi/2$ **crossings** of phase shift \neq # of poles




$\bar{K}N$
 $\pi\Sigma$


Summary

 Pole structure of the $\Lambda(1405)$ region is now well constrained by the experimental data. Nominal $\Lambda(1405)$ can be a superposition of **two states**.

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881, 98 (2012);
M. Tanabashi, *et al.* (Particle Data Group), PRD 98, 030001 (2018)

 K^-p correlation function will be a new **accurate constraint** on the $\Lambda(1405)$

ALICE collaboration, arXiv:1905.13470 [nucl-ex];
Y. Kamiya, T. Hyodo, K Morita, A. Ohnishi, in preparation

 Finite volume spectrum does **not directly indicate** the pole structure.

Y. Tsuchida, T. Hyodo, Phys. Rev. C97, 0552113 (2018)