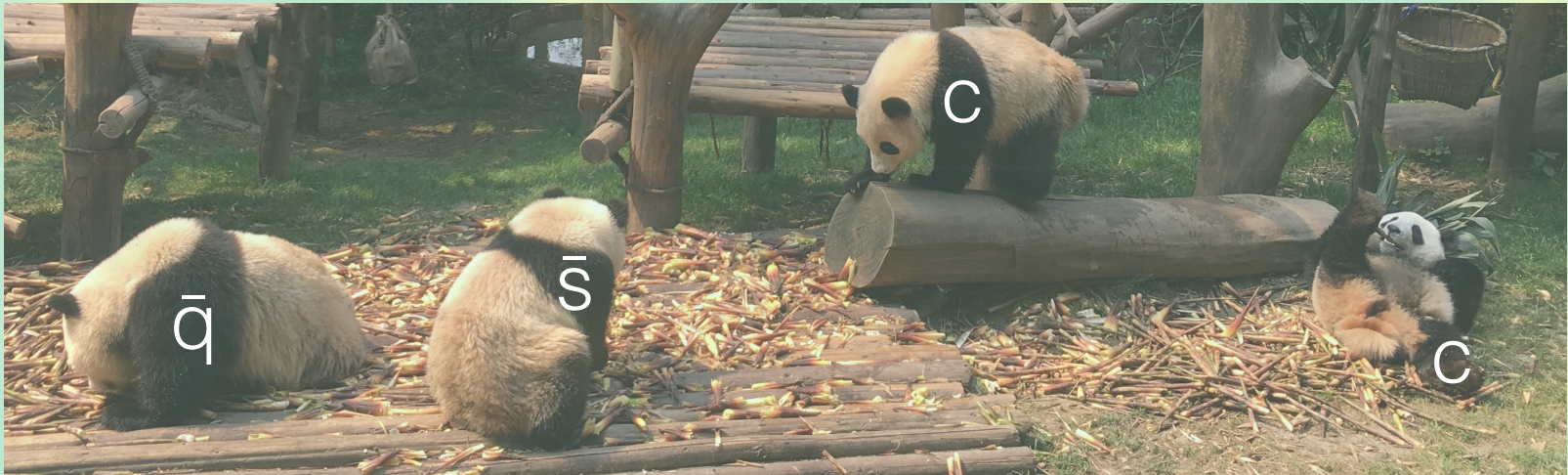


# Exotic hadrons and emergent long range correlation in QCD



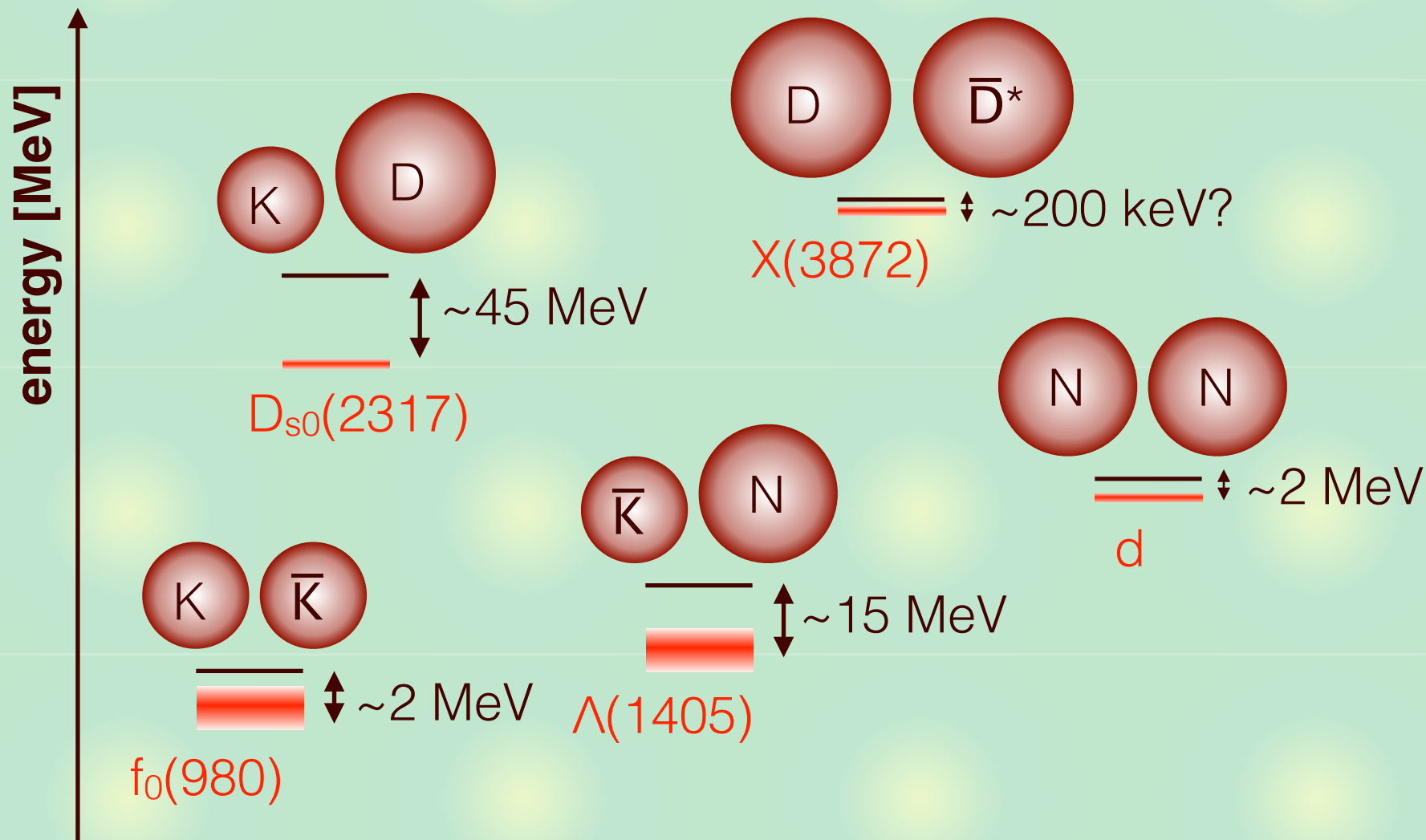
**Tetsuo Hyodo**

*Tokyo Metropolitan Univ.*



# Hadron clusters

## Hadrons near an s-wave two-body threshold



“hadronic molecules” (various flavors, baryon numbers, ...)

# Two-body universal physics

## Near-threshold s-wave state: **universal physics**

E. Braaten, H.-W. Hammer, *Phys. Rept.* **428**, 259 (2006);

P. Naidon, S. Endo, *Rept. Prog. Phys.* **80**, 056001 (2017)

- **scattering length**  $|a| \gg$  **interaction range**  $r_e$
- **size of (quasi-)bound state**  $\sim |a|$  : **loosely bound**
- **relation with eigenenergy**  $E$

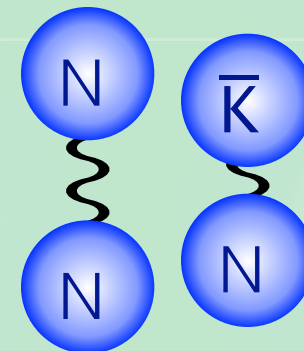
$$a(E) \sim \frac{i}{k} = \frac{i}{\sqrt{2\mu E}}$$

**vdW**

**Examples:** d,  $\Lambda(1405)$ ,  $^4\text{He}$  **dimer**

	NN [fm]	$\bar{K}N$ [fm]	$^4\text{He}$ [ $a_0$ ]
$a(E)$	<b>4.3</b>	<b>1.2-0.8i</b>	<b>178</b>
$a_{\text{emp}}$	<b>5.1</b>	<b>1.4-0.9i</b>	<b>189</b>
$r_e$	<b>1.4</b>	<b>0.4</b>	<b>10</b>

**strong**



$^4\text{He}$



# Classification of hadrons

Observed hadrons

PDG2018 : <http://pdg.lbl.gov/>

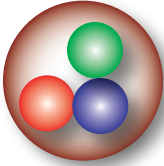
Only **color singlet** states are observed.

—> Color confinement problem

Flavor quantum numbers are described by  $qqq/q\bar{q}$ .

Why no  $qqq\bar{q}$ ,  $qqqq\bar{q}$ , ... states (**exotic hadrons**)?

—> Exotic hadron problem, as nontrivial as confinement!

$\Lambda(2700)$ $13/2^+ **$	$\Lambda(1710)$ $1/2^+ *$	$\Sigma(3000)$ *	 <p>~ 150 baryons</p>
$\Lambda(1800)$ $1/2^- ***$	$\Lambda(1810)$ $1/2^+ ***$	$\Sigma(3170)$ *	
$\Lambda(1820)$ $5/2^+ ****$	$\Lambda(1830)$ $5/2^- ****$		
$\Lambda(1890)$ $3/2^+ ****$	$\Lambda(2000)$ *		
$\Lambda(2020)$ $7/2^+ *$	$\Lambda(2050)$ $3/2^- *$		
$\Lambda(2100)$ $7/2^- ****$	$\Lambda(2110)$ $5/2^+ ***$		
$\Lambda(2325)$ $3/2^- *$	$\Lambda(2350)$ $9/2^+ ***$		
$\Lambda(2585)$ **			
$\Sigma_b^-$ $1/2^+ ***$	$\Sigma_b^0$ $1/2^+ ***$		
$\Sigma_b^+$ $3/2^+ ***$	$\Xi_b^-$ $1/2^+ ***$		
$\Xi_b^0$ $1/2^+ ***$	$\Xi_b^0(5935)^-$ $1/2^+ ***$		
$\Xi_b^+$ $3/2^+ ***$	$\Xi_b^0(5945)^0$ $3/2^+ ***$		
$\Xi_b^+$ $3/2^+ ***$	$\Xi_b^+(5955)$ $3/2^+ ***$		
$\Omega_b^-$ $1/2^+ ***$			

LIGHT UNFLAVORED ( $\bar{u}, \bar{d}, \bar{s}$ )	STRANGE ( $\bar{c}, \bar{b}, \bar{t}$ )	CHARMED, STRANGE ( $\bar{c}, \bar{s}, \bar{b}$ )	$c\bar{c}$ ( $\bar{b}\bar{b}$ )
$a_1(1640)$ $1^-(1^{++})$	$a_0(2450)$ $1^-(6^{++})$	$D_0^*(2400)^0$ $1/2^-(0^+)$	BOTTOM, CHARMED ( $B=C=+1$ )
$f_2(1640)$ $0^+(2^{++})$	$f_0(2510)$ $0^+(6^{++})$	$D_0^*(2400)^+$ $1/2^-(0^+)$	$B_c^+$ $0^-(0^-)$
$\rho(1645)$ $0^+(2^-)$		$D_1(2420)^0$ $1/2^-(1^+)$	$B_c(2S)^+$ $?^?(?^?)$
$\omega(1650)$ $0^-(1^-)$		$D_1(2420)^+$ $1/2^-(1^+)$	
$\omega_3(1670)$ $0^-(3^-)$		$D_1(2430)^0$ $1/2^-(1^+)$	
$\pi_2(1670)$ $1^-(2^-)$		$D_2^*(2460)^0$ $1/2^-(2^+)$	
		$D_2^*(2460)^+$ $1/2^-(2^+)$	
		$D(2550)^0$ $1/2^-(0^-)$	
		$D(2600)$ $1/2^-(?)$	
		$D^*(2640)^+$ $1/2^-(?)$	
		$D(2750)$ $1/2^-(?)$	
			$\chi_{b1}(1P)$ $0^-(1^{--})$
			$\chi_{b2}(1P)$ $0^-(2^{--})$
			$\eta_b(2S)$ $0^+(0^-)$
			$\Upsilon(2S)$ $0^-(1^{--})$
			$\Upsilon(1D)$ $0^-(2^{--})$
			$\chi_{b0}(2P)$ $0^+(0^{++})$
			$\chi_{b1}(2P)$ $0^+(1^{++})$
			$\chi_{b2}(2P)$ $0^+(2^{++})$
			$\Upsilon(3S)$ $0^-(1^{--})$
			$\chi_{b1}(3P)$ $0^+(1^{++})$
			$\Upsilon(4S)$ $0^-(1^{--})$
			$X(10610)^+$ $1^+(1^+)$
			$X(10610)^0$ $1^+(1^+)$
			$X(10650)^+$ $?^+(1^+)$
			$\Upsilon(10860)$ $0^-(1^{--})$
			$\Upsilon(11020)$ $0^-(1^{--})$

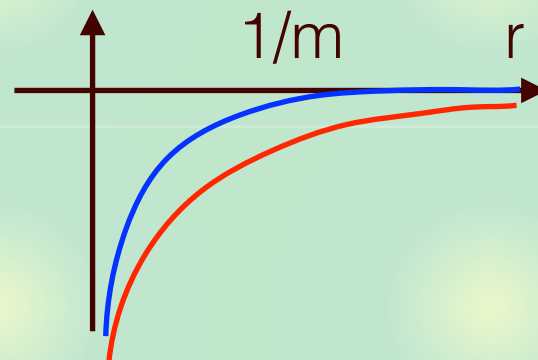
All ~ 360 hadrons emerge from single QCD Lagrangian.

# Long range correlation in QCD?

## Two-body potential

$$V(r) \propto \frac{1}{r} \quad : \text{long (infinite) range}$$

$$V(r) \propto \frac{e^{-mr}}{r} \quad : \text{finite } (\sim 1/m) \text{ range}$$



Hadron-hadron interaction is considered to be **finite range**.

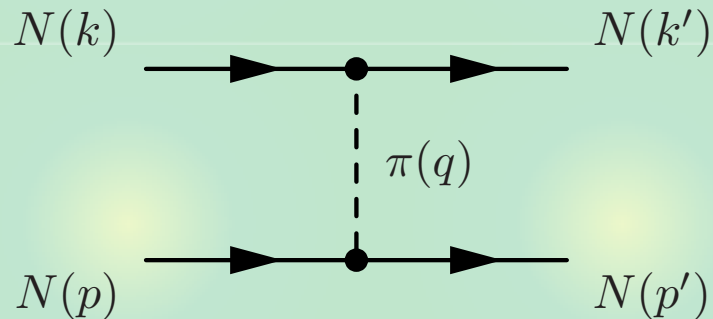
- Longest interaction range  
← exchange of lightest particle ( $\pi$ )  $\sim 1$  fm
- Absence of the long range force is the basis for the (standard) scattering theory, Lüscher/HALQCD method, etc.

There can be (quasi) **long range** force beyond 1 fm.

M. Sanchez Sanchez, L.S. Geng, J. Lu, T. Hyodo, M.P. Valderrama, PRD98, 054001 (2018)

# NN potential

## Low energy NN interaction : $\pi$ exchange



- **Static approx.**  $p^\mu = (M_N, \mathbf{p})$ ,  $p'^\mu = (M_N, \mathbf{p}')$ ,  $q^\mu = p'^\mu - p^\mu = (0, \mathbf{q})$

- **Coupling**  $g\bar{N}i\gamma_5\pi N \sim g\chi^\dagger \boldsymbol{\sigma} \cdot \mathbf{q}\chi$  **(isospin ignored)**

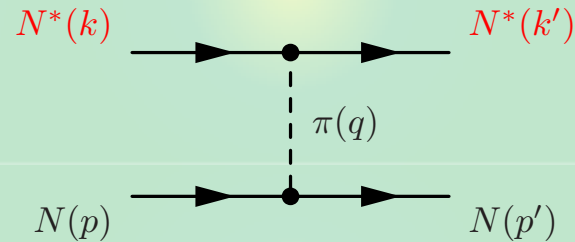
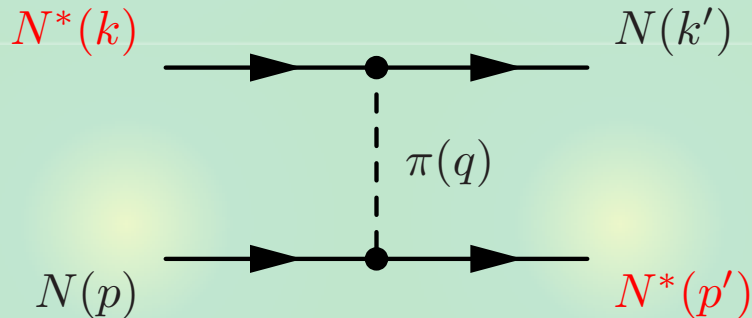
## Potential

$$V(\mathbf{r}) \sim \text{F.T.} \left\{ \underbrace{g^2 (\boldsymbol{\sigma}_1 \cdot \mathbf{q})(\boldsymbol{\sigma}_2 \cdot \mathbf{q})}_{\text{Tensor op.}} \underbrace{\frac{-1}{q^2 + m_\pi^2}}_{\text{Yukawa}} \right\} \frac{1}{(q^0)^2 - \mathbf{q}^2 - m_\pi^2}$$

**Tensor op.** **Yukawa**  $\frac{e^{-m_\pi r}}{r}$

# NN\* potential (exchange)

## NN\*(J<sup>P</sup>=1/2-) interaction



**Mass difference  
= energy transfer**

$$\Delta = M_{N^*} - M_N$$

- **Static approx.**  $p^\mu = (M_N, \mathbf{p})$ ,  $p'^\mu = (M_{N^*}, \mathbf{p}')$ ,  $q^\mu = (\Delta, \mathbf{q})$

- **Coupling**  $\tilde{g} \bar{N}^* \pi N + \text{h.c.} \sim \tilde{g} \chi^\dagger \mathbf{1} \chi$

**Potential (P<sub>σ</sub>: spin exchange factor)**

$$\mu = \sqrt{m_\pi^2 - \Delta^2}$$

$$V(r) \sim \text{F.T.} \left\{ \tilde{g}^2 \frac{1}{\Delta^2 - \mathbf{q}^2 - m_\pi^2} \right\} P_\sigma = \text{F.T.} \left\{ \tilde{g}^2 \frac{-1}{\mathbf{q}^2 + \mu^2} \right\} P_\sigma \sim \tilde{g}^2 P_\sigma \frac{e^{-\mu r}}{r}$$

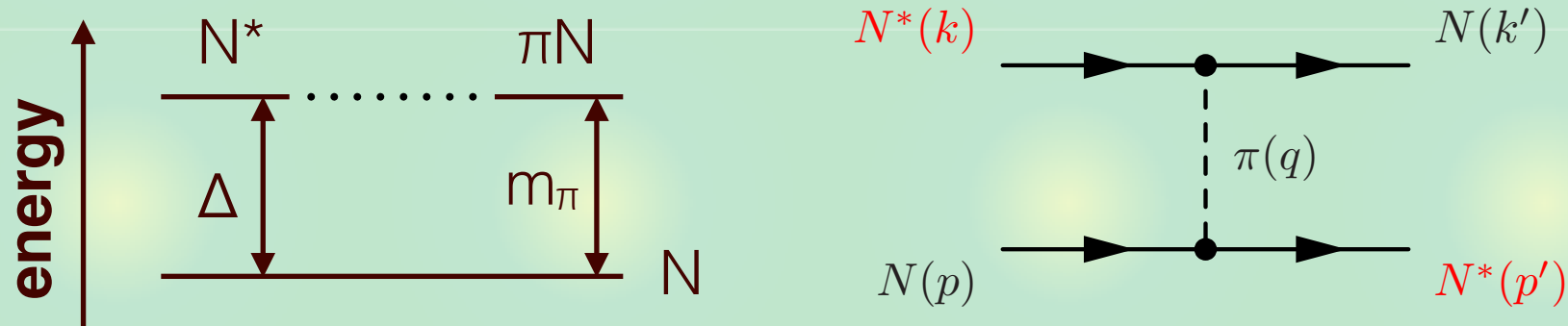
- **Sign of V(r) is fixed and attractive (c.f. σ exchange in NN)**

- **Effective mass μ=0 → long range force (Coulomb like)**

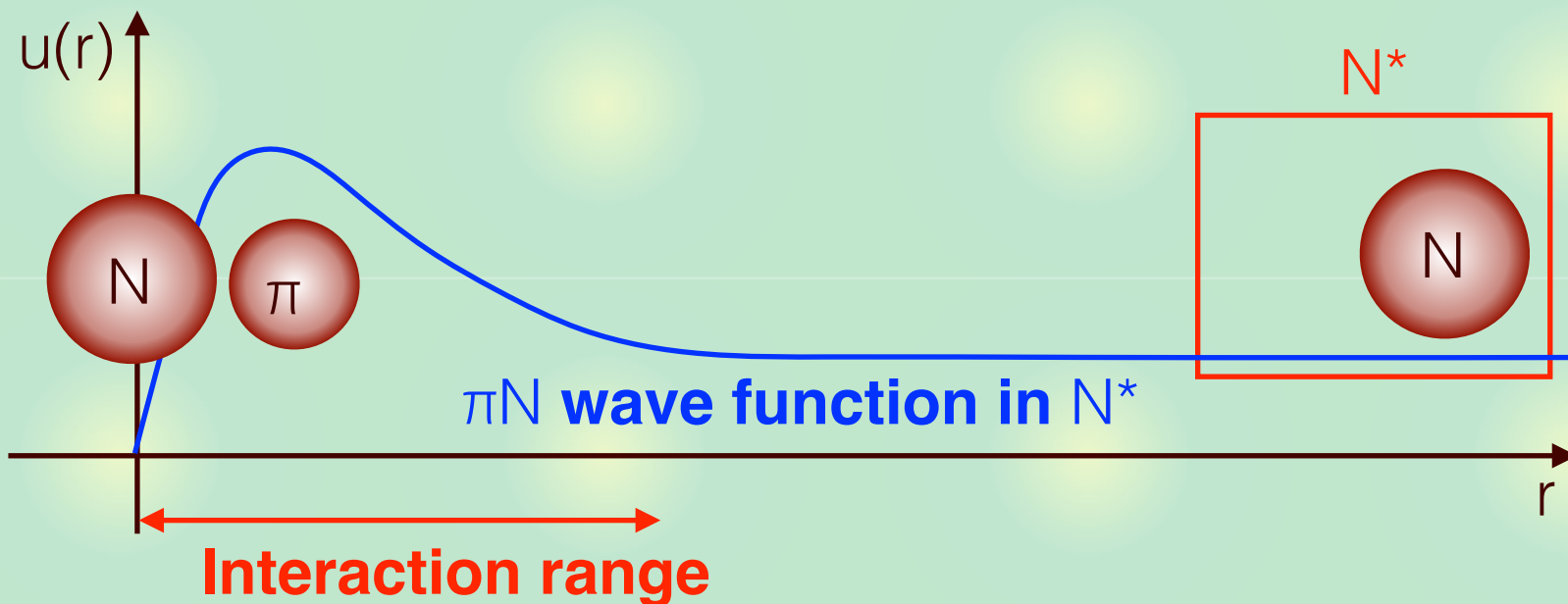


# Unitary limit and zero-energy resonance

What does  $\mu = (m_\pi^2 - \Delta^2)^{1/2} = 0 \iff \Delta = m_\pi$  mean?

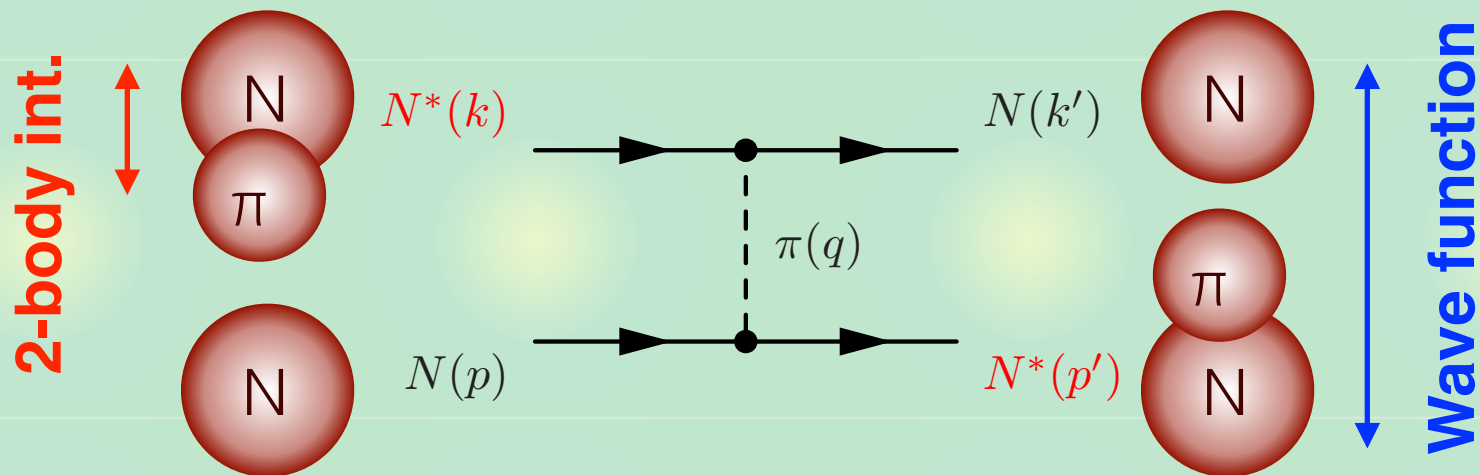


-  $\Delta = m_\pi$  :  $N^*$  lies on top of the  $\pi N$  threshold  $\rightarrow a_{\pi N} = \infty$



# Remarks and toward physical realization

$N^*N \sim \pi NN$  : effective description of three-body system



## Similarity with the Efimov effect

- spatially large three-body system via unitary two-body int.
- $1/r$  attraction (not  $1/r^2$ )?

## Realization in physical hadron systems

- No system with exact  $\mu=0$  ( $N^*$ :  $\Delta \sim 595$  MeV /  $m_\pi \sim 140$  MeV)
- Is there any system with **small**  $\mu$ ?

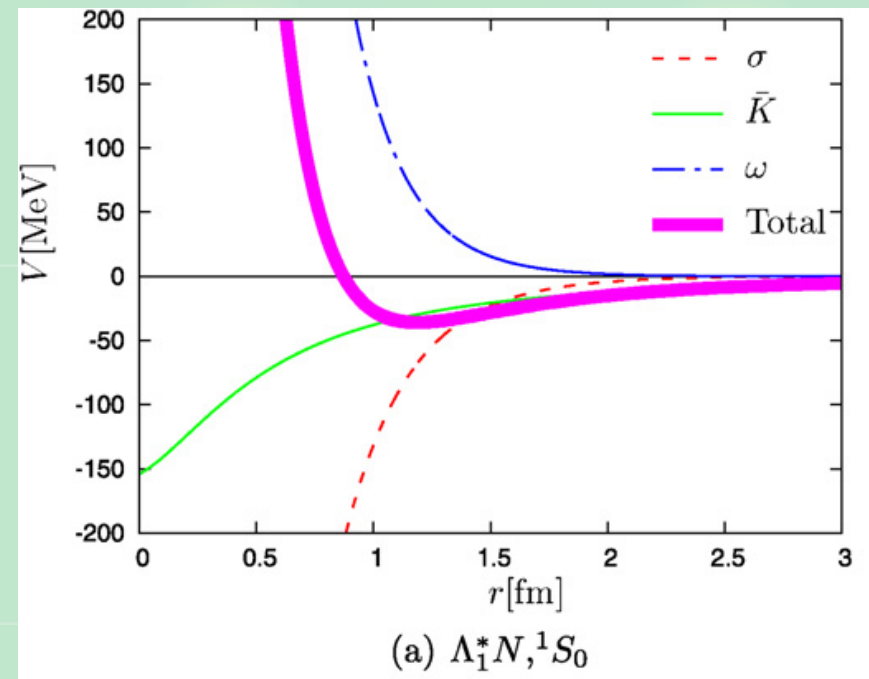
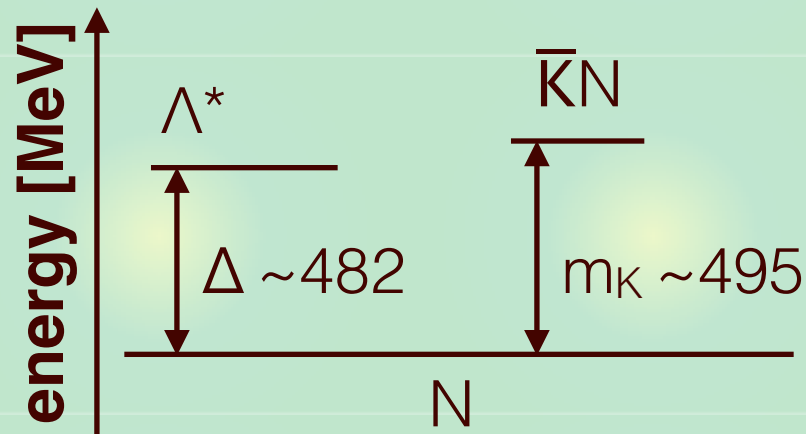
# Strange dibaryon

$\Lambda(1405)=\Lambda^*$ :  $\bar{K}N$  quasibound state near the threshold

T. Hyodo, D. Jido, Prog. Part. Nucl. Phys. 67, 55 (2012)

- $\bar{K}$  exchange between  $\Lambda^*$  and N

$\Lambda^*$  (at 1420 MeV),  $\bar{K}N$  threshold



- $\mu \sim 91$  MeV:  $\bar{K}$  exchange has longer tail than expected
- attractive in spin singlet channel  $\rightarrow \bar{K}NN$  as  $\Lambda^*N$  system

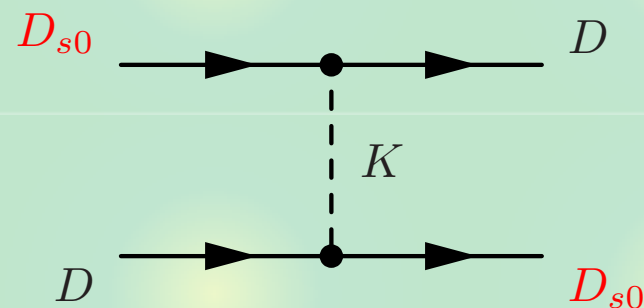
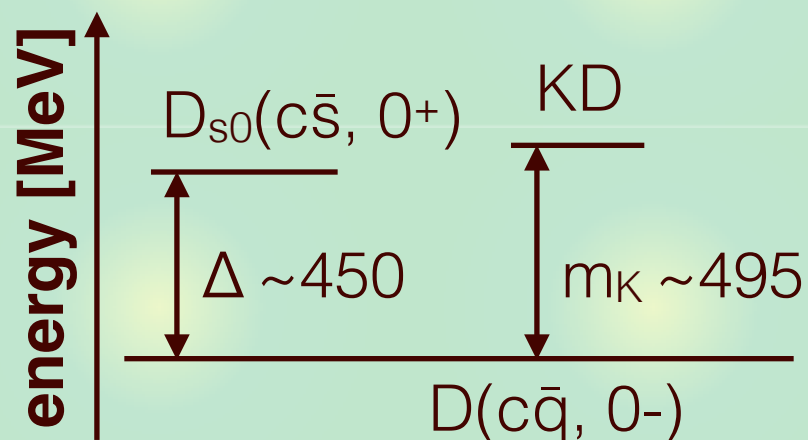
T. Uchino, T. Hyodo, M. Oka, Nucl. Phys. A, 868-869, 53 (2011)

# Doubly charmed exotic meson

We consider  $D_{s0}(c\bar{s}, 0^+)D(c\bar{q}, 0^-)$  system via K exchange

- Charm  $C=2$ : manifestly **exotic** ( $cc\bar{q}\bar{s}$ )

$D_{s0}(2317)$ , KD threshold



- K exchange gives **quasi-long range** ( $\mu \sim 200$  MeV) attraction

Can the attraction generate a bound state?

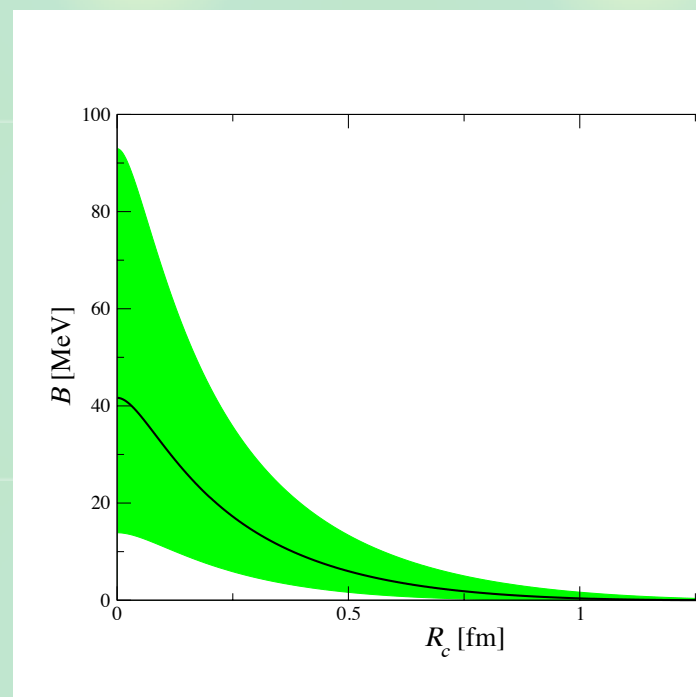


# Prediction of binding energy

## Effective Lagrangian for $D_{s0}DK$ (and HQ partners) coupling

$$\mathcal{L} = \frac{h}{2} \text{Tr}[\bar{H}_a S_b A_{ab} \gamma_5] + \text{C.C.}$$

- coupling constant  $h$  :  $D_0 \rightarrow D\pi$  decay + SU(3) symmetry
- Short range cutoff  $R_c \leftarrow$  hadron size

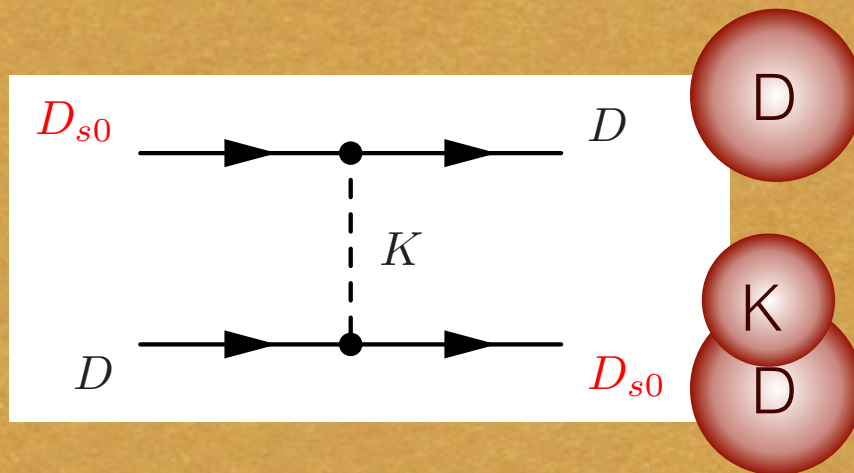


- $R_c \sim 0.5$  fm  $\rightarrow$   $\sim 6$  MeV binding

# Summary



**Long range correlation** among hadrons emerges when the mass difference  $\Delta$  matches with the mass of the exchange particle  $m$ .



$$V(r) \sim \frac{e^{-\mu r}}{r}, \quad \mu = \sqrt{m^2 - \Delta^2}$$



**K exchange in  $D_{s0}(0^+)D(0^-)$  system:  $\mu \sim 200$  MeV**  
**—> prediction of exotic charmed tetraquark**

M. Sanchez Sanchez, L.S. Geng, J. Lu, T. Hyodo, M.P. Valderrama,  
Phys. Rev. D98, 054001 (2018)