

# Compositeness of hadrons from effective field theory



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2018, Dec. 12th <sub>1</sub>

# Contents

## Introduction: exotic hadron resonances

## Compositeness of hadron resonances

S. Weinberg, Phys. Rev. 137, B672 (1965);

T. Hyodo, Int. J. Mod. Phys. A 28, 1330045 (2013)

### - Weak binding relation from EFT

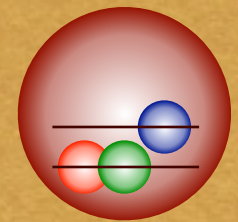
Y. Kamiya, T. Hyodo, Phys. Rev. C93, 035203 (2016);

Y. Kamiya, T. Hyodo, PTEP2017, 023D02 (2017)

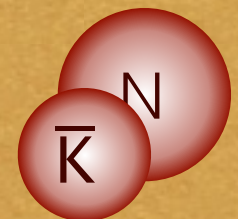
### - Analysis for $\Lambda(1405)$

Y. Ikeda, T. Hyodo, W. Weise, Phys. Lett. B 706, 63 (2011);

Y. Ikeda, T. Hyodo, W. Weise, Nucl. Phys. A 881 98 (2012)



or



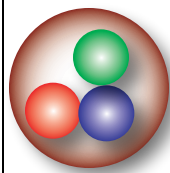
## Summary

# Classification of hadrons

## Observed hadrons

PDG2018 : <http://pdg.lbl.gov/>

$p$	$1/2^+$ ****	$\Delta(1232)$	$3/2^+$ ****	$\Sigma^+$	$1/2^+$ ****	$\Xi^0$	$1/2^+$ ****	$\Lambda_c^+$	$1/2^+$ ****
$n$	$1/2^+$ ****	$\Delta(1600)$	$3/2^+$ ***	$\Sigma^0$	$1/2^+$ ****	$\Xi^-$	$1/2^+$ ****	$\Lambda_c(2595)^+$	$1/2^-$ ***
$N(1440)$	$1/2^+$ ****	$\Delta(1620)$	$1/2^-$ ****	$\Sigma^-$	$1/2^+$ ****	$\Xi(1530)$	$3/2^+$ ****	$\Lambda_c(2625)^+$	$3/2^-$ ***
$N(1520)$	$3/2^-$ ****	$\Delta(1700)$	$3/2^-$ ****	$\Sigma(1385)$	$3/2^+$ ****	$\Xi(1620)$	*	$\Lambda_c(2765)^+$	*
$N(1535)$	$1/2^-$ ****	$\Delta(1750)$	$1/2^+$ *	$\Sigma(1480)$	*	$\Xi(1690)$	***	$\Lambda_c(2880)^+$	$5/2^+$ ***
$N(1650)$	$1/2^-$ ****	$\Delta(1900)$	$1/2^-$ **	$\Sigma(1560)$	**	$\Xi(1820)$	$3/2^-$ ***	$\Lambda_c(2940)^+$	***
$N(1675)$	$5/2^-$ ****	$\Delta(1905)$	$5/2^+$ ****	$\Sigma(1580)$	$3/2^-$ **	$\Xi(1950)$	***	$\Sigma_c(2455)$	$1/2^+$ ****
$N(1680)$	$5/2^+$ ****	$\Delta(1910)$	$1/2^+$ ****	$\Sigma(1620)$	$1/2^-$ *	$\Xi(2030)$	$\geq 5/2^?$ ***	$\Sigma_c(2520)$	$3/2^+$ ****
$N(1685)$	*	$\Delta(1920)$	$3/2^+$ ***	$\Sigma(1660)$	$1/2^+$ ***	$\Xi(2120)$	**	$\Sigma_c(2800)$	***
$N(1700)$	$3/2^-$ ***	$\Delta(1930)$	$5/2^-$ ***	$\Sigma(1670)$	$3/2^-$ ****	$\Xi(2250)$	*	$\Xi_c$	$1/2^+$ ***
$N(1710)$	$1/2^+$ ***	$\Delta(1940)$	$3/2^-$ **	$\Sigma(1690)$	**	$\Xi(2370)$	**	$\Xi_c$	$1/2^+$ ***
$N(1720)$	$3/2^+$ ****	$\Delta(1950)$	$7/2^+$ ****	$\Sigma(1730)$	$3/2^+$ *	$\Xi(2500)$	*	$\Xi_c$	$1/2^+$ ***
$N(1860)$	$5/2^+$ **	$\Delta(2000)$	$5/2^+$ **	$\Sigma(1750)$	$1/2^-$ ***			$\Xi_c$	$1/2^+$ ***
$N(1875)$	$3/2^-$ ***	$\Delta(2150)$	$1/2^-$ *	$\Sigma(1770)$	$1/2^+$ *	$\Omega^-$	$3/2^+$ ****	$\Xi_c$	$1/2^+$ ***
$N(1880)$	$1/2^+$ **	$\Delta(2200)$	$7/2^-$ *	$\Sigma(1775)$	$5/2^-$ ****	$\Omega(2250)^-$	***	$\Xi_c(2645)$	$3/2^+$ ***
$N(1895)$	$1/2^-$ **	$\Delta(2300)$	$9/2^+$ **	$\Sigma(1840)$	$3/2^+$ **	$\Omega(2380)^-$	**	$\Xi_c(2790)$	$1/2^-$ ***
$N(1900)$	$3/2^+$ ***	$\Delta(2350)$	$5/2^-$ **	$\Sigma(1880)$	$1/2^+$ **	$\Omega(2470)^-$	**	$\Xi_c(2815)$	$3/2^-$ ***
$N(1990)$	$7/2^+$ **	$\Delta(2390)$	$7/2^+$ *	$\Sigma(1900)$	$1/2^-$ *			$\Xi_c(2930)$	*
$N(2000)$	$5/2^+$ **	$\Delta(2400)$	$9/2^-$ **	$\Sigma(1915)$	$5/2^+$ ****			$\Xi_c(2980)$	**
$N(2040)$	$3/2^+$ **	$\Delta(2420)$	$11/2^+$ ****	$\Sigma(1940)$	$3/2^+$ **			$\Xi_c(3055)$	***
$N(2060)$	$5/2^-$ **	$\Delta(2750)$	$13/2^-$ **	$\Sigma(1940)$	$3/2^-$ ***			$\Xi_c(3080)$	***
$N(2100)$	$1/2^+$ *	$\Delta(2950)$	$15/2^+$ **	$\Sigma(2000)$	$1/2^-$ *			$\Xi_c(3123)$	*
$N(2120)$	$3/2^-$ **			$\Sigma(2030)$	$7/2^+$ ****			$\Omega_c^0$	$1/2^+$ ***
$N(2190)$	$7/2^-$ ****	$\Lambda$	$1/2^+$ ****	$\Sigma(2070)$	$5/2^+$ **			$\Omega_c(2770)^0$	$3/2^+$ ***
$N(2220)$	$9/2^+$ ****	$\Lambda(1405)$	$1/2^-$ ****	$\Sigma(2080)$	$3/2^+$ **			$\Xi_{cc}^+$	*
$N(2250)$	$9/2^-$ ****	$\Lambda(1520)$	$3/2^-$ ****	$\Sigma(2100)$	$7/2^-$ *				
$N(2300)$	$1/2^+$ **	$\Lambda(1600)$	$1/2^+$ ***	$\Sigma(2250)$	***			$\Lambda_b^0$	$1/2^+$ ***
$N(2570)$	$5/2^-$ **	$\Lambda(1670)$	$1/2^-$ ****	$\Sigma(2455)$	**			$\Lambda_b(5912)^0$	$1/2^-$ ***
$N(2600)$	$11/2^-$ ***	$\Lambda(1690)$	$3/2^-$ ****	$\Sigma(2620)$	**			$\Lambda_b(5920)^0$	$3/2^-$ ***
$N(2700)$	$13/2^+$ **	$\Lambda(1710)$	$1/2^+$ *	$\Sigma(3000)$	*			$\Sigma_b$	$1/2^+$ ***
		$\Lambda(1800)$	$1/2^-$ ***	$\Sigma(3170)$	*			$\Sigma_b^+$	$3/2^+$ ***
		$\Lambda(1810)$	$1/2^+$ ***					$\Xi_b^-$	$1/2^+$ ***
		$\Lambda(1820)$	$5/2^+$ ****					$\Xi_b(5935)^-$	$1/2^+$ ***
		$\Lambda(1830)$	$5/2^-$ ****					$\Xi_b(5945)^0$	$3/2^+$ ***
		$\Lambda(1890)$	$3/2^+$ ****					$\Xi_b(5955)^-$	$3/2^+$ ***
		$\Lambda(2000)$	*					$\Omega_b$	$1/2^+$ ***
		$\Lambda(2020)$	$7/2^+$ *						
		$\Lambda(2050)$	$3/2^-$ *						
		$\Lambda(2100)$	$7/2^-$ ****						
		$\Lambda(2110)$	$5/2^+$ ***						
		$\Lambda(2325)$	$3/2^-$ *						
		$\Lambda(2350)$	$9/2^+$ ***						
		$\Lambda(2585)$	**						



~ 150 baryons

LIGHT UNFLAVORED (S = C = B = 0)		STRANGE (S = ±1, C = B = 0)		CHARMED, STRANGE (C = S = ±1)		$c\bar{c}$ $F_1(F_2)$			
$F_1(F_2)$	$F_1(F_2)$	$F_1(F_2)$	$F_1(F_2)$	$F_1(F_2)$	$F_1(F_2)$	$F_1(F_2)$	$F_1(F_2)$		
$\pi^+$	$1^-(0^-)$	$\rho(1680)$	$0^-(1^-)$	$K^+$	$1/2(0^-)$	$D_s^+$	$0^-(0^-)$	$\eta_c(1S)$	$0^+(0^-)$
$\pi^0$	$1^-(0^-)$	$\rho_3(1690)$	$1^+(3^-)$	$K^0$	$1/2(0^-)$	$D_s^0$	$0^-(0^-)$	$J/\psi(1S)$	$0^-(1^-)$
$\eta$	$0^+(0^-)$	$\rho(1700)$	$1^+(1^-)$	$K_S^0$	$1/2(0^-)$	$D_{s1}^0(2317)^0$	$0^+(0^-)$	$\chi_{c0}(1P)$	$0^+(0^+)$
$\eta(500)$	$0^+(0^+)$	$a_2(1700)$	$1^-(2^+)$	$K_L^0$	$1/2(0^-)$	$D_{s1}^+(2460)^+$	$0^+(1^+)$	$\chi_{c1}(1P)$	$0^+(1^+)$
$\rho(770)$	$1^+(1^+)$	$\omega(1710)$	$0^+(0^+)$	$K_S^0(800)$	$1/2(0^+)$	$D_{s1}^+(2536)^+$	$0^+(1^+)$	$\chi_{c2}(1P)$	$0^+(2^+)$
$\omega(782)$	$0^-(1^-)$	$\eta(1760)$	$0^+(0^-)$	$K^*(892)$	$1/2(1^-)$	$D_{s1}^+(2573)^+$	$0^+(1^+)$	$\eta_c(2S)$	$0^+(0^+)$
$\eta(958)$	$0^+(0^+)$	$\eta(1800)$	$1^-(0^-)$	$K_1^*(1270)$	$1/2(1^+)$	$D_{s1}^+(2700)^+$	$0^-(1^-)$	$\psi(2S)$	$0^-(1^-)$
$\eta(980)$	$0^+(0^+)$	$\eta(1810)$	$0^+(2^+)$	$K_1^*(1400)$	$1/2(1^+)$	$D_{s1}^+(2860)^+$	$0^-(0^-)$	$\psi(3770)$	$0^-(1^-)$
$\omega(980)$	$1^-(0^+)$	$X(1835)$	$2^?(2^-)$	$K^*(1410)$	$1/2(1^-)$	$D_{s1}^+(3040)^+$	$0^-(0^-)$	$X(3823)$	$2^?(2^-)$
$\phi(1020)$	$0^-(1^-)$	$X(1840)$	$2^?(2^?)$	$K_0^*(1430)$	$1/2(0^+)$			$X(3872)$	$0^+(1^+)$
$\phi(1170)$	$0^-(1^+)$	$\omega_3(1850)$	$0^-(3^-)$	$K_2^*(1430)$	$1/2(2^+)$	BOTTOM (B = ±1)		$X(3900)^0$	$2^?(1^+)$
$h_1(1235)$	$1^+(1^+)$	$\eta_2(1870)$	$0^+(2^+)$	$K(1460)$	$1/2(0^-)$	$B^+$	$1/2(0^-)$	$X(3900)^0$	$2^?(1^+)$
$a_1(1260)$	$1^+(1^+)$	$\eta_3(1880)$	$1^-(2^-)$	$K_0^*(1580)$	$1/2(2^-)$	$B^0$	$1/2(0^-)$	$\chi_{c0}(3915)$	$0^+(0^+)$
$f_2(1270)$	$0^+(2^+)$	$\rho(1900)$	$1^+(1^-)$	$K(1630)$	$1/2(2^-)$	$B^0/B^0/B^0$ ADMIXTURE	$1/2(0^-)$	$\chi_{c2}(3940)$	$0^+(0^+)$
$f_1(1285)$	$0^+(1^+)$	$\rho(1910)$	$1^+(2^+)$	$K(1630)$	$1/2(2^-)$	$B^+$ ADMIXTURE	$1/2(0^-)$	$X(4020)$	$2^?(2^?)$
$\eta(1295)$	$0^+(0^+)$	$\rho_3(1900)$	$0^+(2^+)$	$K_1(1650)$	$1/2(1^+)$	$B^+$ ADMIXTURE	$1/2(0^-)$	$\psi(4040)$	$0^-(1^-)$
$\pi(1300)$	$1^-(0^+)$	$\rho_3(1950)$	$0^+(2^+)$	$K_1^*(1680)$	$1/2(1^-)$	$V_{cb}$ and $V_{cb}$ CKM Matrix Elements	$1/2(0^-)$	$X(4050)^0$	$2^?(2^?)$
$a_2(1320)$	$1^-(2^+)$	$\rho_3(1990)$	$1^+(3^-)$	$K_2^*(1770)$	$1/2(2^-)$	$B_c^+$	$1/2(0^-)$	$X(4140)$	$0^+(2^+)$
$f_0(1370)$	$0^+(0^+)$	$f_0(2010)$	$0^+(2^+)$	$K_3^*(1780)$	$1/2(3^-)$	$B_c^0$	$1/2(0^-)$	$\psi(4160)$	$0^-(1^-)$
$h_1(1380)$	$1^-(1^+)$	$f_0(2020)$	$0^+(2^+)$	$K_0^*(1820)$	$1/2(2^-)$	$B_c^0$	$1/2(0^-)$	$X(4160)$	$0^-(1^-)$
$\eta_1(1400)$	$1^-(1^+)$	$a_0(2040)$	$1^-(4^+)$	$K(1830)$	$1/2(0^-)$	$B_1(5721)^+$	$1/2(1^+)$	$X(4160)$	$0^-(1^-)$
$\eta_3(1400)$	$1^-(1^+)$	$f_0(2050)$	$0^+(4^+)$	$K_0^*(1950)$	$1/2(0^+)$	$B_1(5721)^0$	$1/2(1^+)$	$X(4230)$	$2^?(1^-)$
$\eta(1405)$	$0^+(0^+)$	$f_0(2100)$	$1^-(2^+)$	$K_2^*(1980)$	$1/2(2^+)$	$B_1^*(5732)$	$2^?(2^?)$	$X(4240)^0$	$2^?(1^-)$
$f_1(1420)$	$0^+(1^+)$	$f_0(2100)$	$0^+(0^+)$	$K_2^*(1980)$	$1/2(2^+)$	$B_2^*(5747)^0$	$1/2(2^+)$	$X(4250)^0$	$2^?(1^-)$
$\omega(1420)$	$0^-(1^-)$	$f_2(2150)$	$0^+(2^+)$	$K_4^*(2045)$	$1/2(4^+)$	$B_2^*(5747)^0$	$1/2(2^+)$	$X(4250)^0$	$2^?(1^-)$
$f_2(1430)$	$0^+(2^+)$	$\rho(2150)$	$1^-(1^-)$	$K_0^*(2050)$	$1/2(2^+)$	$B(5970)^0$	$2^?(2^?)$	$X(4260)$	$2^?(1^-)$
$a_0(1450)$	$1^-(0^+)$	$\phi(2170)$	$0^-(1^-)$	$K_0^*(2380)$	$1/2(3^-)$	$B(5970)^0$	$2^?(2^?)$	$X(4360)$	$2^?(1^-)$
$\rho(1450)$	$1^+(1^-)$	$f_0(2200)$	$0^+(0^+)$	$K_2^*(2380)$	$1/2(3^-)$	$B(5970)^0$	$2^?(2^?)$	$X(4415)$	$0^-(1^-)$
$\eta(1475)$	$0^+(0^+)$	$f_1(2220)$	$0^+(2^+)$	$K_3^*(2380)$	$1/2(3^-)$	$B(5970)^0$	$2^?(2^?)$	$X(4430)^0$	$2^?(1^-)$
$f_0(1500)$	$0^+(0^+)$	$\eta(2225)$	$0^+(0^-)$	$K_4^*(2500)$	$1/2(4^+)$	BOTTOM, STRANGE (B = ±1, S = ±1)		$X(4660)$	$2^?(1^-)$
$f_1(1510)$	$0^+(1^+)$	$\rho_3(2250)$	$1^+(3^-)$	$K(3100)$	$2^?(2^?)$	CHARMED (C = ±1)		$b\bar{b}$	
$f_2(1525)$	$0^+(2^+)$	$f_0(2300)$	$0^+(2^+)$	$D^+$	$1/2(0^-)$	$B_c^+$	$0^-(0^-)$	$\eta_b(1S)$	$0^+(0^+)$
$f_3(1565)$	$0^+(2^+)$	$f_0(2300)$	$0^+(4^+)$	$D^0$	$1/2(0^-)$	$B_c^0$	$0^-(0^-)$	$\Upsilon(1S)$	$0^-(1^-)$
$\rho(1570)$	$1^+(1^+)$	$f_2(2330)$	$0^+(0^+)$	$D^*$	$1/2(1^-)$	$B_{c1}(5830)^0$	$0^+(1^+)$	$\chi_{b0}(1P)$	$0^+(0^+)$
$h_1(1595)$	$0^-(1^+)$	$f_2(2340)$	$0^+(2^+)$	$D^*$	$1/2(1^-)$	$B_{c1}(5840)^0$	$0^+(2^+)$	$\chi_{b1}(1P)$	$0^+(1^+)$
$\eta_1(1600)$	$1^-(1^+)$	$\rho_3(2350)$	$1^+(5^-)$	$D^*$	$1/2(1^-)$	$B_{c1}(5850)$	$2^?(2^?)$	$\chi_{b2}(1P)$	$0^+(2^+)$
$a_1(1640)$	$1^-(1^+)$	$a_2(2450)$	$1^-(6^+)$	$D_0(2400)^0$	$1/2(0^+)$	BOTTOM, CHARMED (B = C = ±1)		$h_b(1P)$	$2^?(1^+)$
$a_2(1640)$	$0^+(2^+)$	$f_0(2510)$	$0^+(6^+)$	$D_0(2400)^0$	$1/2(0^+)$	$B_c^+$	$0^-(0^-)$	$\eta_b(2S)$	$0^+(0^+)$
$\eta_2(1645)$	$0^+(2^+)$			$D_1(2420)^0$	$1/2(1^+)$	$B_c^0$	$0^-(0^-)$	$\Upsilon(2S)$	$0^-(1^-)$
$\omega(1650)$	$0^-(1^-)$			$D_1(2420)^0$	$1/2(1^+)$			$\Upsilon(3S)$	$0^-(1^-)$
$\omega_3(1670)$	$0^-(3^-)$			$D_1(2430)^0$	$1/2(1^+)$			$\chi_{b0}(2P)$	$0^+(0^+)$
$\pi_2(1670)$	$1^-(2^-)$			$D_2(2460)^0$	$1/2(2^+)$			$\chi_{b1}(2P)$	$0^+(1^+)$
				$D_2(2460)^0$	$1/2(2^+)$			$h_b(2P)$	$2^?(1^+)$
				$D(2550)^0$ </					



# Exotic candidates beyond $qqq/q\bar{q}$

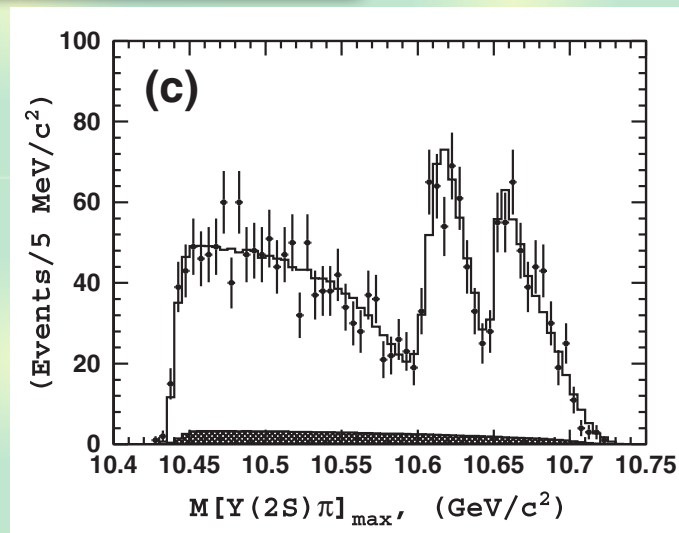
## Tetraquark candidate (Belle)

:  $Z_b(10610)$ ,  $Z_b(10650)$

$$Y(5S) \longrightarrow \pi^\pm + Z_b$$

$$\hookrightarrow Y(nS)(b\bar{b}) + \pi^\mp(u\bar{d}/d\bar{u})$$

A. Bondar, *et al.*, Phys. Rev. Lett. 108, 122001 (2012)



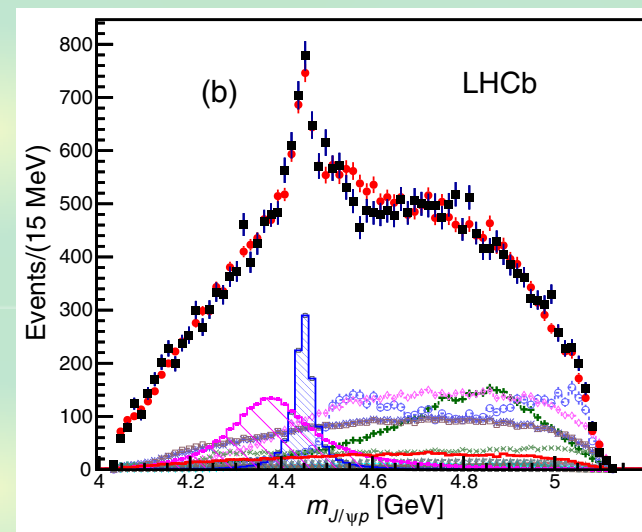
## Pentaquark candidate (LHCb)

:  $P_c(4450)$ ,  $P_c(4380)$

$$\Lambda_b \longrightarrow K^- + P_c$$

$$\hookrightarrow J/\psi(c\bar{c}) + p(uud)$$

R. Aaij, *et al.*, Phys. Rev. Lett. 115, 072001 (2015)



Only a few are observed. **Why only a few?**

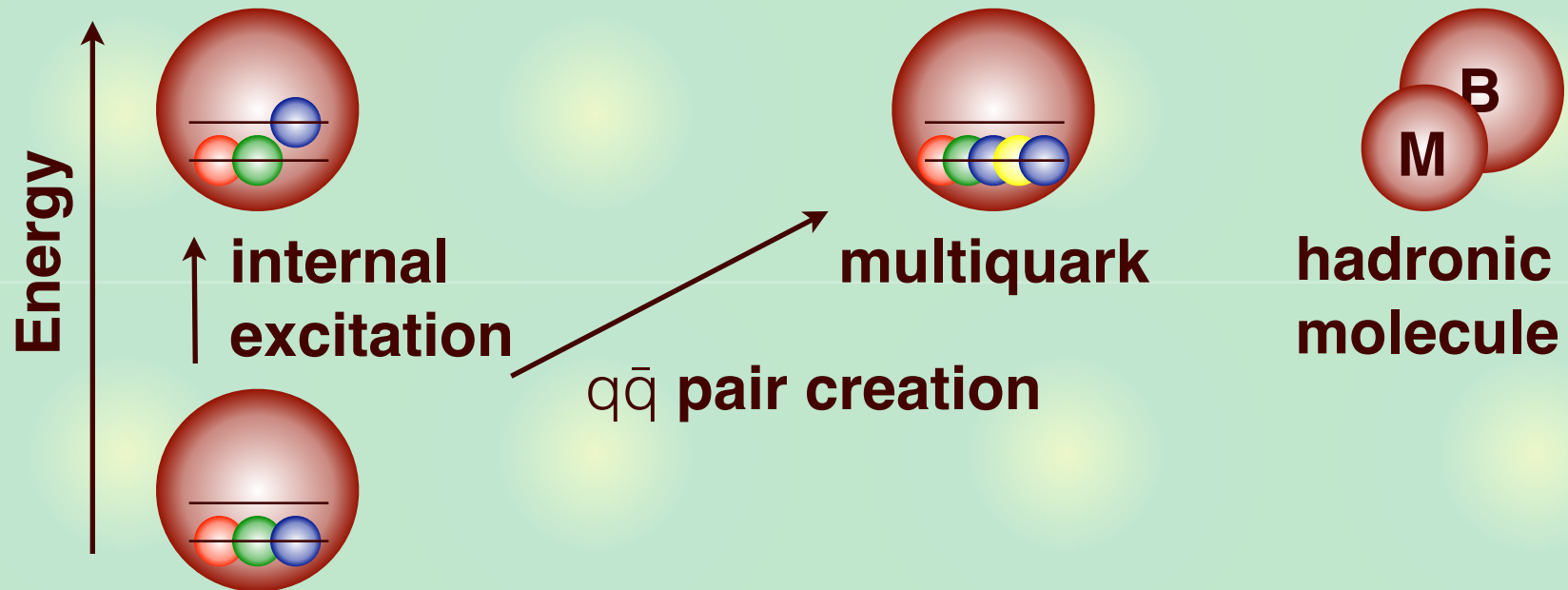


# Various hadronic excitations

## Description of excited baryons

### Conventional structure

### Exotic structures



In QCD, non- $qqq$  structures naturally arise.

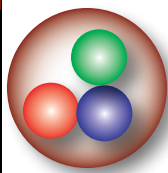
- Baryons: superposition of  $qqq$  + exotic structures
- > How can we disentangle different components?

# Unstable states via strong interaction

## Hadron resonances

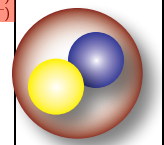
PDG2018 : <http://pdg.lbl.gov/>

$p$	$1/2^+$ ****	$\Delta(1232)$	$3/2^+$ ****	$\Sigma^+$	$1/2^+$ ****	$\Xi^0$	$1/2^+$ ****	$\Lambda_c^+$	$1/2^+$ ****
$n$	$1/2^+$ ****	$\Delta(1600)$	$3/2^+$ ***	$\Sigma^0$	$1/2^+$ ****	$\Xi^-$	$1/2^+$ ****	$\Lambda_c(2595)^+$	$1/2^-$ ***
$N(1440)$	$1/2^+$ ****	$\Delta(1620)$	$1/2^-$ ****	$\Sigma^-$	$1/2^+$ ****	$\Xi(1530)$	$3/2^+$ ****	$\Lambda_c(2625)^+$	$3/2^-$ ***
$N(1520)$	$3/2^-$ ****	$\Delta(1700)$	$3/2^-$ ****	$\Sigma(1385)$	$3/2^+$ ****	$\Xi(1620)$	***	$\Lambda_c(2765)^+$	*
$N(1535)$	$1/2^-$ ****	$\Delta(1750)$	$1/2^+$ *	$\Sigma(1480)$	*	$\Xi(1690)$	***	$\Lambda_c(2880)^+$	$5/2^+$ ***
$N(1650)$	$1/2^-$ ****	$\Delta(1900)$	$1/2^-$ **	$\Sigma(1560)$	**	$\Xi(1820)$	$3/2^-$ ***	$\Lambda_c(2940)^+$	***
$N(1675)$	$5/2^-$ ****	$\Delta(1905)$	$5/2^+$ ****	$\Sigma(1580)$	$3/2^-$ **	$\Xi(1950)$	***	$\Sigma_c(2455)$	$1/2^+$ ****
$N(1680)$	$5/2^+$ ****	$\Delta(1910)$	$1/2^+$ ****	$\Sigma(1620)$	$1/2^-$ *	$\Xi(2030)$	$\geq 5/2^?$ ***	$\Sigma_c(2520)$	$3/2^+$ ****
$N(1685)$	*	$\Delta(1920)$	$3/2^+$ ***	$\Sigma(1660)$	$1/2^+$ ***	$\Xi(2120)$	*	$\Sigma_c(2800)$	***
$N(1700)$	$3/2^-$ ***	$\Delta(1930)$	$5/2^-$ **	$\Sigma(1670)$	$3/2^-$ ****	$\Xi(2250)$	**	$\Xi_c$	$1/2^+$ ***
$N(1710)$	$1/2^+$ ***	$\Delta(1940)$	$3/2^-$ **	$\Sigma(1690)$	**	$\Xi(2370)$	**	$\Xi_c$	$1/2^+$ ***
$N(1720)$	$3/2^+$ ****	$\Delta(1950)$	$7/2^+$ ****	$\Sigma(1730)$	$3/2^+$ **	$\Xi(2500)$	*	$\Xi_c$	$1/2^+$ ***
$N(1860)$	$5/2^+$ **	$\Delta(2000)$	$5/2^+$ **	$\Sigma(1750)$	$1/2^-$ ***			$\Xi_c$	$1/2^+$ ***
$N(1875)$	$3/2^-$ ***	$\Delta(2150)$	$1/2^-$ *	$\Sigma(1770)$	$1/2^+$ **	$\Omega^-$	$3/2^+$ ****	$\Xi_c$	$1/2^+$ ***
$N(1880)$	$1/2^+$ **	$\Delta(2200)$	$7/2^-$ *	$\Sigma(1775)$	$5/2^-$ ****	$\Omega(2250)$	***	$\Xi_c$	$3/2^+$ ****
$N(1895)$	$1/2^-$ **	$\Delta(2300)$	$9/2^+$ **	$\Sigma(1840)$	$3/2^+$ **	$\Omega(2380)^-$	**	$\Xi_c$	$1/2^-$ ****
$N(1900)$	$3/2^+$ ***	$\Delta(2350)$	$5/2^-$ **	$\Sigma(1880)$	$1/2^+$ **	$\Omega(2470)^-$	**	$\Xi_c$	$3/2^-$ ****
$N(1990)$	$7/2^+$ **	$\Delta(2390)$	$7/2^+$ *	$\Sigma(1900)$	$1/2^-$ **			$\Xi_c$	***
$N(2000)$	$5/2^+$ **	$\Delta(2400)$	$9/2^-$ **	$\Sigma(1915)$	$5/2^+$ ****			$\Xi_c$	***
$N(2040)$	$3/2^+$ *	$\Delta(2420)$	$11/2^+$ ****	$\Sigma(1940)$	$3/2^+$ **			$\Xi_c$	***
$N(2060)$	$5/2^-$ **	$\Delta(2750)$	$13/2^-$ **	$\Sigma(1940)$	$3/2^-$ ***			$\Xi_c$	***
$N(2100)$	$1/2^+$ *	$\Delta(2950)$	$15/2^+$ **	$\Sigma(2000)$	$1/2^-$ *			$\Xi_c$	***
$N(2120)$	$3/2^-$ **			$\Sigma(2030)$	$7/2^+$ ****			$\Omega_c^0$	$1/2^+$ ***
$N(2190)$	$7/2^-$ ****	$\Lambda$	$1/2^+$ ****	$\Sigma(2070)$	$5/2^+$ *			$\Omega_c(2770)^0$	$3/2^+$ ***
$N(2220)$	$9/2^+$ ****	$\Lambda(1405)$	$1/2^-$ ****	$\Sigma(2080)$	$3/2^+$ **			$\Xi_{cc}^+$	*
$N(2250)$	$9/2^-$ ****	$\Lambda(1520)$	$3/2^-$ ****	$\Sigma(2100)$	$7/2^-$ **			$\Lambda_b^0$	$1/2^+$ ***
$N(2300)$	$1/2^+$ **	$\Lambda(1600)$	$1/2^+$ ***	$\Sigma(2250)$	***			$\Lambda_b(5912)^0$	$1/2^-$ ***
$N(2570)$	$5/2^-$ **	$\Lambda(1670)$	$1/2^-$ ****	$\Sigma(2455)$	**			$\Lambda_b(5920)^0$	$3/2^-$ ***
$N(2600)$	$11/2^-$ ***	$\Lambda(1690)$	$3/2^-$ ****	$\Sigma(2620)$	**			$\Sigma_b$	$1/2^+$ ***
$N(2700)$	$13/2^+$ **	$\Lambda(1710)$	$1/2^+$ *	$\Sigma(3000)$	*			$\Sigma_b^+$	$3/2^+$ ***
		$\Lambda(1800)$	$1/2^-$ ***	$\Sigma(3170)$	*			$\Xi_b^-, \Xi_b^-$	$1/2^+$ ***
		$\Lambda(1810)$	$1/2^+$ ***					$\Xi_b^-, \Xi_b^-$	$1/2^+$ ***
		$\Lambda(1820)$	$5/2^+$ ****					$\Xi_b^-, \Xi_b^-$	$1/2^+$ ***
		$\Lambda(1830)$	$5/2^-$ ****					$\Xi_b^-, \Xi_b^-$	$3/2^+$ ***
		$\Lambda(1890)$	$3/2^+$ ****					$\Xi_b^-, \Xi_b^-$	$3/2^+$ ***
		$\Lambda(2000)$	*					$\Omega_b$	$1/2^+$ ***
		$\Lambda(2020)$	$7/2^+$ *						
		$\Lambda(2050)$	$3/2^-$ *						
		$\Lambda(2100)$	$7/2^-$ ****						
		$\Lambda(2110)$	$5/2^+$ ***						
		$\Lambda(2325)$	$3/2^-$ *						
		$\Lambda(2350)$	$9/2^+$ ***						
		$\Lambda(2585)$	**						



~ 150 baryons

LIGHT UNFLAVORED (S = C = B = 0)		STRANGE (S = ±1, C = B = 0)		CHARMED, STRANGE (C = S = ±1)	$c\bar{c}$ $F_c(F_c^c)$
$F_c(F_c^c)$	$F_c(F_c^c)$	$F_c(F_c^c)$	$F_c(F_c^c)$	$F_c(F_c^c)$	$F_c(F_c^c)$
$\pi^+$	$1^-(0^-)$	$\rho(1680)$	$0^-(1^-)$	$D_s^+$	$0^-(0^-)$
$\pi^0$	$1^-(0^-)$	$\rho(1690)$	$1^+(3^-)$	$D_s^0$	$0^-(0^-)$
$\eta$	$0^+(0^-)$	$\rho(1700)$	$1^+(1^-)$	$D_s^-(2317)^-$	$0^+(0^+)$
$\eta(500)$	$0^+(0^+)$	$\rho(1700)$	$1^-(2^+)$	$D_{s1}(2460)^+$	$0^+(1^+)$
$\rho(770)$	$1^+(1^-)$	$\eta(1710)$	$0^+(0^+)$	$D_{s1}(2536)^+$	$0^+(1^+)$
$\omega(782)$	$0^-(1^-)$	$\eta(1760)$	$0^+(0^-)$	$D_{s2}(2573)$	$0^?(2^?)$
$\eta(958)$	$0^+(0^+)$	$\eta(1800)$	$1^-(0^-)$	$D_{s2}(2700)^+$	$0^-(1^-)$
$\eta(980)$	$0^+(0^+)$	$\eta(1810)$	$0^+(2^+)$	$D_{s1}(2860)^+$	$0^?(2^?)$
$\omega(980)$	$1^-(0^+)$	$X(1835)$	$2^?(2^?)$	$D_{s1}(3040)^+$	$0^?(2^?)$
$\phi(1020)$	$0^-(1^-)$	$X(1840)$	$2^?(2^?)$		
$h_1(1170)$	$0^-(1^+)$	$\eta_3(1850)$	$0^-(3^-)$		
$h_1(1235)$	$1^+(1^+)$	$\eta_3(1880)$	$0^+(2^+)$		
$\omega(1260)$	$1^+(1^+)$	$\eta_3(1890)$	$1^-(2^+)$		
$f_2(1270)$	$0^+(2^+)$	$\eta_3(1900)$	$1^+(1^-)$		
$f_1(1285)$	$0^+(1^+)$	$f_2(1910)$	$1^+(1^-)$		
$\eta(1295)$	$0^+(0^+)$	$f_2(1950)$	$0^+(2^+)$		
$\pi(1300)$	$1^-(0^+)$	$\eta_3(1990)$	$1^+(3^-)$		
$\omega(1320)$	$1^-(2^+)$	$f_2(2010)$	$0^+(2^+)$		
$f_0(1370)$	$0^+(0^+)$	$f_2(2020)$	$0^+(0^+)$		
$h_1(1380)$	$1^-(1^+)$	$\omega(2040)$	$1^-(4^+)$		
$\eta_3(1400)$	$1^-(1^+)$	$f_0(2050)$	$0^+(4^+)$		
$\eta(1405)$	$0^+(0^+)$	$\eta_3(2100)$	$1^-(2^+)$		
$f_1(1420)$	$0^+(1^+)$	$f_2(2100)$	$0^+(0^+)$		
$\omega(1420)$	$0^-(1^-)$	$f_2(2150)$	$0^+(2^+)$		
$f_2(1430)$	$0^+(2^+)$	$\rho(2150)$	$1^-(1^-)$		
$\omega(1450)$	$1^-(0^+)$	$\rho(2170)$	$0^-(1^-)$		
$\rho(1450)$	$1^+(1^-)$	$f_2(2200)$	$0^+(0^+)$		
$\eta(1475)$	$0^+(0^+)$	$f_2(2220)$	$0^+(2^+)$		
$f_0(1500)$	$0^+(0^+)$	$\eta(2225)$	$0^+(0^+)$		
$f_1(1510)$	$0^+(1^+)$	$\eta_3(2250)$	$1^+(3^-)$		
$f_2(1525)$	$0^+(2^+)$	$f_2(2300)$	$0^+(2^+)$		
$f_2(1565)$	$0^+(2^+)$	$f_2(2300)$	$0^+(4^+)$		
$\omega(1570)$	$1^+(1^+)$	$f_2(2330)$	$0^+(0^+)$		
$h_1(1595)$	$0^-(1^+)$	$f_2(2340)$	$0^+(2^+)$		
$\eta_3(1600)$	$1^-(1^+)$	$\eta_3(2350)$	$1^+(5^-)$		
$\omega(1640)$	$1^-(1^+)$	$\omega(2450)$	$1^-(6^+)$		
$f_2(1640)$	$0^+(2^+)$	$\omega(2510)$	$0^+(6^+)$		
$\eta_3(1645)$	$0^+(2^+)$				
$\omega(1650)$	$0^-(1^-)$				
$\omega_3(1670)$	$0^-(3^-)$				
$\eta_3(1670)$	$1^-(2^+)$				



~ 210 mesons

- **stable/unstable** via strong interaction
- Excited states are **mostly unstable**. → resonances

# Difficulty of resonances

## Resonance as an “eigenstate” of Hamiltonian

### - complex energy

G. Gamow, *Z. Phys.* **51**, 204 (1928)

Zur Quantentheorie des Atomkernes.

Von G. Gamow, z. Zt. in Göttingen.

Mit 5 Abbildungen. (Eingegangen am 2. August 1928.)

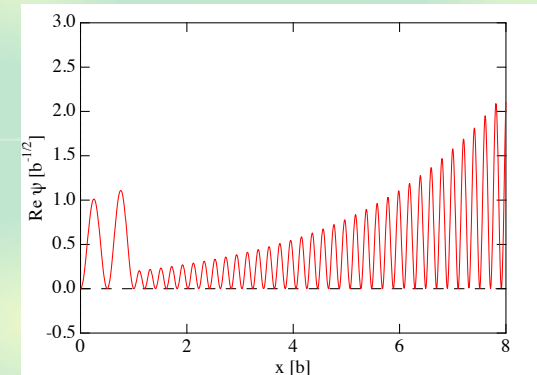
Um diese Schwierigkeit zu überwinden, müssen wir annehmen, daß die Schwingungen gedämpft sind, und  $E$  komplex setzen:

$$E = E_0 + i \frac{\hbar \lambda}{4\pi},$$

wo  $E_0$  die gewöhnliche Energie ist und  $\lambda$  das Dämpfungsdekrement (Zerfallskonstante). Dann sehen wir aber aus den Relationen (2a) und (2b),

### - diverging wave function ( $\text{Im } k < 0$ )

$$\langle R | R \rangle = \int d\mathbf{r} |\psi_R(\mathbf{r})|^2 \sim \int_0^\infty dr e^{-2\text{Im}[k]r} \rightarrow \infty$$



## Bi-orthogonal basis (Gamow vectors): normalizable!

N. Hokkyo, *Prog. Theor. Phys.* **33**, 1116 (1965)

T. Berggren, *Nucl. Phys. A* **109**, 265 (1968)

$$|\tilde{R}\rangle = |R^*\rangle, \quad |\langle \tilde{R} | R \rangle| = \left| \int d\mathbf{r} [\psi_R(\mathbf{r})]^2 \right| < \infty$$

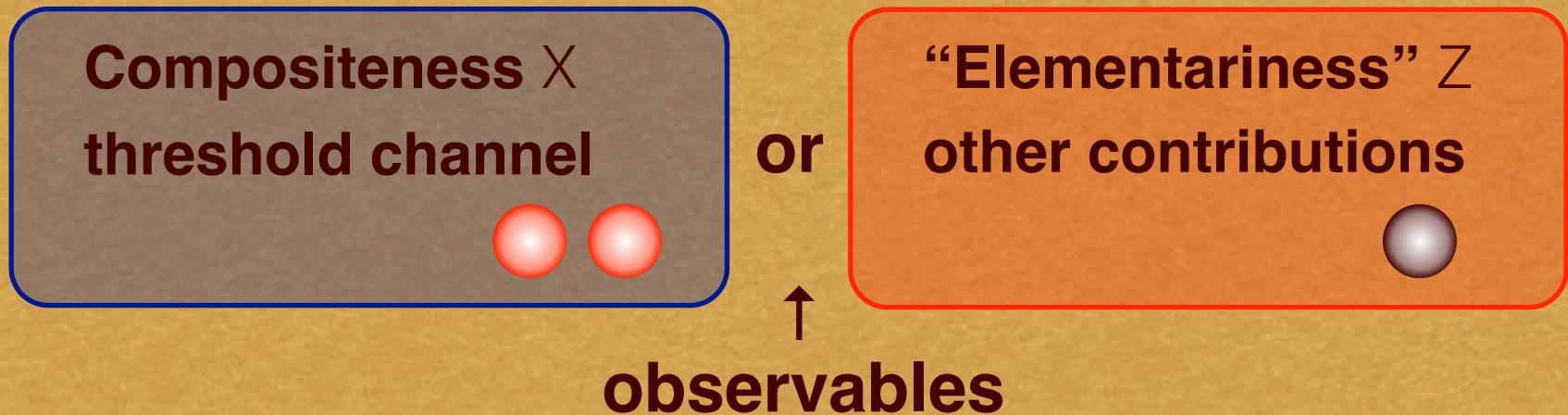
### - Complex expectation value (norm, $\langle r^2 \rangle$ ) $\rightarrow$ interpretation?



# Compositeness of hadrons

- Structure of unstable state is **nontrivial**.
- Weak binding relation for stable bound states

*S. Weinberg, Phys. Rev. 137, B672 (1965)*



- Effective field theory  $\rightarrow$  description of low-energy scattering amplitude, generalization to **unstable** resonances

# Weak binding relation for stable states

**Compositeness  $X$  of s-wave weakly bound state ( $R \gg R_{\text{typ}}$ )**

**S. Weinberg, Phys. Rev. 137, B672 (1965);**

**T. Hyodo, Int. J. Mod. Phys. A 28, 1330045 (2013)**

$$|d\rangle = \sqrt{X} |NN\rangle + \sqrt{1-X} |\text{others}\rangle$$

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\}, \quad r_e = R \left\{ \frac{X-1}{X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\}$$

**$a_0$ : scattering length,  $r_e$ : effective range**

**$R = (2\mu B)^{-1/2}$ : radius of wave function**

**$R_{\text{typ}}$ : length scale of interaction**

- **Deuteron is NN composite ( $a_0 \sim R \gg r_e$ )  $\rightarrow X \sim 1$**
- **Internal structure from observable**

**Problem: applicable only for stable states.**

# Effective field theory

## Low-energy scattering with near-threshold bound state

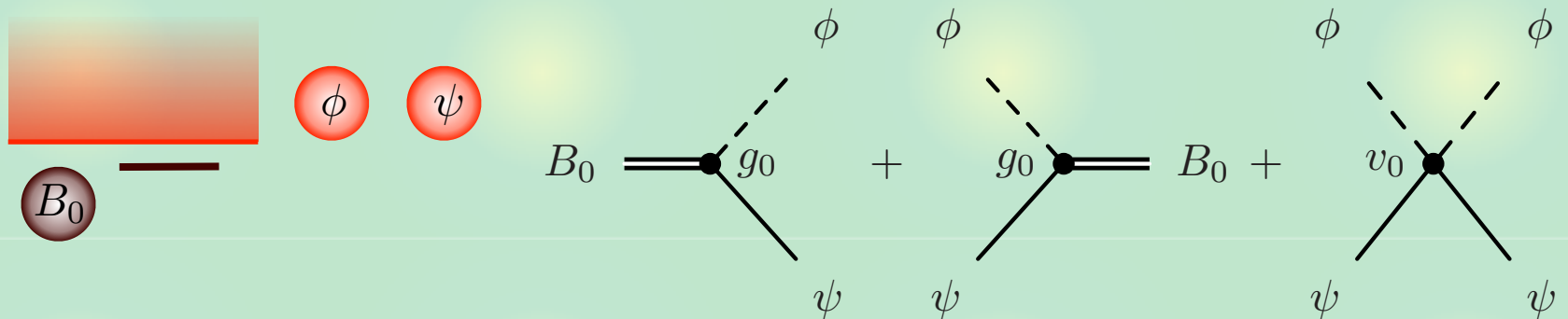
### - Nonrelativistic EFT with contact interaction

D.B. Kaplan, Nucl. Phys. B494, 471 (1997)

E. Braaten, M. Kusunoki, D. Zhang, Annals Phys. 323, 1770 (2008)

$$H_{\text{free}} = \int d\mathbf{r} \left[ \frac{1}{2M} \nabla \psi^\dagger \cdot \nabla \psi + \frac{1}{2m} \nabla \phi^\dagger \cdot \nabla \phi + \frac{1}{2M_0} \nabla B_0^\dagger \cdot \nabla B_0 + \nu_0 B_0^\dagger B_0 \right],$$

$$H_{\text{int}} = \int d\mathbf{r} \left[ g_0 \left( B_0^\dagger \phi \psi + \psi^\dagger \phi^\dagger B_0 \right) + v_0 \psi^\dagger \phi^\dagger \phi \psi \right]$$



- **cutoff** :  $\Lambda \sim 1/R_{\text{typ}}$  (interaction range of microscopic theory)
- At low energy  $p \ll \Lambda$ , interaction  $\sim$  contact



# Compositeness and “elementariness”

## Eigenstates

$$H_{\text{free}}|B_0\rangle = \nu_0|B_0\rangle, \quad H_{\text{free}}|\mathbf{p}\rangle = \frac{p^2}{2\mu}|\mathbf{p}\rangle \quad \text{free (discrete + continuum)}$$

$$(H_{\text{free}} + H_{\text{int}})|B\rangle = -B|B\rangle \quad \text{full (bound state)}$$

### - normalization of $|B\rangle$ + completeness relation

$$\langle B|B\rangle = 1, \quad 1 = |B_0\rangle\langle B_0| + \int \frac{d\mathbf{p}}{(2\pi)^3} |\mathbf{p}\rangle\langle \mathbf{p}|$$

### - projections onto free eigenstates

$$1 = Z + X, \quad \underline{Z \equiv |\langle B_0|B\rangle|^2}, \quad \underline{X \equiv \int \frac{d\mathbf{p}}{(2\pi)^3} |\langle \mathbf{p}|B\rangle|^2}$$

“elementariness”      compositeness



$Z, X$ : real and nonnegative  $\rightarrow$  interpreted as **probability**

# Weak binding relation

## $\Psi\Phi$ scattering amplitude (exact result)

$$f(E) = -\frac{\mu}{2\pi} \frac{1}{[v(E)]^{-1} - G(E)}$$

$$v(E) = v_0 + \frac{g_0^2}{E - \nu_0}, \quad G(E) = \frac{1}{2\pi^2} \int_0^\Lambda dp \frac{p^2}{E - p^2/(2\mu) + i0^+}$$

## Compositeness $X \leftarrow -v(E), G(E)$

$$X = \frac{G'(-B)}{G'(-B) - [1/v(-B)]'}$$

$1/R = (2\mu B)^{1/2}$  expansion: leading term  $\leftarrow X$

$$a_0 = -f(E=0) = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\} \text{ renormalization dependent}$$

renormalization independent

If  $R \gg R_{\text{typ}}$ , correction terms neglected:  $X \leftarrow (B, a_0)$

# Introduction of decay channel

## Introduce decay channel

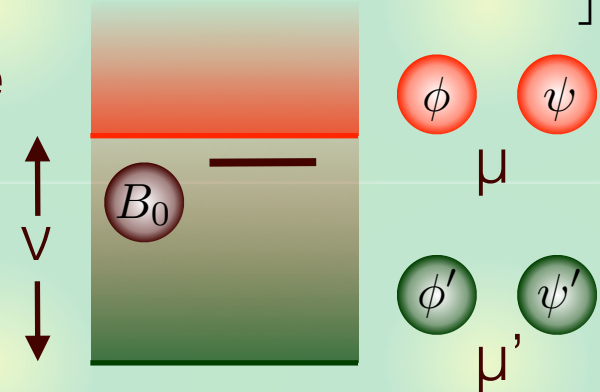
$$H'_{\text{free}} = \int d\mathbf{r} \left[ \frac{1}{2M'} \nabla \psi'^{\dagger} \cdot \nabla \psi' - \nu_{\psi} \psi'^{\dagger} \psi' + \frac{1}{2m'} \nabla \phi'^{\dagger} \cdot \nabla \phi' - \nu_{\phi} \phi'^{\dagger} \phi' \right],$$

$$H'_{\text{int}} = \int d\mathbf{r} \left[ g'_0 \left( B_0^{\dagger} \phi' \psi' + \psi'^{\dagger} \phi'^{\dagger} B_0 \right) + v'_0 \psi'^{\dagger} \phi'^{\dagger} \phi' \psi' + v_0^t (\psi^{\dagger} \phi^{\dagger} \phi' \psi' + \psi'^{\dagger} \phi'^{\dagger} \phi \psi) \right],$$

## Quasi-bound state: complex eigenvalue

$$H = H_{\text{free}} + H'_{\text{free}} + H_{\text{int}} + H'_{\text{int}}$$

$$H |QB\rangle = E_{QB} |QB\rangle, \quad E_{QB} \in \mathbb{C}$$



## Generalized relation: **correction term** ← threshold difference

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O} \left( \left| \frac{R_{\text{typ}}}{R} \right| \right) + \mathcal{O} \left( \left| \frac{l}{R} \right|^3 \right) \right\}, \quad R = \frac{1}{\sqrt{-2\mu E_{QB}}} \in \mathbb{C}, \quad l \equiv \frac{1}{\sqrt{2\mu\nu}}$$

Y. Kamiya, T. Hyodo, Phys. Rev. C93, 035203 (2016)

c.f. V. Baru, *et al.*, Phys. Lett. B586, 53 (2004),...

If  $|R| \gg (R_{\text{typ}}, l)$  correction terms neglected:  $X \leftarrow (E_{QB}, a_0)$



# Complex compositeness

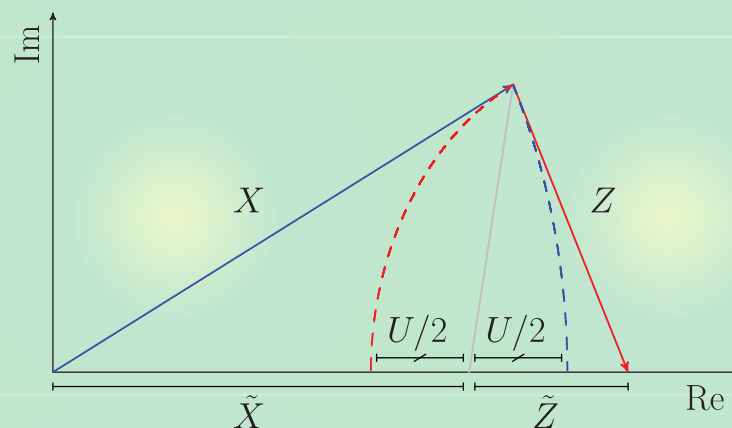
Unstable states  $\rightarrow$  complex  $Z$  and  $X$

$$Z + X = 1, \quad Z, X \in \mathbb{C}$$

- Probabilistic interpretation?

New definition

$$\tilde{Z} = \frac{1 - |X| + |Z|}{2}, \quad \tilde{X} = \frac{1 - |Z| + |X|}{2}$$



- interpreted as **probabilities**  $\tilde{Z} + \tilde{X} = 1, \quad \tilde{Z}, \tilde{X} \in [0, 1]$

- reduces to  $Z$  and  $X$  in the bound state limit

$U/2$ : uncertainty of interpretation

$$U = |Z| + |X| - 1$$

c.f. T. Berggren, Phys. Lett. 33B, 547 (1970)

- Sensible interpretation only for small  $U/2$  case

# Application: $\Lambda(1405)$

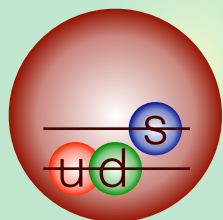
## Generalized weak binding relation

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\left|\frac{R_{\text{typ}}}{R}\right|\right) + \mathcal{O}\left(\left|\frac{l}{R}\right|^3\right) \right\}, \quad R = \frac{1}{\sqrt{-2\mu E_{QB}}}, \quad l \equiv \frac{1}{\sqrt{2\mu\nu}}$$

- We can determine  $X$  from  $(E_{QB}, a_0)$

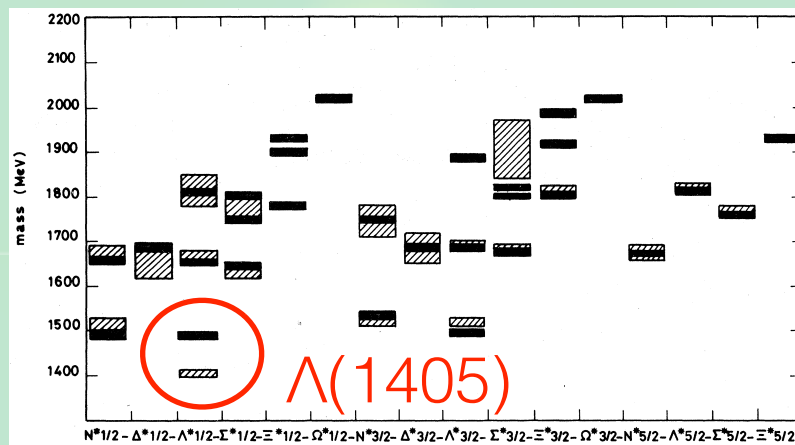
## $\Lambda(1405)$ : exotic candidate

N. Isgur and G. Karl, Phys. Rev. D18, 4187 (1978)



— : th.

▨ : exp.



energy ↑

—  $\bar{K}N$  threshold  
 ▨  $\Lambda(1405)$   
 —  $\pi\Sigma$  threshold

$(E_{QB}, a_0) \leftarrow$  Recent theoretical analysis

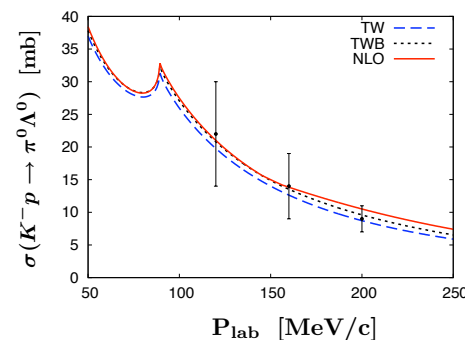
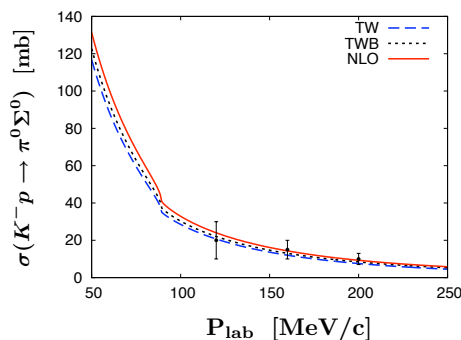
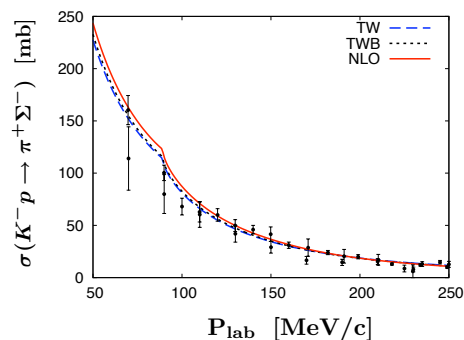
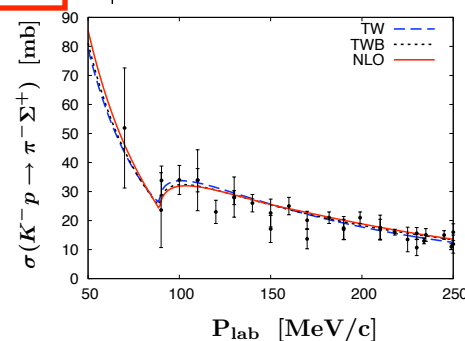
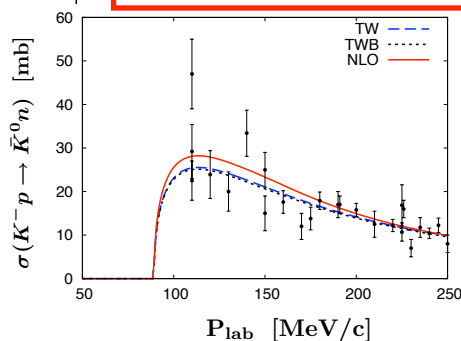
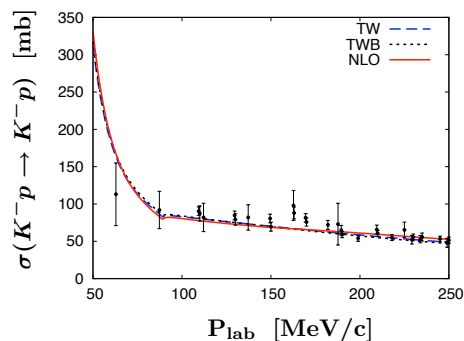
# Fit to experiments: NLO chiral SU(3) dynamics

**SIDDHARTA**

**Branching ratios**

	TW	TWB	NLO	Experiment
$\Delta E$ [eV]	373	377	306	$283 \pm 36 \pm 6$ [10]
$\Gamma$ [eV]	495	514	591	$541 \pm 89 \pm 22$ [10]
$\gamma$	2.36	2.36	2.37	$2.36 \pm 0.04$ [11]
$R_n$	0.20	0.19	0.19	$0.189 \pm 0.015$ [11]
$R_c$	0.66	0.66	0.66	$0.664 \pm 0.011$ [11]
$\chi^2/\text{d.o.f}$	1.12	1.15	0.96	

**cross sections**



Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881 98 (2012)

—> determination of  $(E_{\text{QB}}, a_0)$  for  $\Lambda(1405)$



# Evaluation of compositeness

## Generalized weak-binding relation

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O} \left( \left| \frac{R_{\text{typ}}}{R} \right| \right) + \mathcal{O} \left( \left| \frac{l}{R} \right|^3 \right) \right\}, \quad R = \frac{1}{\sqrt{-2\mu E_{QB}}}, \quad l \equiv \frac{1}{\sqrt{2\mu\nu}}$$

## $(E_{QB}, a_0)$ determinations by several groups

### - neglecting correction terms:

	$E_h$ [MeV]	$a_0$ [fm]	$X_{\bar{K}N}$	$\tilde{X}_{\bar{K}N}$	$U/2$
Set 1 [35]	$-10 - i26$	$1.39 - i0.85$	$1.2 + i0.1$	1.0	0.3
Set 2 [36]	$-4 - i8$	$1.81 - i0.92$	$0.6 + i0.1$	0.6	0.0
Set 3 [37]	$-13 - i20$	$1.30 - i0.85$	$0.9 - i0.2$	0.9	0.1
Set 4 [38]	$2 - i10$	$1.21 - i1.47$	$0.6 + i0.0$	0.6	0.0
Set 5 [38]	$-3 - i12$	$1.52 - i1.85$	$1.0 + i0.5$	0.8	0.3

### - In all cases, $X \sim 1$ with small $U/2$ (complex nature)

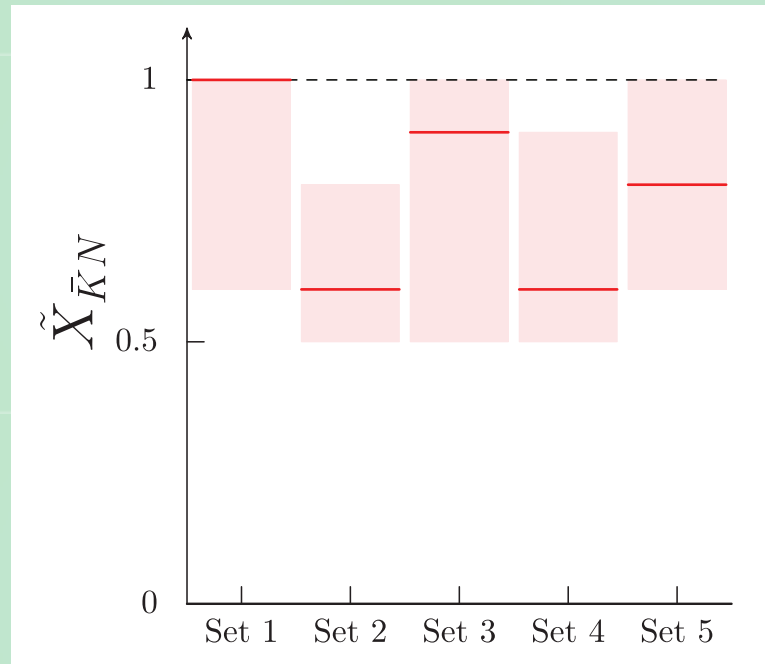
$\Lambda(1405)$  :  $\bar{K}N$  composite dominance  $\leftarrow$  observables

# Uncertainty estimation

Estimation of correction terms :  $|R| \sim 2$  fm

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\left|\frac{R_{\text{typ}}}{R}\right|\right) + \mathcal{O}\left(\left|\frac{l}{R}\right|^3\right) \right\}, \quad R = \frac{1}{\sqrt{-2\mu E_{QB}}}, \quad l \equiv \frac{1}{\sqrt{2\mu\nu}}$$

- $\rho$  meson exchange picture :  $R_{\text{typ}} \sim 0.25$  fm
- energy difference from  $\pi\Sigma$  :  $l \sim 1.08$  fm



$\bar{K}N$  composite dominance holds even **with correction terms.** 18

# Summary

- Compositeness of **near-threshold** bound state can be determined only by **observables**.

S. Weinberg, Phys. Rev. 137, B672 (1965)

- Weak binding relation can be **generalized to unstable states** with effective field theory.

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O} \left( \left| \frac{R_{\text{typ}}}{R} \right| \right) + \mathcal{O} \left( \left| \frac{l}{R} \right|^3 \right) \right\}, \quad R = \frac{1}{\sqrt{-2\mu E_{QB}}}, \quad l \equiv \frac{1}{\sqrt{2\mu\nu}}$$

- Recent determination of  $R$  and  $a_0$  shows that high-mass pole of  $\Lambda(1405)$  is dominated by  **$\bar{K}N$  composite component**.

Y. Kamiya, T. Hyodo, Phys. Rev. C93, 035203 (2016);

Y. Kamiya, T. Hyodo, PTEP2017, 023D02 (2017)

