# $\wedge(1405)$ as a Feshbach resonance 



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## Excitation of hadrons

## Excitation mechanisms

Standard (quark model)



## Exotic excitation


multiquark

hadronic molecule

Exotic structure : excitation inherent in QCD (qव̄ pair creation), different from standard (shell-like) excitation

- Verification? -> compositeness (poster by Kamiya)
- Are they stable particles?


## Unstable states via strong interaction

## Many hadron states



PDG2018 : http://pdg.lbl.gov/


- stable/unstable via strong interaction
- Excited states are mostly unstable. $\rightarrow$ resonances

Current status of $\wedge(1405)$

## $\Lambda(1405)$ and $\bar{K} N$ scattering

$\wedge(1405)$ does not fit in standard picture $\rightarrow>$ exotic candidate
N. Isgur and G. Karl, Phys. Rev. D18, 4187 (1978)


Resonance in coupled-channel scattering

- coupling to MB states


Detailed analysis of $\bar{K} N-\pi \Sigma$ scattering is necessary.

Current status of $\wedge(1405)$

## KN scattering by NLO chiral SU(3) dynamics


Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 88198 (2012)

Accurate description of all existing data ( $\mathrm{x}^{2 / d . o . f . ~} \sim 1$ )

Current status of $\wedge(1405)$

## Subthreshold extrapolation

## Uncertainty of $\bar{K} N \rightarrow \overline{\mathrm{~K} N}(\mathrm{I}=0)$ amplitude below threshold


Y. Kamiya, K. Miyahara, S. Ohnishi, Y. Ikeda, T. Hyodo, E. Oset, W. Weise, NPA 954, 41 (2016)

- c.f. without SIDDHARTA
R. Nissler, Doctoral Thesis (2007)



SIDDHARTA is essential for subthreshold extrapolation.

## Update in PDG

## $\wedge(1405)$ in Particle Data Group (PDG)

## M. Tanabashi, et al., PRD 98, 030001 (2018), http://pdg. lbl.gov/



## 105. Pole Structure of the $\boldsymbol{\Lambda}(1405)$ Region

Written November 2015 by Ulf-G. Meißner (Bonn Univ. / FZ Jülich) and Tetsuo Hyodo (YITP, Kyoto Univ.).

The $\Lambda(1405)$ resonance emerges in the meson-baryon scattering amplitude with the strangeness $S=-1$ and isospin $I=0$. It is the archetype of what is called a dynamically generated resonance, as pioneered by Dalitz and Tuan [1]. The most powerful and


In the 1998 Note on the $\Lambda(1405)$ in PDG 98, R.H. Dalit< uscusseu the S-shaped cusp behavior of the intensity at the $N-\bar{K}$ threshold observed in THOMAS 73 and HEMINGWAY 85. He commented that this behavior "is characteristic of $S$-wave coupling; the other below threshold hyperon, the $\Sigma(1385)$, has no such threshold distortion because its $N-\bar{K}$ coupling is $P$-wave. For $\Lambda(1405)$ this asymmetry is the sole direct evidence that $J^{P}=1 / 2^{-}$."
A recent measurement by the CLAS collaboration, MORIYA 14, definitively established the long-assumed $J^{P}=1 / 2^{-}$spin-parity assignment of the $\Lambda(1405)$. The experiment produced the $\Lambda(1405)$ spin-polarized in the photoproduction process $\gamma p \rightarrow$ $K^{+} \Lambda(1405)$ and measured the decay of the $\Lambda(1405)$ (polarized) $\rightarrow$ $\Sigma^{+}($polarized $) \pi^{-}$. The observed isotropic decay of $\Lambda(1405)$ is consistent with spin $J=1 / 2$. The polarization transfer to the $\Sigma^{+}$(polarized) direction revealed negative parity, and thus established $J^{P}=1 / 2^{-}$
See the related review(s):
Pole Structure of the $\Lambda(1405)$ Region
$\Lambda(1405)$ REGION POLE POSITIONS
REAL PART
$\qquad$

-     - We do not use the following data for averages, fits, limits, etc. - . .

| $1429+8$ | 1 MAI | 15 | DPWA |
| :---: | :---: | :---: | :---: |
| $1325{ }_{-15}^{+15}$ | 2 MAI | 15 | DPWA |
| $1434{ }_{-}^{+} 2$ | ${ }^{3} \mathrm{MAI}$ | 15 | DPWA |
| $1330 \pm 4$ | ${ }^{4} \mathrm{MAI}$ | 15 | DPWA |
| $1421-3$ | ${ }^{5}$ GUO | 13 | DPWA |
| $1388+9$ | 6 Guo | 13 | DPWA |
| $1424{ }_{-23}^{+7}$ | 7 IKEDA | 12 | DPWA |
| $1381+18$ | 8 IKEDA | 12 | DPWA |

- Two-pole structure is confirmed.

Two-pole structure

## Two-pole structure

## Pole position = complex eigenvalue of Hamiltonian

J.R. Taylor, Scattering theory (Wiley, New York, 1972);
T. Hyodo, Intensive lecture at Tohoku Univ. (2018)

- two poles $\rightarrow \wedge(1405)$ is a superposition of two states

Origin in SU(3) basis and implication in spectrum
D. Jido, J.A. Oller, E. Oset, A. Ramos, U.G. Meissner, NPA 723, 205 (2003)

- attraction in 1 and 8 channels of $\operatorname{SU}(3)$ basis
- different channel coupling $\rightarrow$ different $\pi \Sigma$ spectum



Two-pole structure

## Origin in physical basis

## Attraction exists both in $K N$ and $\pi \Sigma$ channels

T. Hyodo, W. Weise, Phys. Rev. C 77, 035204 (2008)



- strong attraction in $\bar{K} N$ : bound state
- attraction in $\pi \Sigma$ : resonance
T. Hyodo, Intensive lecture at SNP school (2017)

Two-pole structure

## Spectrum and pole

## (standard) Feshbach resonance

resonance


Im $E$
Two-pole structure of $\wedge(1405)$
$\wedge(1405)$


Two-pole structure

## Corresponding cold atom system?

${ }^{6} \mathrm{Li}$ atom : large background scattering length
I. Bloch, J. Dalibard, W. Zwerger, Rev. Mod. Phys. 80, 885 (2008)

- theoretical study on large abg
B. Marcelis, et al., Phys. Rev. A 70, 012701 (2004)


FIG. 2. Magnetic field dependence of the scattering length between the two lowest magnetic substates of ${ }^{6} \mathrm{Li}$ with a Feshbach resonance at $B_{0}=834 \mathrm{G}$ and a zero crossing at $B_{0}+\Delta B$ $=534 \mathrm{G}$. The background scattering length $a_{\mathrm{bg}}=-1405 a_{B} \mathrm{~s}$ exceptionally large in this case ( $a_{B}$ the Bohr radius).

- vanishing of scattering length near the unitary limit
$1405!$


CDD zero near pole $\rightarrow$ non-composite nature
Y. Kamiya, T. Hyodo, Phys. Rev. D97, 054019 (2018)

## Summary

On clustering in hadrons

- exotic ( $\overline{\mathrm{q}} \mathrm{q}$ ) excitation: QCD inherent
- hadrons are mostly unstable

Two-pole structure of $\wedge(1405)$

- verified by recent analysis
- Feshbach resonance in resonating continuum
$\because$
$\because$
 $\wedge(1405)$


