

# Status of $\Lambda(1405)$ in chiral dynamics



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# Contents



## $\Lambda(1405)$ in chiral SU(3) dynamics

- **Precise experimental constraint**
- **Determination of pole positions**

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881 98 (2012)



## Kaonic nuclei

- **Local  $\bar{K}N$  potential and  $\Lambda(1405)$  wave function**

K. Miyahara, T. Hyodo, PRC93, 015201 (2016)

- **Density of kaonic nuclei**
- **$\bar{K}N$  v.s. NN correlations**

S. Ohnishi, W. Horiuchi, T. Hoshino, K. Miyahara, T. Hyodo, PRC95, 065202 (2017)

# $\bar{K}$ meson and $\bar{K}N$ interaction

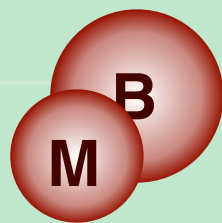
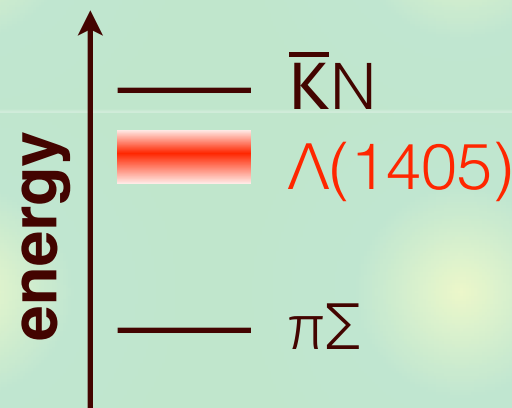
## Two aspects of K( $\bar{K}$ ) meson

- **NG boson** of chiral  $SU(3)_R \otimes SU(3)_L \rightarrow SU(3)_V$
- **Massive** by strange quark:  $m_K \sim 496$  MeV
- $\rightarrow$  **Spontaneous/explicit** symmetry breaking

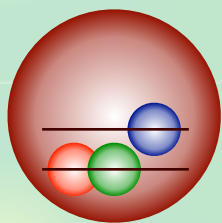
## $\bar{K}N$ interaction ...

T. Hyodo, D. Jido, Prog. Part. Nucl. Phys. 67, 55 (2012)

- is coupled with  $\pi\Sigma$  channel
- generates  $\Lambda(1405)$  below threshold



molecule



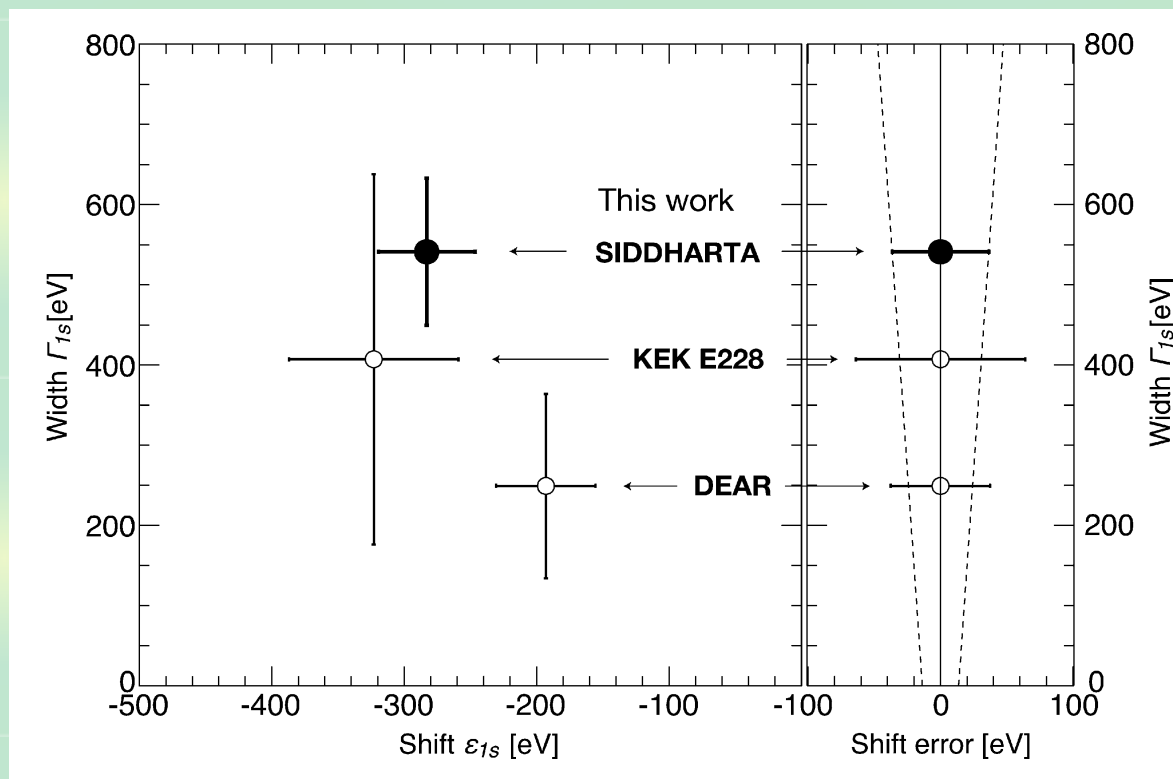
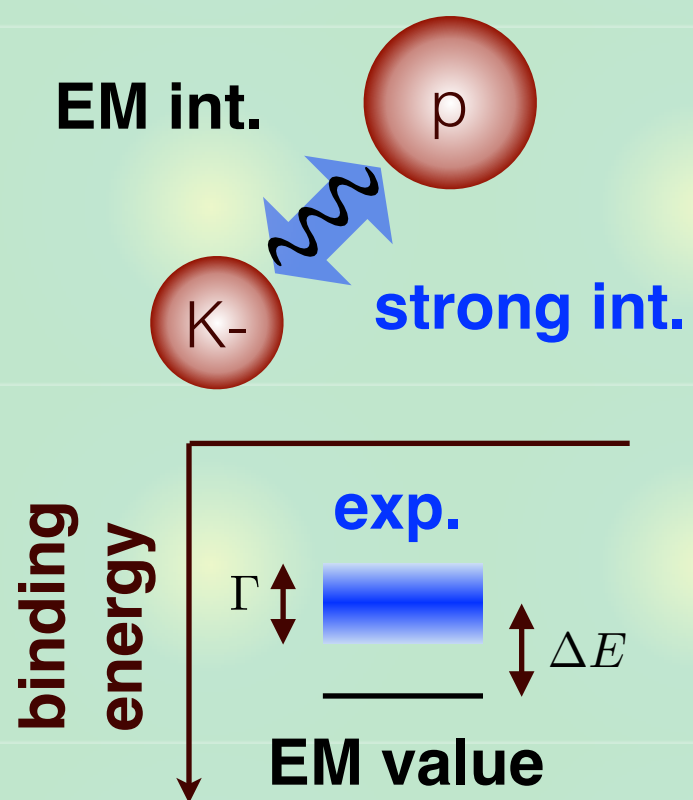
three-quark

- is fundamental building block for  $\bar{K}$ -nuclei,  $\bar{K}$  in medium, ...<sub>3</sub>

# SIDDHARTA measurement

## Precise measurement of the kaonic hydrogen X-rays

M. Bazzi, *et al.*, PLB 704, 113 (2011); NPA 881, 88 (2012)



- Shift and width of atomic state  $\leftrightarrow$   $\bar{K}$ -p scattering length

U.G. Meissner, U. Raha, A. Rusetsky, Eur. Phys. J. C35, 349 (2004)

Quantitative constraint on the  $\bar{K}N$  interaction at fixed energy

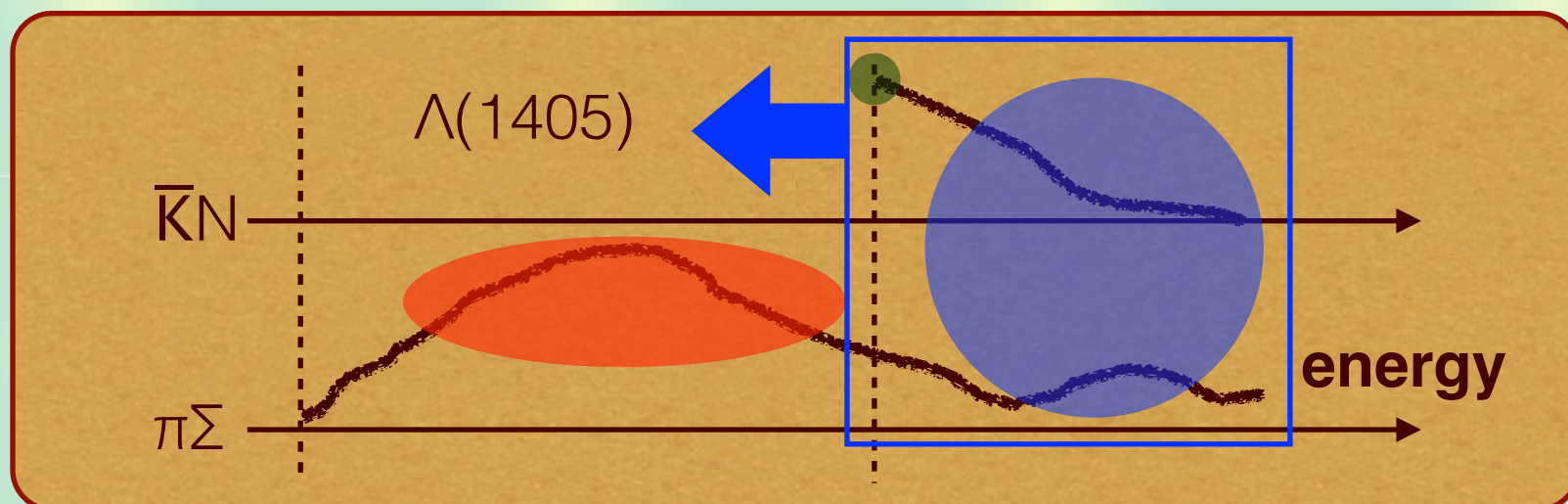
## Strategy for $\bar{K}N$ interaction

**Above the  $\bar{K}N$  threshold: direct constraints**

- $K$ - $p$  **total cross sections** (old data)
- $\bar{K}N$  **threshold branching ratios** (old data)
- $K$ - $p$  **scattering length** (new data: SIDDHARTA)

**Below the  $\bar{K}N$  threshold: indirect constraints**

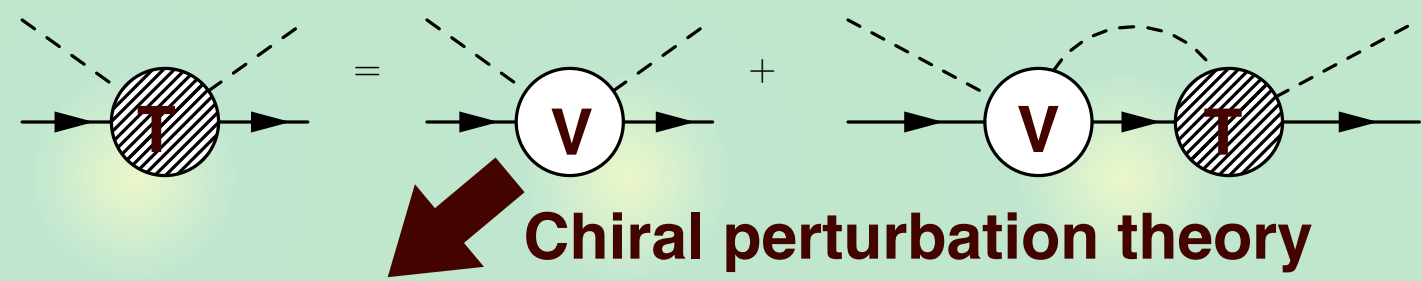
- $\pi\Sigma$  **mass spectra** (new data: LEPS, CLAS, HADES,...)



# Construction of the realistic amplitude

Chiral coupled-channel approach with systematic  $\chi^2$  fitting

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881 98 (2012)



<p><b>1) TW term</b></p> <p><math>\mathcal{O}(p)</math></p> <p><b>6 cutoffs</b></p> <p><b>TW model</b></p>	<p><b>2) Born terms</b></p> <p><math>\mathcal{O}(p)</math></p> <p><b>TWB model</b></p>	<p><b>3) NLO terms</b></p> <p><math>\mathcal{O}(p^2)</math></p> <p><b>7 LECs</b></p> <p><b>NLO model</b></p>
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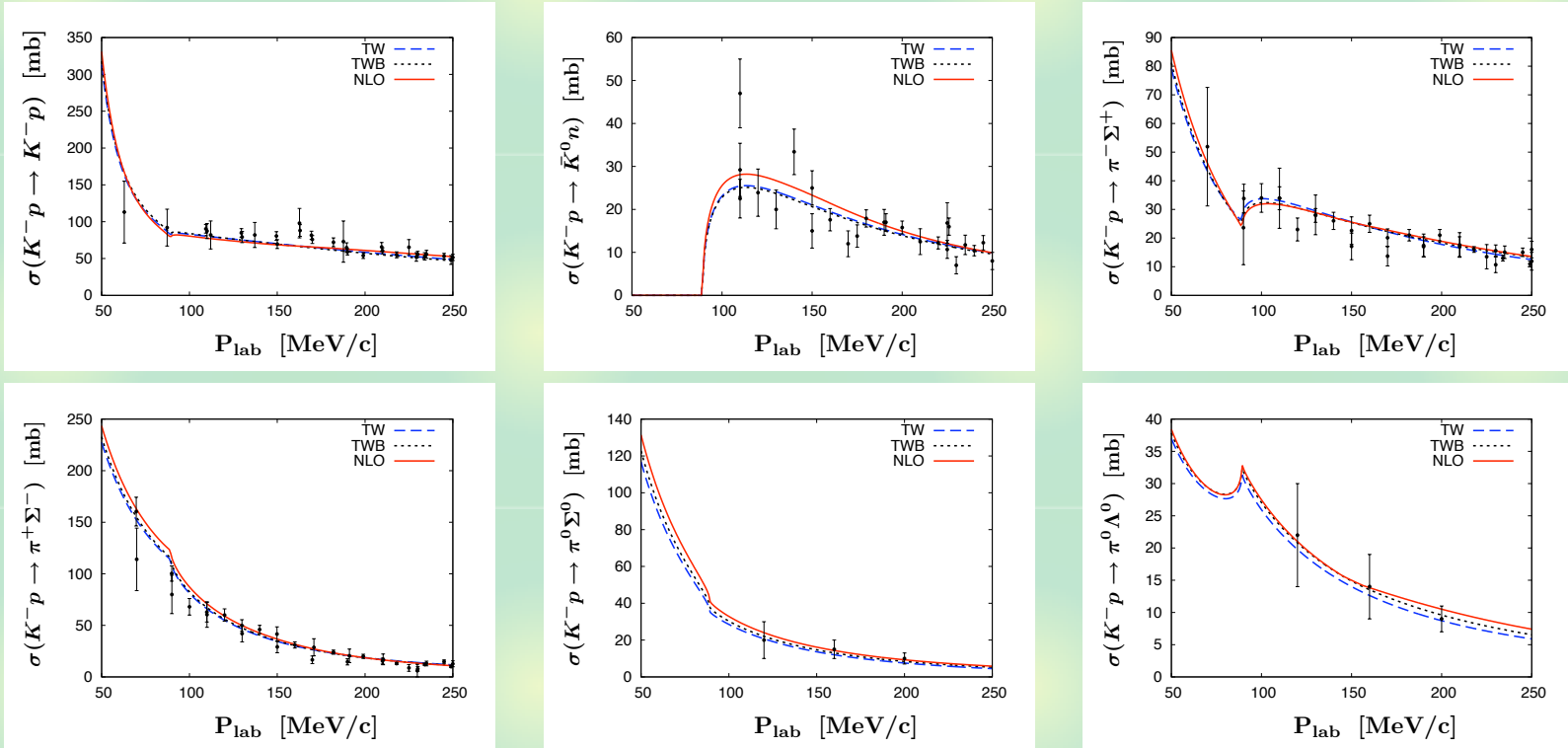
# Best-fit results

**SIDDHARTA**

**Branching ratios**

	TW	TWB	NLO	Experiment
$\Delta E$ [eV]	373	377	306	$283 \pm 36 \pm 6$ [10]
$\Gamma$ [eV]	495	514	591	$541 \pm 89 \pm 22$ [10]
$\gamma$	2.36	2.36	2.37	$2.36 \pm 0.04$ [11]
$R_n$	0.20	0.19	0.19	$0.189 \pm 0.015$ [11]
$R_c$	0.66	0.66	0.66	$0.664 \pm 0.011$ [11]
$\chi^2/\text{d.o.f}$	1.12	1.15	0.96	

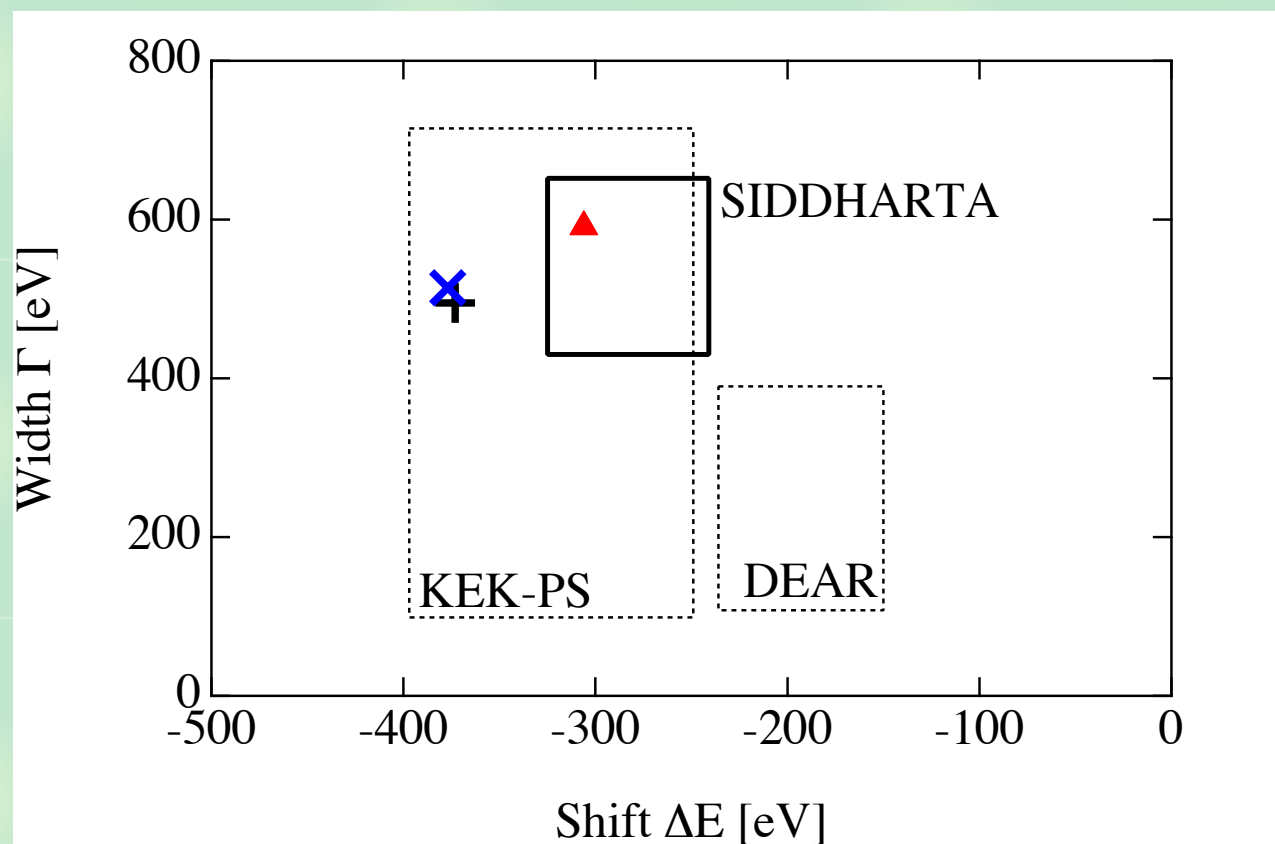
**cross sections**



Accurate description of all existing data ( $\chi^2/\text{d.o.f.} \sim 1$ )

# Comparison with SIDDHARTA

	<b>TW</b>	<b>TWB</b>	<b>NLO</b>
$\chi^2/\text{d.o.f.}$	<b>1.12</b>	<b>1.15</b>	<b>0.957</b>

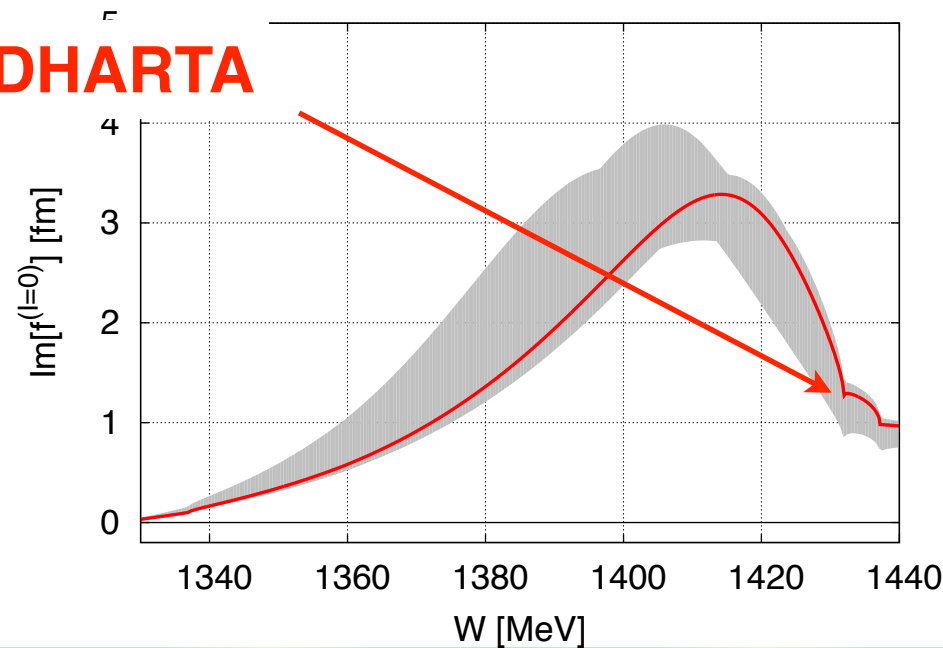
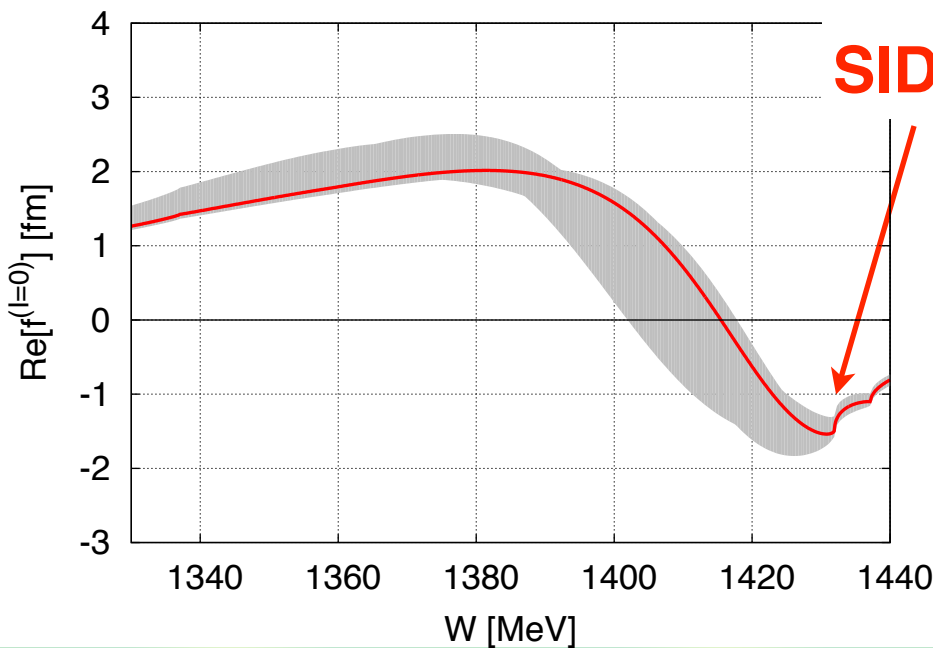


**TW** and **TWB** are reasonable, while best-fit requires **NLO**.



# Subthreshold extrapolation

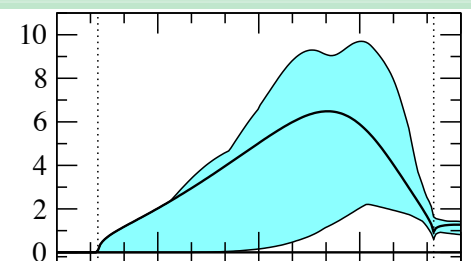
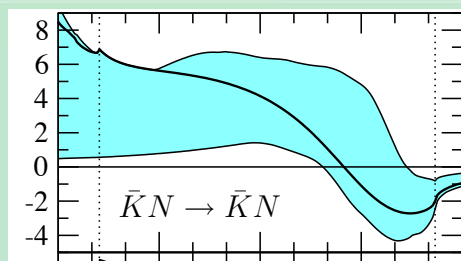
Uncertainty of  $\bar{K}N \rightarrow \bar{K}N$  ( $l=0$ ) amplitude below threshold



Y. Kamiya, K. Miyahara, S. Ohnishi, Y. Ikeda, T. Hyodo, E. Oset, W. Weise, NPA 954, 41 (2016)

- c.f. without SIDDHARTA

R. Nissler, Doctoral Thesis (2007)



SIDDHARTA is essential for **subthreshold** extrapolation.

# Extrapolation to complex energy: two poles

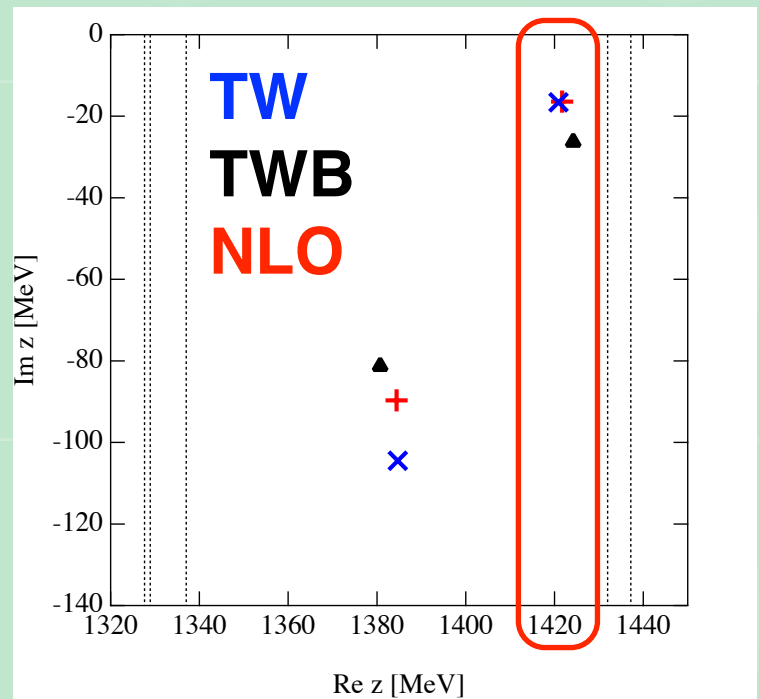
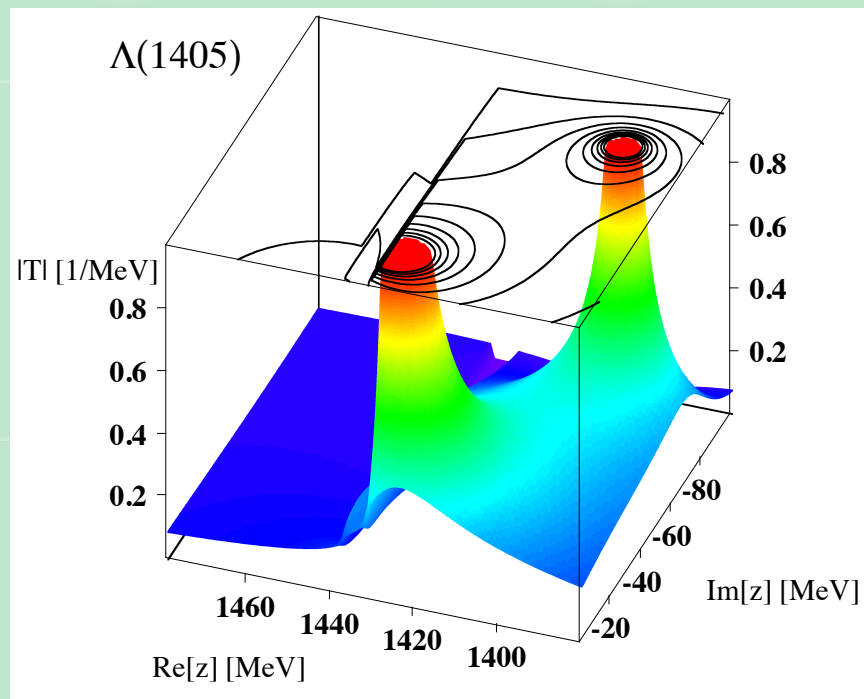
## Two poles: superposition of two states

J.A. Oller, U.G. Meissner, PLB 500, 263 (2001);

D. Jido, J.A. Oller, E. Oset, A. Ramos, U.G. Meissner, NPA 723, 205 (2003);

T. Hyodo, W. Weise, PRC 77, 035204 (2008)

- Higher energy pole at **1420 MeV**, not at 1405 MeV
- Attractions of TW in 1 and 8 ( $\bar{K}N$  and  $\pi\Sigma$ ) channels



NLO analysis confirms the two-pole structure.

# PDG changes

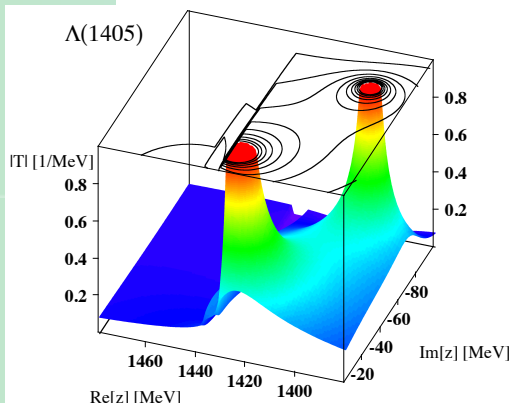
## PDG particle listing of $\Lambda(1405)$

M. Tanabashi, et al., PRD 98, 030001 (2018), <http://pdg.lbl.gov/>

$\Lambda(1405) 1/2^-$

$I(J^P) = 0(\frac{1}{2}^-)$  Status **2014**

The nature of the  $\Lambda(1405)$  has been a puzzle for decades: t... quark state or hybrid; two poles or one. We cannot here survey the rather extensive literature. See, for example, CIEPLY 10, KISSLINGER 11, SEKIHARA 11, and SHEVCHENKO 12A for discussions and earlier references.



of the chiral-unitary community 400-MeV region. ZYCHOR 08 inst the two-pole model, but this REVAI 09, which finds little basis  $\omega$ -pole models; and IKEDA 12,

1405) fits nicely into a  $J^P =$  mer members are the  $\Lambda_c(2595)^+$ , Fig. 1 of our note on "Charmed

MASS

TECN COMMENT

$\Lambda(1405) 1/2^-$

$I(J^P) = 0(\frac{1}{2}^-)$  S **2018**

In the 1998 Note on the  $\Lambda(1405)$  in PDG 98, R.H. Dalitz discussed the S-shaped cusp behavior of the intensity at the  $N\bar{K}$  threshold observed in THOMAS 73 and HEMINGWAY 85. He commented that this behavior "is characteristic of S-wave coupling; the other below threshold hyperon, the  $\Sigma(1385)$ , has no such threshold distortion because its  $N\bar{K}$  coupling is P-wave. For  $\Lambda(1405)$  this asymmetry is the sole direct evidence that  $J^P = 1/2^-$ ."

A recent measurement by the CLAS collaboration, MORIYA 14, definitively established the long-assumed  $J^P = 1/2^-$  spin-parity assignment of the  $\Lambda(1405)$ . The experiment produced the  $\Lambda(1405)$  spin-polarized in the photoproduction process  $\gamma p \rightarrow K^+ \Lambda(1405)$  and measured the decay of the  $\Lambda(1405)$  (polarized)  $\rightarrow \Sigma^+ (\text{polarized}) \pi^-$ . The observed isotropic decay of  $\Lambda(1405)$  is consistent with spin  $J = 1/2$ . The polarization transfer to the  $\Sigma^+$  (polarized) direction revealed negative parity, and thus established  $J^P = 1/2^-$ .

See the related review(s): Pole Structure of the  $\Lambda(1405)$  Region

### $\Lambda(1405)$ REGION POLE POSITIONS

#### REAL PART

VALUE (MeV)	DOCUMENT ID	TECN
1429 <sup>+8</sup> <sub>-7</sub>	1 MAI	15 DPWA
1325 <sup>+15</sup> <sub>-15</sub>	2 MAI	15 DPWA
1434 <sup>+2</sup> <sub>-2</sub>	3 MAI	15 DPWA
1330 <sup>+4</sup> <sub>-5</sub>	4 MAI	15 DPWA
1421 <sup>+3</sup> <sub>-2</sub>	5 GUO	13 DPWA
1388 <sup>+9</sup> <sub>-9</sub>	6 GUO	13 DPWA
1424 <sup>+7</sup> <sub>-23</sub>	7 IKEDA	12 DPWA
1381 <sup>+18</sup> <sub>-6</sub>	8 IKEDA	12 DPWA

### 105. Pole Structure of the $\Lambda(1405)$ Region

Written November 2015 by Ulf-G. Meißner (Bonn Univ. / FZ Jülich) and Tetsuo Hyodo (YITP, Kyoto Univ.).

The  $\Lambda(1405)$  resonance emerges in the meson-baryon scattering amplitude with the strangeness  $S = -1$  and isospin  $I = 0$ . It is the archetype of what is called a dynamically generated resonance, as pioneered by Dalitz and Tuan [1]. The most powerful and

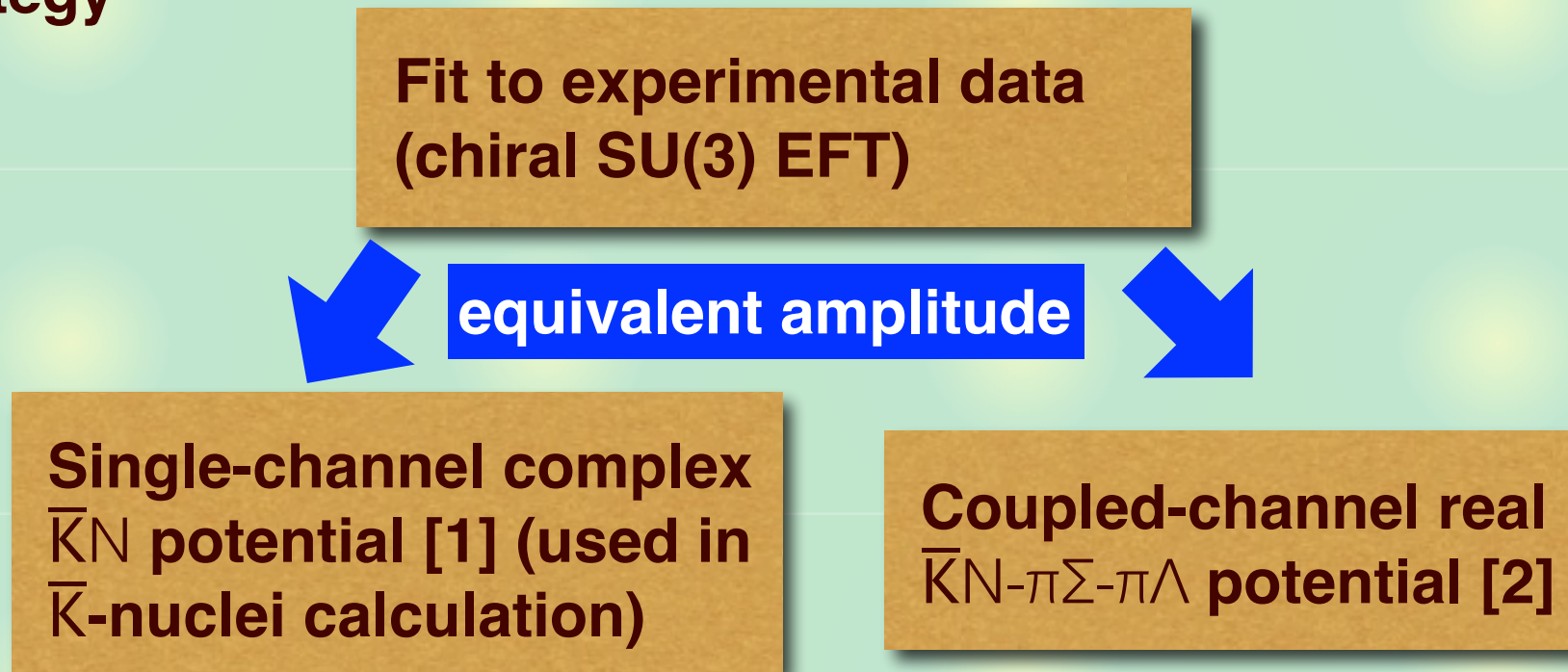
- Our analysis (+ 2 other groups) included
- Pole positions are now tabulated, prior to mass/width.

# Construction of $\bar{K}N$ potential

**Local  $\bar{K}N$  potential** is useful for

- extraction of the wave function of  $\Lambda(1405)$
- application to few-body Kaonic nuclei/atoms

## Strategy



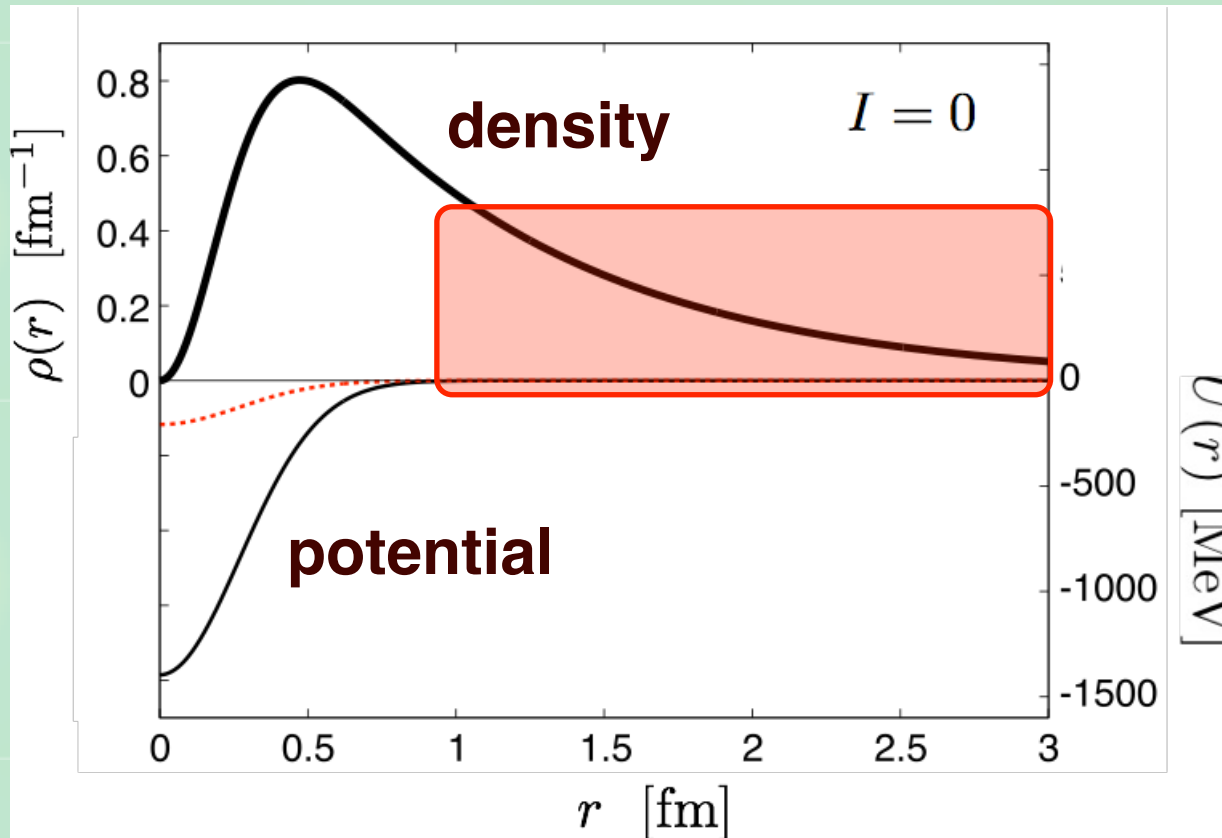
[1] K. Miyahara, T. Hyodo, PRC 93, 015201 (2016);

[2] K. Miyahara, T. Hyodo, W. Weise, PRC 98, 025201 (2018).

# Structure of $\Lambda(1405)$

## $\bar{K}N$ wave function at $\Lambda(1405)$ pole

K. Miyahara, T. Hyodo, PRC93, 015201 (2016)



- substantial distribution at  $r > 1$  fm
- root mean squared radius  $\sqrt{\langle r^2 \rangle} = 1.44$  fm

The **size** of  $\Lambda(1405)$  is much **larger** than ordinary hadrons.

# Kaonic nuclei : current status

**Recent experiment for  $\bar{K}NN$  (J-PARC E15,  ${}^3\text{He}(K^-, \Lambda p)n$ )**

*S. Ajimura, et al., arXiv:1805.12275 [nucl-ex].*

$$B = 47 \pm 3_{-6}^{+3} \text{ MeV}, \quad \Gamma = 115 \pm 7_{-9}^{+10} \text{ MeV}$$

**Theoretical calculation with realistic  $\bar{K}N$  interaction**

- **Fit to K-p cross sections and branching ratios**
- **SIDDHARTHA constraint of Kaonic hydrogen**

[1] J. Revai, N.V. Shevchenko, PRC 90, 034004 (2014),

[2] S. Ohnishi, W. Horiuchi, T. Hoshino, K. Miyahara, T. Hyodo, PRC95, 065202 (2017).

	$V^1$ [1]	$V^2$ [1]	$V^{\text{Chiral}}$ [1]	[2]
B [MeV]	53.3	47.4	32.2	25-28
$\Gamma_{\pi YN}$ [MeV]	64.8	49.8	48.6	31-59

- **2N absorption ( $\Gamma_{YN}$ ) is NOT included.**

# Kaonic nuclei

## Rigorous few-body approach to $\bar{K}$ nuclear systems

S. Ohnishi, W. Horiuchi, T. Hoshino, K. Miyahara, T. Hyodo, PRC95, 065202 (2017).

- Stochastic variational method with correlated gaussians

$$\hat{V} = \hat{V}^{\bar{K}N}(\text{Kyoto } \bar{K}N) + \hat{V}^{NN}(\text{AV4}') \quad (\text{single channel})$$

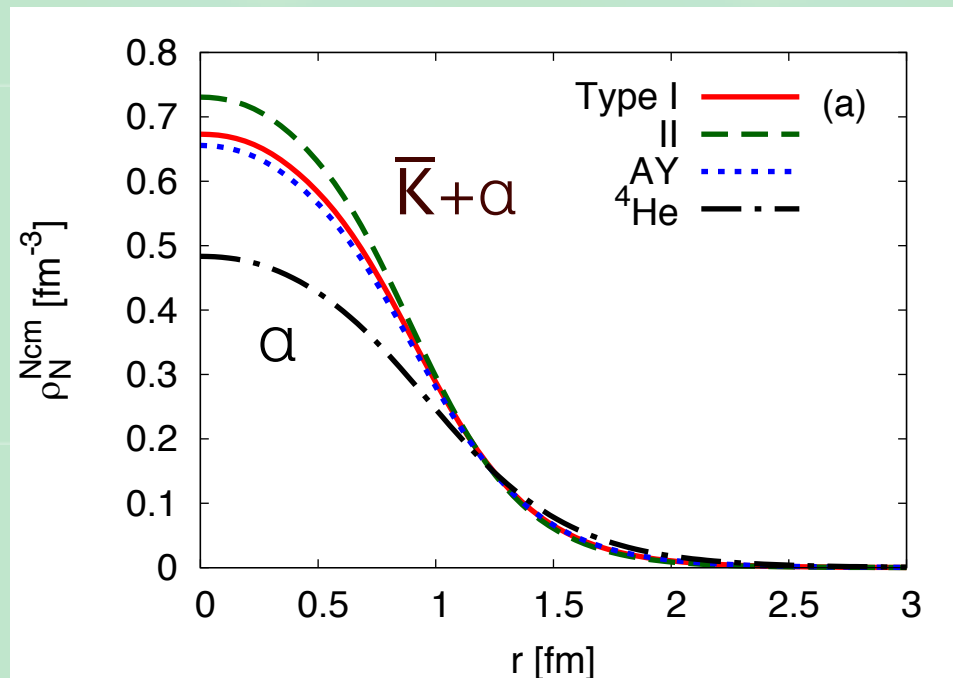
## Results for $A = 2, 3, 4, 6$

	$\bar{K}NN$	$\bar{K}NNN$	$\bar{K}NNNN$	$\bar{K}NNNNN$
B [MeV]	25-28	45-50	68-76	70-81
$\Gamma$ [MeV]	31-59	26-70	28-74	24-76

- **quasi-bound** state below the lowest threshold
- **decay** width (without multi-N absorption)  $\sim$  binding energy

# High density?

## Nucleon density distribution in four-nucleon system



- central density increases (not substantially  $\leftarrow$  NN core)
- $B = 68-76$  MeV (Kyoto  $\bar{K}N$ )
- $B = 85-87$  MeV (AY)

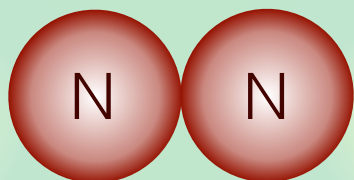
Central density is **not always** proportional to  $B \leftarrow$  tail of w.f.<sub>16</sub>



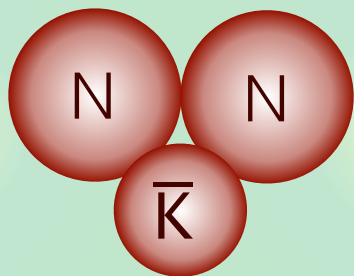
# Interplay between NN and $\bar{K}N$ correlations 1

## Two-nucleon system

$^1S_0$  ( $I_{NN}=1$ )



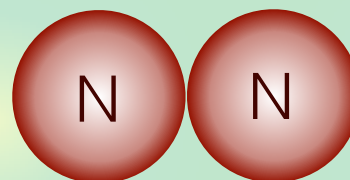
unbound



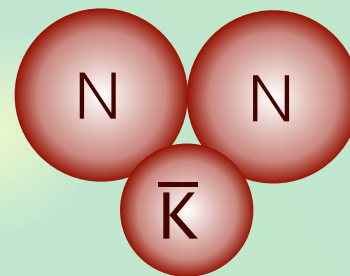
(quasi-)bound

$$\bar{K}N(I=0) : \bar{K}N(I=1) = 3:1$$

$^3S_1$  ( $I_{NN}=0$ )



bound (d)



$\Lambda(1405)$

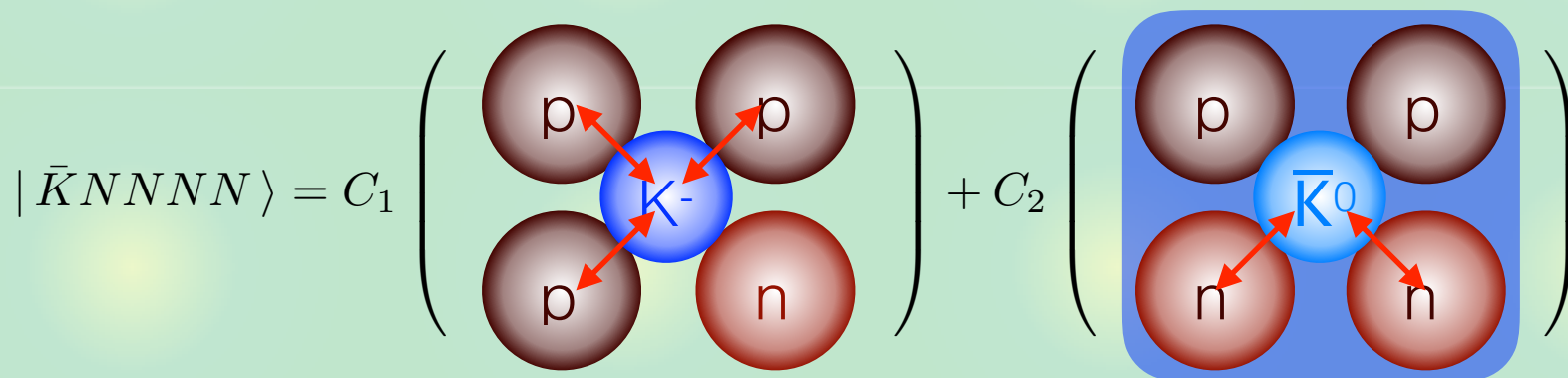
unbound

$$\bar{K}N(I=0) : \bar{K}N(I=1) = 1:3$$

NN correlation  $<$   $\bar{K}N$  correlation (also in  $A=6$ )

# Interplay between NN and $\bar{K}N$ correlations 2

Four-nucleon system with  $J^\pi=0^-$ ,  $I=1/2$ ,  $I_3=+1/2$



-  $\bar{K}N$  correlation

$I=0$  pair in  $K^-p$  (3 pairs) or  $\bar{K}^0n$  (2 pairs) :  $C_1 > C_2$

- NN correlation


ppnn forms  $\alpha$  :  $C_1 < C_2$


- Numerical result

$|C_1|^2 = 0.08$ ,  $|C_2|^2 = 0.92$

NN correlation  $>$   $\bar{K}N$  correlation

**Summary:  $\Lambda(1405)$** 

 **SIDDHARTA measurement of kaonic hydrogen reduces the ambiguity of  $\bar{K}N$  amplitude.**


 **Pole positions of  $\Lambda(1405)$  are determined by fitting all existing data with  $\chi^2/\text{d.o.f.} \sim 1$ .**

$$z_1 = (1424_{-23}^{+7} - i26_{-14}^{+3}) \text{ MeV}, \quad z_2 = (1381_{-6}^{+18} - i81_{-8}^{+19}) \text{ MeV}$$

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881 98 (2012)

 **Realistic  $\bar{K}N$  potential is constructed.**

K. Miyahara, T. Hyodo, PRC93, 015201 (2016)

 **Structure of few-body kaonic nuclei reflects the **interplay** between  $NN$  and  $\bar{K}N$  correlations.**

S. Ohnishi, W. Horiuchi, T. Hoshino, K. Miyahara, T. Hyodo, PRC95, 065202 (2017)