

Compositeness of hadrons from effective field theory



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Contents



Introduction

- Hadron structure and exotics
- Resonance in hadron physics



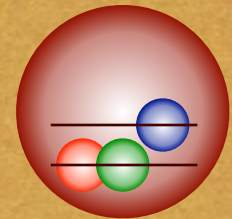
Compositeness of near-threshold resonances

[T. Hyodo, Int. J. Mod. Phys. A 28, 1330045 \(2013\)](#)

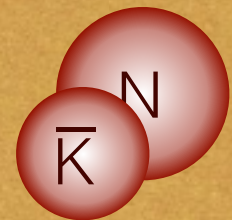
[T. Hyodo, Phys. Rev. Lett. 111, 132002 \(2013\)](#)

- Weak binding relation from EFT
- Generalization to unstable states
- **Application:** $\Lambda(1405)$

[Y. Kamiya, T. Hyodo, arXiv:1509.00146 \[hep-ph\]](#)



or

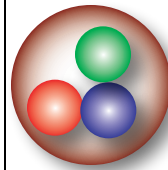


Classification of hadrons

Observed hadrons

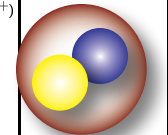
PDG2015 : <http://pdg.lbl.gov/>

p	$1/2^+$ ****	$\Delta(1232)$	$3/2^+$ ****	Σ^+	$1/2^+$ ****	Ξ^0	$1/2^+$ ****	Λ_c^+	$1/2^+$ ****
n	$1/2^+$ ****	$\Delta(1600)$	$3/2^+$ ***	Σ^0	$1/2^+$ ****	Ξ^-	$1/2^+$ ****	$\Lambda_c(2595)^+$	$1/2^-$ ***
$N(1440)$	$1/2^+$ ****	$\Delta(1620)$	$1/2^-$ ****	Σ^-	$1/2^+$ ****	$\Xi(1530)$	$3/2^+$ ****	$\Lambda_c(2625)^+$	$3/2^-$ ***
$N(1520)$	$3/2^-$ ****	$\Delta(1700)$	$3/2^-$ ****	$\Sigma(1385)$	$3/2^+$ ****	$\Xi(1620)$	*	$\Lambda_c(2765)^+$	*
$N(1535)$	$1/2^-$ ****	$\Delta(1750)$	$1/2^+$ *	$\Sigma(1480)$	*	$\Xi(1690)$	*	$\Lambda_c(2940)^+$	$5/2^+$ ***
$N(1650)$	$1/2^-$ ****	$\Delta(1900)$	$1/2^-$ **	$\Sigma(1560)$	**	$\Xi(1820)$	$3/2^-$ ***	$\Lambda_c(2455)$	$1/2^+$ ****
$N(1675)$	$5/2^-$ ****	$\Delta(1905)$	$5/2^+$ ****	$\Sigma(1580)$	$3/2^-$ **	$\Xi(1950)$	*	$\Sigma_c(2520)$	$3/2^+$ ****
$N(1680)$	$5/2^+$ ****	$\Delta(1910)$	$1/2^+$ ****	$\Sigma(1620)$	$1/2^-$ *	$\Xi(2030)$	$\geq 5/2^?$ ***	$\Sigma_c(2800)$	***
$N(1685)$	*	$\Delta(1920)$	$3/2^+$ ***	$\Sigma(1660)$	$1/2^+$ ***	$\Xi(2120)$	*	$\Xi_c(2800)$	***
$N(1700)$	$3/2^-$ ***	$\Delta(1930)$	$5/2^-$ **	$\Sigma(1670)$	$3/2^-$ ****	$\Xi(2250)$	**	Ξ_c	$1/2^+$ ***
$N(1710)$	$1/2^+$ ***	$\Delta(1940)$	$3/2^-$ **	$\Sigma(1690)$	***	Ξ_c	$1/2^+$ ***	Ξ_c	$1/2^+$ ***
$N(1720)$	$3/2^+$ ****	$\Delta(1950)$	$7/2^+$ ****	$\Sigma(1730)$	$3/2^+$ **	Ξ_c	$1/2^+$ ***	Ξ_c	$1/2^+$ ***
$N(1860)$	$5/2^+$ **	$\Delta(2000)$	$5/2^+$ **	$\Sigma(1750)$	$1/2^-$ ***	Ξ_c	$1/2^+$ ***	Ξ_c	$1/2^+$ ***
$N(1875)$	$3/2^-$ ***	$\Delta(2150)$	$1/2^-$ *	$\Sigma(1770)$	$1/2^+$ **	Ω^-	$3/2^+$ ****	Ξ_c	$1/2^+$ ***
$N(1880)$	$1/2^+$ **	$\Delta(2200)$	$7/2^-$ *	$\Sigma(1775)$	$5/2^-$ ****	$\Omega(2250)^-$	***	$\Xi_c(2645)$	$3/2^+$ ***
$N(1895)$	$1/2^-$ **	$\Delta(2300)$	$9/2^+$ **	$\Sigma(1840)$	$3/2^+$ **	$\Omega(2380)^-$	**	$\Xi_c(2790)$	$1/2^-$ ***
$N(1900)$	$3/2^+$ ***	$\Delta(2350)$	$5/2^-$ **	$\Sigma(1880)$	$1/2^+$ **	$\Omega(2470)^-$	**	$\Xi_c(2815)$	$3/2^-$ ***
$N(1990)$	$7/2^+$ **	$\Delta(2390)$	$7/2^+$ *	$\Sigma(1900)$	$1/2^-$ *			$\Xi_c(2930)$	*
$N(2000)$	$5/2^+$ **	$\Delta(2400)$	$9/2^-$ **	$\Sigma(1915)$	$5/2^+$ ****			$\Xi_c(2980)$	**
$N(2040)$	$3/2^+$ *	$\Delta(2420)$	$11/2^+$ ****	$\Sigma(1940)$	$3/2^+$ **			$\Xi_c(3055)$	***
$N(2060)$	$5/2^-$ **	$\Delta(2750)$	$13/2^-$ **	$\Sigma(1940)$	$3/2^-$ ***			$\Xi_c(3080)$	***
$N(2100)$	$1/2^+$ *	$\Delta(2950)$	$15/2^+$ **	$\Sigma(2000)$	$1/2^-$ *			$\Xi_c(3123)$	*
$N(2120)$	$3/2^-$ **			$\Sigma(2030)$	$7/2^+$ ****			Ω_c^0	$1/2^+$ ***
$N(2190)$	$7/2^-$ ****	Λ	$1/2^+$ ****	$\Sigma(2070)$	$5/2^+$ *			$\Omega_c(2770)^0$	$3/2^+$ ***
$N(2220)$	$9/2^+$ ****	$\Lambda(1405)$	$1/2^-$ ****	$\Sigma(2080)$	$3/2^+$ **			Ξ_{cc}^+	*
$N(2250)$	$9/2^-$ ****	$\Lambda(1520)$	$3/2^-$ ****	$\Sigma(2100)$	$7/2^-$ **				
$N(2300)$	$1/2^+$ **	$\Lambda(1600)$	$1/2^+$ **	$\Sigma(2250)$	***			Λ_b^0	$1/2^+$ ***
$N(2570)$	$5/2^-$ **	$\Lambda(1670)$	$1/2^-$ ****	$\Sigma(2455)$	**			$\Lambda_b(5912)^0$	$1/2^-$ ***
$N(2600)$	$11/2^-$ ***	$\Lambda(1690)$	$3/2^-$ ****	$\Sigma(2620)$	**			$\Lambda_b(5920)^0$	$3/2^-$ ***
$N(2700)$	$13/2^+$ **	$\Lambda(1710)$	$1/2^+$ *	$\Sigma(3000)$	**			Σ_b	$1/2^+$ ***
		$\Lambda(1800)$	$1/2^-$ ***	$\Sigma(3170)$	*			Σ_b^+	$3/2^+$ ****
		$\Lambda(1810)$	$1/2^+$ ***					Ξ_b^-	$1/2^+$ ***
		$\Lambda(1820)$	$5/2^+$ ****					Ξ_b'	$1/2^+$ ***
		$\Lambda(1830)$	$5/2^-$ ****					$\Xi_b'(5935)^-$	$1/2^+$ ***
		$\Lambda(1890)$	$3/2^+$ ****					$\Xi_b(5945)^0$	$3/2^+$ ***
		$\Lambda(2000)$	*					$\Xi_b(5955)^-$	$3/2^+$ ***
		$\Lambda(2020)$	$7/2^+$ *					$\Xi_b(5955)^-$	$3/2^+$ ***
		$\Lambda(2050)$	$3/2^-$ *					Ω_b	$1/2^+$ ***
		$\Lambda(2100)$	$7/2^-$ ****						
		$\Lambda(2110)$	$5/2^+$ ***						
		$\Lambda(2325)$	$3/2^-$ *						
		$\Lambda(2350)$	$9/2^+$ ***						
		$\Lambda(2585)$	**						



~ 150 baryons

LIGHT UNFLAVORED (S = C = B = 0)		STRANGE (S = ±1, C = B = 0)		CHARMED, STRANGE (C = S = ±1)		cc	
$F^C(F^C)$		$F^S(F^S)$		$F^C(F^S)$		$F^C(F^C)$	
π^\pm	$1^-(0^-)$	$\rho(1680)$	$0^-(1^-)$	K^{*2}	$1/2(0^-)$	D_s^{*2}	$0(0^-)$
π^0	$1^-(0^-)$	$\rho_3(1690)$	$1^+(3^-)$	K^{*0}	$1/2(0^-)$	D_s^{*0}	$0(0^-)$
η	$0^+(0^+)$	$\omega(1700)$	$1^+(1^-)$	K_2^*	$1/2(0^-)$	D_{s1}^*	$0^+(0^+)$
$\eta(500)$	$0^+(0^+)$	$a_2(1700)$	$1^-(2^+)$	K_1^*	$1/2(0^-)$	$D_{s1}(2460)^0$	$0(1^+)$
$\eta(770)$	$1^+(1^+)$	$f_0(1710)$	$0^+(0^+)$	$K_0^*(800)$	$1/2(0^+)$	$D_{s1}(2536)^0$	$0(1^+)$
$\omega(782)$	$0^-(1^-)$	$\eta(1760)$	$0^+(0^+)$	$K^*(892)$	$1/2(1^-)$	$D_{s2}(2573)$	$0(0^+)$
$\eta(958)$	$0^+(0^+)$	$\omega(1800)$	$1^-(0^+)$	$K_1^*(1270)$	$1/2(1^+)$	$D_{s1}(2700)^0$	$0(1^-)$
$f_0(980)$	$0^+(0^+)$	$f_2(1810)$	$0^+(2^+)$	$K_1^*(1400)$	$1/2(1^+)$	$D_{s1}^*(2860)^0$	$0(0^+)$
$a_0(980)$	$1^-(0^+)$	$X(1835)$	$?^?(2^-)$	$K^*(1410)$	$1/2(1^-)$	$D_{s1}(3040)^0$	$?^?(0^+)$
$\phi(1020)$	$0^-(1^-)$	$X(1840)$	$?^?(2^?)$	$K_0^*(1430)$	$1/2(0^+)$		
$h_1(1170)$	$0^-(1^+)$	$\phi_3(1850)$	$0^-(3^-)$	$K_2^*(1430)$	$1/2(2^+)$	BOTTOM (B = ±1)	
$b_1(1235)$	$1^-(1^+)$	$\eta_2(1870)$	$0^+(2^+)$	$K_1(1460)$	$1/2(0^-)$	B^{*2}	$1/2(0^-)$
$a_1(1260)$	$1^-(1^+)$	$\eta_3(1880)$	$1^-(2^+)$	$K_0^*(1580)$	$1/2(2^-)$	B^0	$1/2(0^-)$
$f_2(1270)$	$0^+(2^+)$	$\rho(1900)$	$1^+(1^-)$	$K(1630)$	$1/2(0^+)$	B^{*2}/B^0	ADMIXTURE
$f_1(1285)$	$0^+(1^+)$	$f_0(1910)$	$0^+(2^+)$	$K_1(1650)$	$1/2(1^+)$	$B^{*2}/B^0/B^0$	b-baryon ADMIXTURE
$\eta(1295)$	$0^+(0^+)$	$f_2(1950)$	$0^+(2^+)$	$K^*(1680)$	$1/2(1^-)$	V_{cb} and V_{cb}	KM Matrix Elements
$\pi(1320)$	$1^-(0^+)$	$\rho_3(1990)$	$1^+(3^-)$	$K_0^*(1770)$	$1/2(1^-)$	B^{*2}	$1/2(1^-)$
$a_2(1320)$	$1^-(2^+)$	$f_2(2010)$	$0^+(2^+)$	$K_3^*(1780)$	$1/2(3^-)$	$B_1(5721)^0$	$1/2(1^+)$
$f_0(1370)$	$0^+(0^+)$	$f_0(2020)$	$0^+(0^+)$	$K_0^*(1820)$	$1/2(2^-)$	$B_1(5721)^+$	$1/2(1^+)$
$\omega(1400)$	$?^-(1^+)$	$a_1(2040)$	$1^-(4^+)$	$K(1830)$	$1/2(0^-)$	$B_1(5721)^+$	$1/2(1^+)$
$\pi_1(1400)$	$1^-(1^+)$	$f_0(2050)$	$0^+(0^+)$	$K_1^*(1950)$	$1/2(0^+)$	$B_1(5732)$	$?^?(0^+)$
$\eta(1405)$	$0^+(0^+)$	$\rho_3(2100)$	$1^-(2^+)$	$K_0^*(1980)$	$1/2(2^+)$	$B_2^*(5747)^0$	$1/2(2^+)$
$f_1(1420)$	$0^+(1^+)$	$f_0(2100)$	$0^+(0^+)$	$K_2^*(1980)$	$1/2(2^+)$	$B_2^*(5747)^0$	$1/2(2^+)$
$\omega(1420)$	$0^-(1^-)$	$f_2(2150)$	$0^+(2^+)$	$K_1^*(2045)$	$1/2(4^+)$	$B(5970)^0$	$?^?(0^+)$
$f_2(1430)$	$0^+(2^+)$	$\rho(2150)$	$1^-(1^-)$	$K_0^*(2250)$	$1/2(3^-)$	$B(5970)^0$	$?^?(0^+)$
$a_1(1450)$	$1^-(0^+)$	$f_0(2200)$	$0^+(0^+)$	$K_3^*(2380)$	$1/2(5^-)$	$B(5970)^0$	$?^?(0^+)$
$\rho(1450)$	$1^+(1^-)$	$f_2(2220)$	$0^+(0^+)$	$K_1^*(2500)$	$1/2(4^-)$	$B(5970)^0$	$?^?(0^+)$
$\eta(1475)$	$0^+(0^+)$	$f_1(2220)$	$0^+(2^+)$	$K(3100)$	$?^?(0^+)$	BOTTOM, STRANGE (B = ±1, S = ±1)	
$f_0(1500)$	$0^+(0^+)$	$\eta(2225)$	$0^+(0^+)$			B_c^+	$0(0^-)$
$f_1(1510)$	$0^+(1^+)$	$\rho_3(2250)$	$1^+(3^-)$	CHARMED (C = ±1)		B_c^0	$0(1^-)$
$f_2(1525)$	$0^+(2^+)$	$f_0(2300)$	$0^+(2^+)$	D^{*2}	$1/2(0^-)$	B_c^+	$0(1^-)$
$f_3(1565)$	$0^+(2^+)$	$f_1(2300)$	$0^+(4^+)$	D^{*0}	$1/2(0^-)$	B_c^0	$0(1^-)$
$\omega(1570)$	$1^-(1^+)$	$f_0(2330)$	$0^+(0^+)$	D^0	$1/2(0^-)$	$B_c^0(5890)^0$	$0(1^+)$
$h_1(1595)$	$0^+(1^+)$	$f_2(2340)$	$0^+(2^+)$	$D^*(2007)^0$	$1/2(1^-)$	$B_c^0(5890)^0$	$0(2^+)$
$\pi_1(1600)$	$1^-(1^+)$	$\rho_3(2350)$	$1^+(5^-)$	$D^*(2010)^0$	$1/2(1^-)$	$B_c^0(5890)^0$	$?^?(0^+)$
$a_1(1640)$	$1^-(1^+)$	$a_1(2450)$	$1^-(6^+)$	$D^*(2400)^0$	$1/2(0^+)$	BOTTOM, CHARMED (B = C = ±1)	
$f_2(1640)$	$0^+(2^+)$	$f_0(2510)$	$0^+(6^+)$	$D_0^*(2400)^0$	$1/2(0^+)$	B_c^+	$0(0^-)$
$\eta_2(1645)$	$0^+(2^-)$			$D_1(2420)^0$	$1/2(1^+)$	B_c^0	$?^?(0^+)$
$\omega(1650)$	$0^-(1^-)$	OTHER LIGHT		$D_1(2430)^0$	$1/2(1^+)$		
$\omega_3(1670)$	$0^-(3^-)$	Further States		$D_2^*(2460)^0$	$1/2(2^+)$		
$\rho_3(1670)$	$1^-(2^-)$			$D_2^*(2460)^0$	$1/2(2^+)$		
				$D(2550)^0$	$1/2(0^-)$		
				$D_2^*(2460)^0$	$1/2(2^+)$		
				$D(2600)$	$1/2(0^+)$		
				$D^*(2640)^0$	$1/2(2^+)$		
				$D(2750)$	$1/2(0^+)$		



~ 200 mesons

All ~ 350 hadrons emerge from single QCD Lagrangian.

All quantum numbers are described by $qqq/q\bar{q}$.

States with more than $qqq/q\bar{q}$

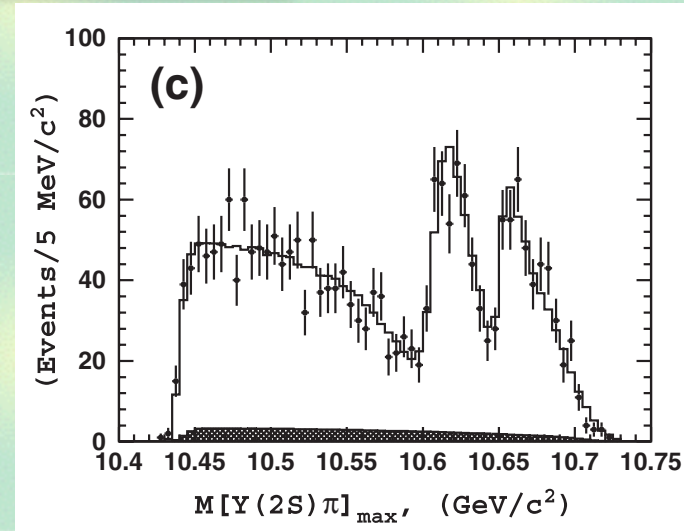
Tetraquark candidate (Belle)

: $Z_b(10610)$, $Z_b(10650)$

$$Y(5S) \longrightarrow \pi^\pm + Z_b$$

$$\hookrightarrow Y(nS)(b\bar{b}) + \pi^\mp(u\bar{d}/d\bar{u})$$

A. Bondar, *et al.*, *Phys. Rev. Lett.* **108**, 122001 (2012)



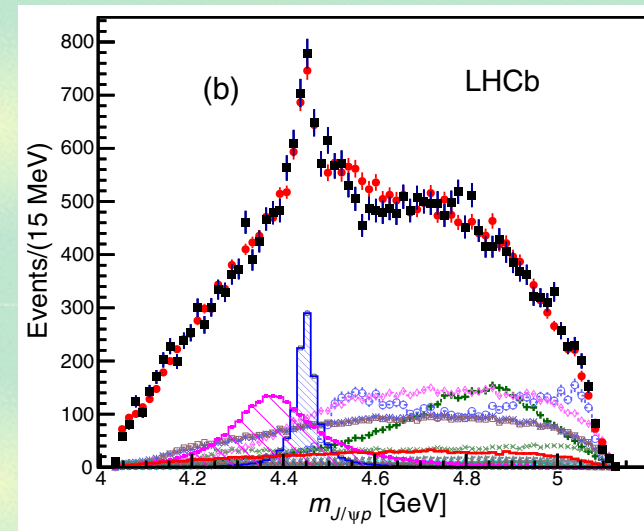
Pentaquark candidate (LHCb)

: $P_c(4450)$, $P_c(4380)$

$$\Lambda_b \longrightarrow K^- + P_c$$

$$\hookrightarrow J/\psi(c\bar{c}) + p(uud)$$

R. Aaij, *et al.*, *Phys. Rev. Lett.* **115**, 072001 (2015)



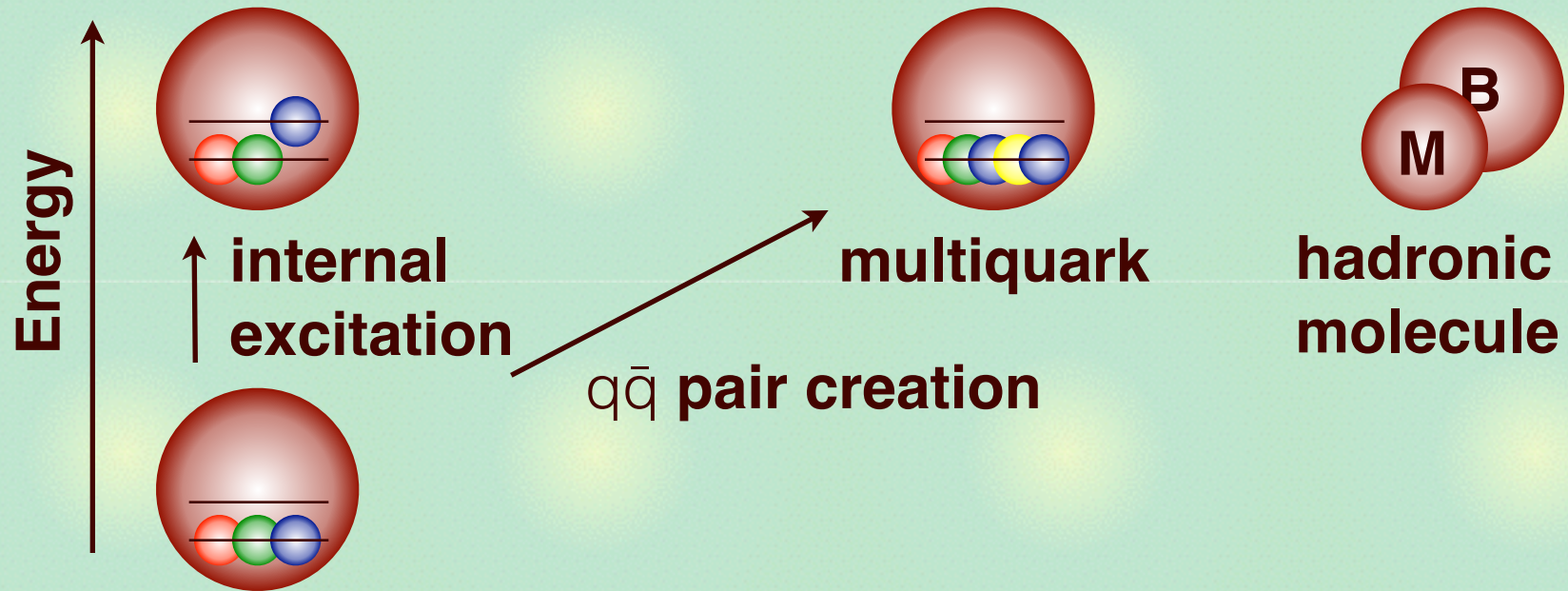
Only a few are observed. **Why only a few?**

Various hadronic excitations

Description of excited baryons

Conventional structure

Exotic structures



In QCD, non- qqq structures naturally arise.

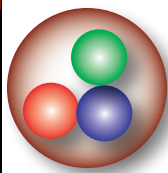
- Baryons: superposition of qqq + exotic structures
- > How can we disentangle different components?

Unstable states via strong interaction

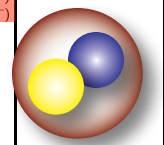
Hadron resonances

PDG2015 : <http://pdg.lbl.gov/>

p	$1/2^+$ ****	$\Delta(1232)$	$3/2^+$ ****	Σ^+	$1/2^+$ ****	Ξ^0	$1/2^+$ ****	Λ_c^+	$1/2^+$ ****
n	$1/2^+$ ****	$\Delta(1600)$	$3/2^+$ ***	Σ^0	$1/2^+$ ****	Ξ^-	$1/2^+$ ****	$\Lambda_c(2595)^+$	$1/2^-$ ***
$N(1440)$	$1/2^+$ ****	$\Delta(1620)$	$1/2^-$ ****	Σ^-	$1/2^+$ ****	$\Xi(1530)$	$3/2^+$ ****	$\Lambda_c(2625)^+$	$3/2^-$ ***
$N(1520)$	$3/2^-$ ****	$\Delta(1700)$	$3/2^-$ ****	$\Sigma(1385)$	$3/2^+$ ****	$\Xi(1620)$	*	$\Lambda_c(2765)^+$	*
$N(1535)$	$1/2^-$ ****	$\Delta(1750)$	$1/2^+$ *	$\Sigma(1480)$	*	$\Xi(1690)$	*	$\Lambda_c(2880)^+$	$5/2^+$ ***
$N(1650)$	$1/2^-$ ****	$\Delta(1900)$	$1/2^-$ **	$\Sigma(1560)$	**	$\Xi(1820)$	$3/2^-$ ***	$\Lambda_c(2940)^+$	***
$N(1675)$	$5/2^-$ ****	$\Delta(1905)$	$5/2^+$ ****	$\Sigma(1580)$	$3/2^-$ **	$\Xi(1950)$	*	$\Sigma_c(2455)$	$1/2^+$ ****
$N(1680)$	$5/2^+$ ****	$\Delta(1910)$	$1/2^+$ ****	$\Sigma(1620)$	$1/2^-$ *	$\Xi(2030)$	$\geq 5/2^?$ ***	$\Sigma_c(2520)$	$3/2^+$ ****
$N(1685)$	*	$\Delta(1920)$	$3/2^+$ ***	$\Sigma(1660)$	$1/2^+$ ***	$\Xi(2120)$	*	$\Sigma_c(2800)$	***
$N(1700)$	$3/2^-$ ***	$\Delta(1930)$	$5/2^-$ **	$\Sigma(1670)$	$3/2^-$ ****	$\Xi(2250)$	**	Ξ_c	$1/2^+$ ***
$N(1710)$	$1/2^+$ ***	$\Delta(1940)$	$3/2^-$ **	$\Sigma(1690)$	***	$\Xi(2370)$	**	Ξ_c	$1/2^+$ ***
$N(1720)$	$3/2^+$ ****	$\Delta(1950)$	$7/2^+$ ****	$\Sigma(1730)$	$3/2^+$ *	$\Xi(2500)$	*	Ξ_c	$1/2^+$ ***
$N(1860)$	$5/2^+$ **	$\Delta(2000)$	$5/2^+$ **	$\Sigma(1750)$	$1/2^-$ ***			Ξ_c	$1/2^+$ ***
$N(1875)$	$3/2^-$ ***	$\Delta(2150)$	$1/2^-$ *	$\Sigma(1770)$	$1/2^+$ *	Ω^-	$3/2^+$ ****	Ξ_c	$3/2^+$ ****
$N(1880)$	$1/2^+$ **	$\Delta(2200)$	$7/2^-$ *	$\Sigma(1775)$	$5/2^-$ ****	$\Omega(2250)$	***	Ξ_c	$3/2^+$ ****
$N(1895)$	$1/2^-$ **	$\Delta(2300)$	$9/2^+$ **	$\Sigma(1840)$	$3/2^+$ **	$\Omega(2380)$	**	Ξ_c	$1/2^-$ ****
$N(1900)$	$3/2^+$ ***	$\Delta(2350)$	$5/2^-$ **	$\Sigma(1880)$	$1/2^+$ **	$\Omega(2470)$	**	Ξ_c	$3/2^-$ ****
$N(1990)$	$7/2^+$ **	$\Delta(2390)$	$7/2^+$ *	$\Sigma(1900)$	$1/2^-$ *			Ξ_c	$2/2^+$ **
$N(2000)$	$5/2^+$ **	$\Delta(2400)$	$9/2^-$ **	$\Sigma(1915)$	$5/2^+$ ****			Ξ_c	$3/2^+$ **
$N(2040)$	$3/2^+$ *	$\Delta(2420)$	$11/2^+$ ****	$\Sigma(1940)$	$3/2^+$ *			Ξ_c	$3/2^+$ **
$N(2060)$	$5/2^-$ **	$\Delta(2750)$	$13/2^-$ **	$\Sigma(1940)$	$3/2^-$ ***			Ξ_c	$3/2^+$ **
$N(2100)$	$1/2^+$ *	$\Delta(2950)$	$15/2^+$ **	$\Sigma(2000)$	$1/2^-$ *			Ξ_c	$3/2^+$ **
$N(2120)$	$3/2^-$ **			$\Sigma(2030)$	$7/2^+$ ****			Ω_c^0	$1/2^+$ ***
$N(2190)$	$7/2^-$ ****	Λ	$1/2^+$ ****	$\Sigma(2070)$	$5/2^+$ *			$\Omega_c(2770)^0$	$3/2^+$ ***
$N(2220)$	$9/2^+$ ****	$\Lambda(1405)$	$1/2^-$ ****	$\Sigma(2080)$	$3/2^+$ **			Ξ_{cc}	*
$N(2250)$	$9/2^-$ ****	$\Lambda(1520)$	$3/2^-$ ****	$\Sigma(2100)$	$7/2^-$ **			Λ_b^0	$1/2^+$ ***
$N(2300)$	$1/2^+$ **	$\Lambda(1600)$	$1/2^+$ ***	$\Sigma(2250)$	***			$\Lambda_b(5912)^0$	$1/2^-$ ***
$N(2570)$	$5/2^-$ **	$\Lambda(1670)$	$1/2^-$ ****	$\Sigma(2455)$	**			$\Lambda_b(5920)^0$	$3/2^-$ ***
$N(2600)$	$11/2^-$ ***	$\Lambda(1690)$	$3/2^-$ ****	$\Sigma(2620)$	**			Σ_b	$1/2^+$ ***
$N(2700)$	$13/2^+$ **	$\Lambda(1710)$	$1/2^+$ *	$\Sigma(3000)$	*			Σ_b^+	$3/2^+$ ***
		$\Lambda(1800)$	$1/2^-$ ***	$\Sigma(3170)$	*			Ξ_b^0, Ξ_b^-	$1/2^+$ ***
		$\Lambda(1810)$	$1/2^+$ ***					$\Xi_b', \Xi_b(5935)$	$1/2^+$ ***
		$\Lambda(1820)$	$5/2^+$ ****					$\Xi_b(5945)^0$	$3/2^+$ ***
		$\Lambda(1830)$	$5/2^-$ ****					$\Xi_b(5955)$	$3/2^+$ ***
		$\Lambda(1890)$	$3/2^+$ ****					Ω_b	$1/2^+$ ***
		$\Lambda(2000)$	*						
		$\Lambda(2020)$	$7/2^+$ *						
		$\Lambda(2050)$	$3/2^-$ *						
		$\Lambda(2100)$	$7/2^-$ ****						
		$\Lambda(2110)$	$5/2^+$ ***						
		$\Lambda(2325)$	$3/2^-$ *						
		$\Lambda(2350)$	$9/2^+$ ***						
		$\Lambda(2585)$	**						



LIGHT UNFLAVORED (S = C = B = 0)		STRANGE (S = ±1, C = B = 0)		CHARMED, STRANGE (C = S = ±1)	$c\bar{c}$ $F_c(F_c)$		
$F_c(F_c)$	$F_c(F_c)$	$F_c(F_c)$	$F_c(F_c)$	$F_c(F_c)$	$F_c(F_c)$		
π^\pm	$1^-(0^-)$	$\rho(1680)$	$0^-(1^-)$	K^{*2}	$1/2(0^-)$	$\eta_c(1S)$	$0^+(0^+)$
π^0	$1^-(0^-)$	$\rho(1690)$	$1^+(3^-)$	K^{*0}	$1/2(0^-)$	$J/\psi(1S)$	$0^-(0^+)$
η	$0^+(0^+)$	$\rho(1700)$	$1^+(1^-)$	K^{*0}	$1/2(0^-)$	$\chi_{c0}(1P)$	$0^+(0^+)$
$\eta(770)$	$1^+(1^-)$	$\omega(1710)$	$0^+(0^+)$	K^{*0}	$1/2(0^-)$	$\chi_{c1}(1P)$	$0^+(1^+)$
$\omega(782)$	$0^-(0^+)$	$\omega(1760)$	$0^+(0^+)$	K^{*0}	$1/2(0^-)$	$\eta_c(1P)$	$0^+(1^+)$
$\eta(958)$	$0^+(0^+)$	$\omega(1800)$	$1^-(0^+)$	$K_1^*(800)$	$1/2(0^-)$	$\chi_{c2}(1P)$	$0^+(2^+)$
$\eta(980)$	$0^+(0^+)$	$\eta(1810)$	$0^+(2^+)$	$K^*(892)$	$1/2(1^-)$	$\eta_c(2S)$	$0^+(0^+)$
$\omega(980)$	$0^+(0^+)$	$\eta(1820)$	$1^-(0^+)$	$K_1^*(1270)$	$1/2(1^+)$	$\psi(2S)$	$0^-(1^-)$
$\omega(980)$	$0^+(0^+)$	$\eta(1835)$	$??(2^+)$	$K_1^*(1400)$	$1/2(1^+)$	$\psi(3770)$	$0^-(1^-)$
$\omega(1020)$	$0^-(1^-)$	$X(1840)$	$??(2^?)$	$K^*(1410)$	$1/2(1^-)$	$X(3823)$	$?^?(2^-)$
$h_1(1170)$	$0^-(1^-)$	$\omega(1850)$	$0^-(3^-)$	$K_0^*(1430)$	$1/2(0^+)$	$X(3872)$	$0^+(1^+)$
$b_1(1235)$	$1^-(1^+)$	$\eta_2(1870)$	$0^+(2^+)$	$K_2^*(1430)$	$1/2(2^+)$	$X(3900)^+$	$?^?(1^+)$
$\omega(1260)$	$1^-(1^+)$	$\eta_2(1880)$	$1^-(2^+)$	$K(1460)$	$1/2(0^-)$	$X(3900)^0$	$?^?(2^+)$
$f_1(1270)$	$0^+(2^+)$	$\rho(1900)$	$1^+(1^-)$	$K_0^*(1580)$	$1/2(2^-)$	B^{*+}	$1/2(0^-)$
$f_1(1285)$	$0^+(1^+)$	$f_0(1910)$	$0^+(2^+)$	$K_1^*(1580)$	$1/2(2^-)$	B^0	$1/2(0^-)$
$\eta(1295)$	$0^+(0^+)$	$f_0(1950)$	$0^+(2^+)$	$K(1630)$	$1/2(2^-)$	B^{*+}/B^0	ADMIXTURE
$f_1(1320)$	$1^-(0^+)$	$\rho_3(1990)$	$1^+(3^-)$	$K_1^*(1650)$	$1/2(2^+)$	$B^{*+}/B^0/B^0/b$	ADMIXTURE
$\omega(1320)$	$1^-(2^+)$	$f_0(2010)$	$0^+(2^+)$	$K_1^*(1680)$	$1/2(1^-)$	$B^{*+}/B^0/B^0/b$	ADMIXTURE
$h_1(1370)$	$0^+(0^+)$	$f_0(2020)$	$0^+(0^+)$	$K_2^*(1770)$	$1/2(2^+)$	V_{cb} and V_{cb}	CKM Matrix Elements
$h_1(1380)$	$1^-(1^+)$	$\rho_3(2050)$	$1^-(4^+)$	$K_3^*(1780)$	$1/2(3^-)$	B^{*+}	$1/2(1^-)$
$\pi_1(1400)$	$1^-(1^+)$	$\omega(2040)$	$1^-(4^+)$	$K_0^*(1830)$	$1/2(0^-)$	$B_1(5721)^+$	$1/2(1^+)$
$\eta(1405)$	$0^+(0^+)$	$\rho_2(2100)$	$1^-(2^+)$	$K_1^*(1950)$	$1/2(0^+)$	$B_1(5721)^0$	$1/2(1^+)$
$f_1(1420)$	$0^+(1^+)$	$f_0(2100)$	$0^+(0^+)$	$K_2^*(1980)$	$1/2(2^+)$	$B_1(5732)$	$?^?(2^+)$
$\omega(1420)$	$0^-(1^-)$	$f_2(2150)$	$0^+(2^+)$	$K_1^*(2045)$	$1/2(4^+)$	$B_2^*(5747)^+$	$1/2(2^+)$
$f_2(1430)$	$0^+(2^+)$	$\rho(2150)$	$1^-(1^-)$	$K_2^*(2250)$	$1/2(2^+)$	$B_2^*(5747)^0$	$1/2(2^+)$
$a_1(1450)$	$1^-(0^+)$	$\omega(2170)$	$0^-(1^-)$	$K_3^*(2380)$	$1/2(3^+)$	$B(5970)^+$	$?^?(2^+)$
$\rho(1450)$	$1^+(1^-)$	$f_0(2200)$	$0^+(0^+)$	$K_3^*(2380)$	$1/2(3^+)$	$B(5970)^0$	$?^?(2^+)$
$\eta(1475)$	$0^+(0^+)$	$f_2(2220)$	$0^+(2^+)$	$K_4^*(2500)$	$1/2(4^-)$	$B(5970)^+$	$?^?(2^+)$
$f_0(1500)$	$0^+(0^+)$	$\eta(2225)$	$0^+(0^+)$	$K(3100)$	$?^?(2^?)$		
$f_1(1510)$	$0^+(1^+)$	$\rho_3(2250)$	$1^+(3^-)$				
$f_2(1525)$	$0^+(2^+)$	$f_0(2300)$	$0^+(2^+)$				
$f_1(1565)$	$0^+(2^+)$	$f_0(2300)$	$0^+(4^+)$				
$\omega(1570)$	$1^-(1^+)$	$f_0(2330)$	$0^+(0^+)$				
$h_1(1595)$	$0^+(1^+)$	$f_2(2340)$	$0^+(2^+)$				
$\pi_1(1600)$	$1^-(1^+)$	$\rho_3(2350)$	$1^+(5^-)$				
$a_1(1640)$	$1^-(1^+)$	$a_1(2450)$	$1^-(6^+)$				
$f_2(1640)$	$0^+(2^+)$	$f_0(2510)$	$0^+(6^+)$				
$\eta_2(1645)$	$0^+(2^+)$						
$\omega(1650)$	$0^-(1^-)$						
$\omega_3(1670)$	$0^-(3^-)$						
$\rho_2(1670)$	$1^-(2^+)$						



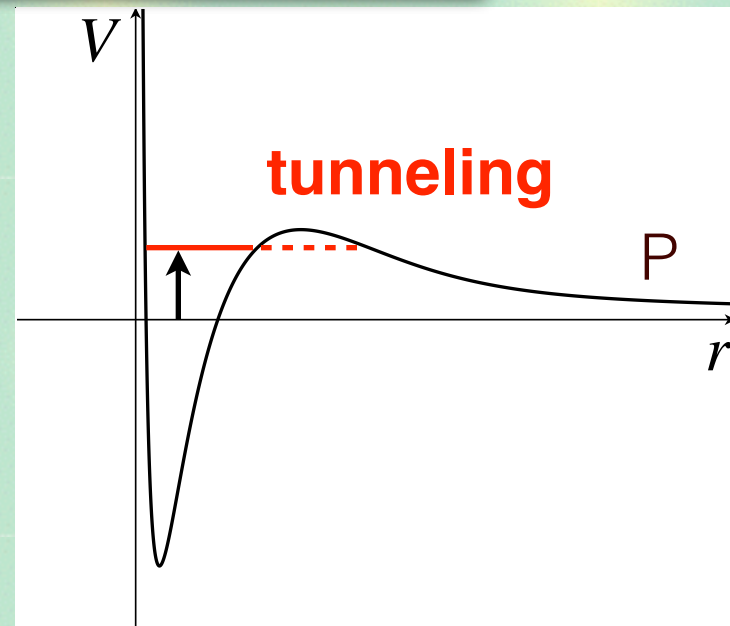
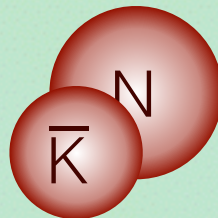
- stable/unstable via strong interaction

- Excited states are mostly unstable. —> resonances

Resonances in quantum mechanics

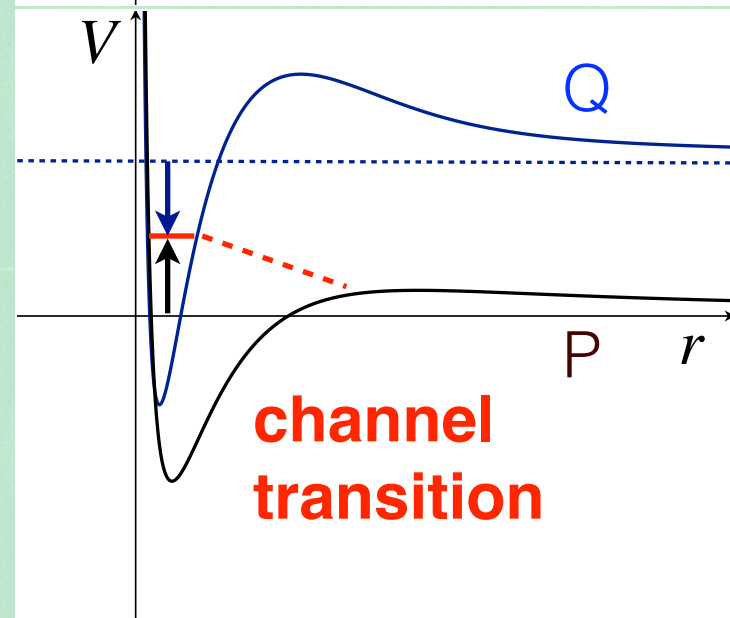
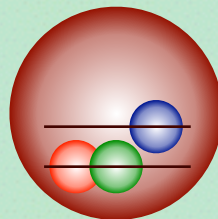
1) Potential (shape) resonance

- 1 channel (P)
- potential barrier : $E > 0$
- unstable via tunneling
- (**composite** of P-channel)



2) Feshbach resonance

- coupled-channel (P+Q)
- bound state of Q: $E_Q < 0$, $E_P > 0$
- unstable via transition
- (**“elementary”**: other than P)



Difficulty of resonances

Resonance as an “eigenstate” of Hamiltonian

- complex energy

G. Gamow, Z. Phys. 51, 204 (1928)

Zur Quantentheorie des Atomkernes.

Von G. Gamow, z. Zt. in Göttingen.

Mit 5 Abbildungen. (Eingegangen am 2. August 1928.)

Um diese Schwierigkeit zu überwinden, müssen wir annehmen, daß die Schwingungen gedämpft sind, und E komplex setzen:

$$E = E_0 + i \frac{\hbar \lambda}{4\pi},$$

wo E_0 die gewöhnliche Energie ist und λ das Dämpfungsdekrement (Zerfallskonstante). Dann sehen wir aber aus den Relationen (2a) und (2b),

- Wave function cannot be normalized.

$$\langle R | R \rangle = \int d\mathbf{r} |\psi_R(\mathbf{r})|^2 \sim \int_0^\infty dr e^{-2\text{Im}[k]r} \rightarrow \infty$$

Normalization by bi-orthogonal basis (Gamow vector)

N. Hokkyo, Prog. Theor. Phys. 33, 1116 (1965)

T. Berggren, Nucl. Phys. A 109, 265 (1968)

$$|\tilde{R}\rangle = |R^*\rangle, \quad \langle \tilde{R} | R \rangle = \int d\mathbf{r} [\psi_R(\mathbf{r})]^2 < \infty$$

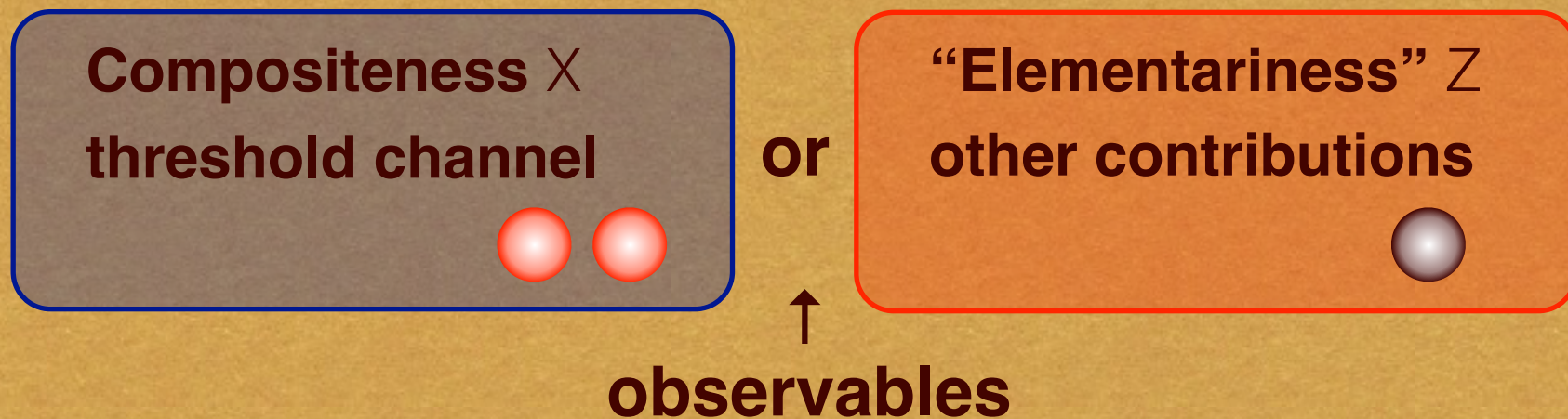
- Complex expectation value (e.g. $\langle r^2 \rangle$) \rightarrow interpretation?

Compositeness of hadrons

Internal structure of excited hadrons

Weak binding relation for stable bound states

S. Weinberg, Phys. Rev. 137, B672 (1965)



Study **compositeness of hadron resonances** by **generalizing** weak binding relation **to unstable states** with effective field theory.

Weak binding relation for stable states

Compositeness of s-wave **weakly bound** state ($R \gg R_{\text{typ}}$)

S. Weinberg, Phys. Rev. 137, B672 (1965);

T. Hyodo, Int. J. Mod. Phys. A 28, 1330045 (2013)

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\}, \quad r_e = R \left\{ \frac{X-1}{X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\}$$

a_0 : **scattering length**, r_e : **effective range**

$R = (2\mu B)^{-1/2}$: **radius of wave function**

R_{typ} : **length scale of interaction**

- **deuteron is NN composite** ($a_0 \sim R \gg r_e$) $\rightarrow X \sim 1$

- **internal structure from **observable****

- **no nuclear force potential / wavefunction of deuteron**

Problem: applicable only for stable states.

Effective field theory

Low-energy scattering with near-threshold bound state

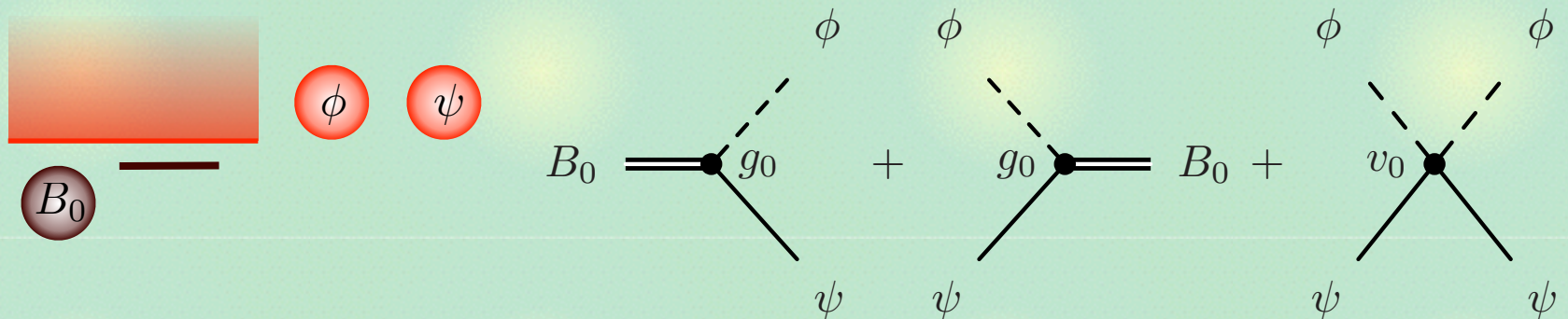
- Nonrelativistic EFT with **contact interaction**

D.B. Kaplan, Nucl. Phys. B494, 471 (1997)

E. Braaten, M. Kusunoki, D. Zhang, Annals Phys. 323, 1770 (2008)

$$H_{\text{free}} = \int d\mathbf{r} \left[\frac{1}{2M} \nabla \psi^\dagger \cdot \nabla \psi + \frac{1}{2m} \nabla \phi^\dagger \cdot \nabla \phi + \frac{1}{2M_0} \nabla B_0^\dagger \cdot \nabla B_0 + \nu_0 B_0^\dagger B_0 \right],$$

$$H_{\text{int}} = \int d\mathbf{r} \left[g_0 \left(B_0^\dagger \phi \psi + \psi^\dagger \phi^\dagger B_0 \right) + v_0 \psi^\dagger \phi^\dagger \phi \psi \right]$$



- **cutoff** : $\Lambda \sim 1/R_{\text{typ}}$ (interaction range of microscopic theory)

- At low energy $p \ll \Lambda$, interaction \sim contact

Compositeness and “elementariness”

Eigenstates

$$H_{\text{free}}|B_0\rangle = \nu_0|B_0\rangle, \quad H_{\text{free}}|\mathbf{p}\rangle = \frac{p^2}{2\mu}|\mathbf{p}\rangle \quad \text{free (discrete + continuum)}$$

$$(H_{\text{free}} + H_{\text{int}})|B\rangle = -B|B\rangle \quad \text{full (bound state)}$$

- normalization of $|B\rangle$ + completeness relation

$$\langle B|B\rangle = 1, \quad 1 = |B_0\rangle\langle B_0| + \int \frac{d\mathbf{p}}{(2\pi)^3} |\mathbf{p}\rangle\langle \mathbf{p}|$$

- projections onto bare states

$$1 = Z + X, \quad \underline{Z \equiv |\langle B_0|B\rangle|^2}, \quad \underline{X \equiv \int \frac{d\mathbf{p}}{(2\pi)^3} |\langle \mathbf{p}|B\rangle|^2}$$

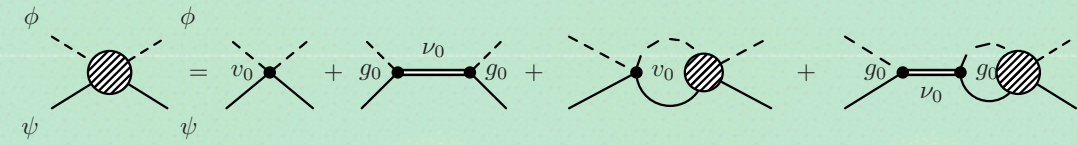
“elementariness” compositeness



Z, X : real and nonnegative \rightarrow interpreted as **probability**

Weak binding relation

$\Psi\Phi$ scattering amplitude (exact result)

$$f(E) = -\frac{\mu}{2\pi} \frac{1}{[v(E)]^{-1} - G(E)}$$


The diagrams show the expansion of the scattering amplitude. The first diagram is a shaded circle with two incoming ψ lines and two outgoing ϕ lines. This is equal to the sum of four diagrams: 1) a vertex v_0 with two ψ lines and two ϕ lines; 2) a loop with two ψ lines and two ϕ lines, with a shaded circle in the middle; 3) a loop with two ψ lines and two ϕ lines, with a shaded circle in the middle and a v_0 vertex on the ψ line; 4) a loop with two ψ lines and two ϕ lines, with a shaded circle in the middle and a v_0 vertex on the ϕ line.

$$v(E) = v_0 + \frac{g_0^2}{E - \nu_0}, \quad G(E) = \frac{1}{2\pi^2} \int_0^\Lambda dp \frac{p^2}{E - p^2/(2\mu) + i0^+}$$

Compositeness $X \leftarrow v(E), G(E)$: renormalization dependent

T. Sekihara, T. Hyodo, D. Jido, PTEP2015, 063D04 (2015)

T. Hyodo, arXiv:1511.00870 [hep-ph]

$$X = \{1 + G^2(-B)v'(-B) [G'(-B)]^{-1}\}^{-1}$$

1/R expansion of scattering length: leading term $\leftarrow X$

$$a_0 = -f(E=0) = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\}$$

renormalization independent

If $R \gg R_{\text{typ}}$, correction terms neglected: $X \leftarrow (B, a_0)$

Generalization to unstable states

Introduce decay channel

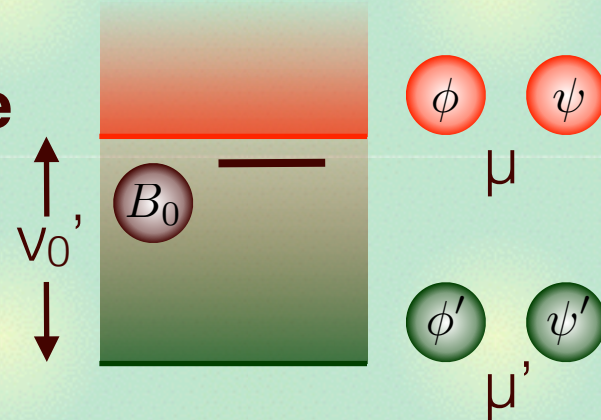
$$H'_{\text{free}} = \int d\mathbf{r} \left[\frac{1}{2M'} \nabla \psi'^{\dagger} \cdot \nabla \psi' - \nu_{\psi} \psi'^{\dagger} \psi' + \frac{1}{2m'} \nabla \phi'^{\dagger} \cdot \nabla \phi' - \nu_{\phi} \phi'^{\dagger} \phi' \right],$$

$$H'_{\text{int}} = \int d\mathbf{r} \left[g'_0 \left(B_0^{\dagger} \phi' \psi' + \psi'^{\dagger} \phi'^{\dagger} B_0 \right) + v'_0 \psi'^{\dagger} \phi'^{\dagger} \phi' \psi' + v_0^t (\psi^{\dagger} \phi^{\dagger} \phi' \psi' + \psi'^{\dagger} \phi'^{\dagger} \phi \psi) \right],$$

Quasi-bound state: complex eigenvalue

$$H = H_{\text{free}} + H'_{\text{free}} + H_{\text{int}} + H'_{\text{int}}$$

$$H |QB\rangle = E_{QB} |QB\rangle, \quad E_{QB} \in \mathbb{C}$$



Generalized relation: **correction term** ← threshold difference

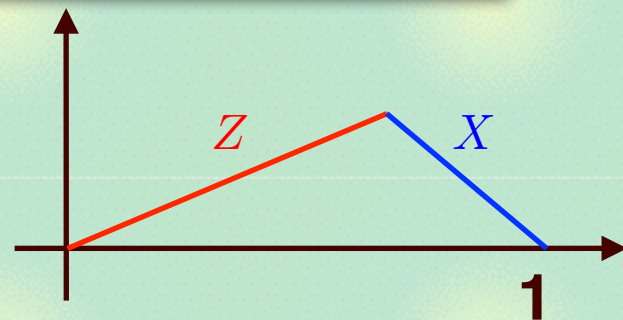
$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O} \left(\left| \frac{R_{\text{typ}}}{R} \right| \right) + \sqrt{\frac{\mu'^3}{\mu^3}} \mathcal{O} \left(\left| \frac{l}{R} \right|^3 \right) \right\}, \quad R = \frac{1}{\sqrt{-2\mu E_{QB}}}, \quad l \equiv \frac{1}{\sqrt{2\mu\nu'_0}}$$

If $|R| \gg (R_{\text{typ}}, l)$ correction terms neglected: $X \leftarrow (E_{QB}, a_0)$

Complex compositeness and interpretation

Unstable states \rightarrow complex Z and X

$$Z + X = 1, \quad Z, X \in \mathbb{C}$$

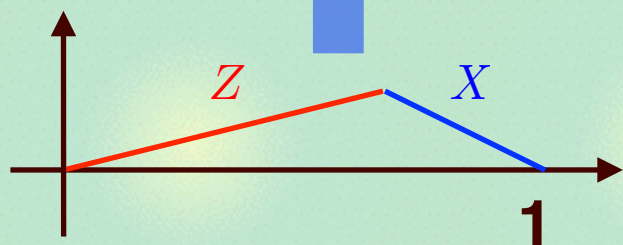
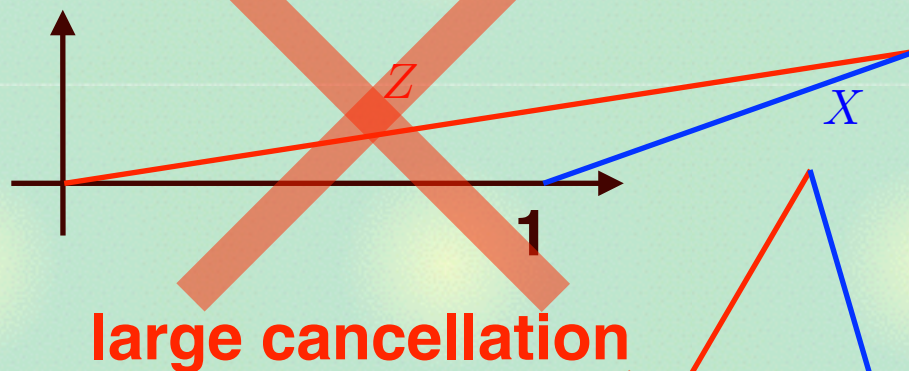


Similarity with bound state

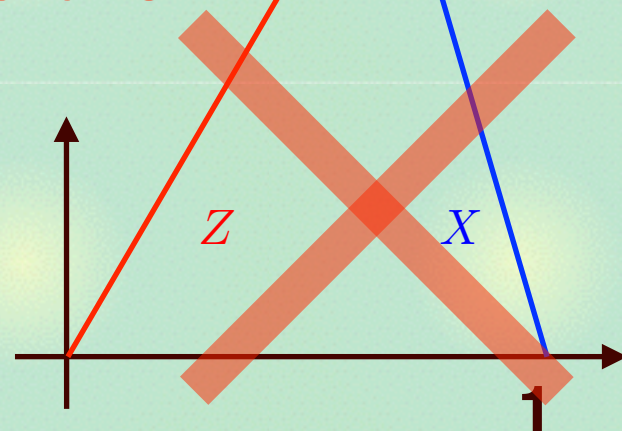
c.f. [T. Sekihara, T. Hyodo, D. Jido, PTEP2015, 063D04 \(2015\)](#)

bound state
: well defined

$$Z + X = 1, \quad Z, X \in [0, 1]$$



small cancellation



New definitions

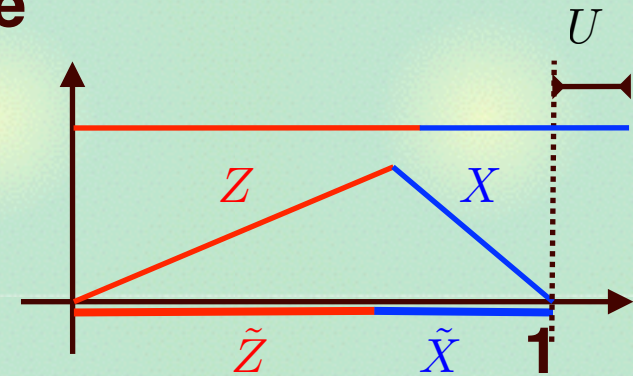
Step 1: quantify the deviation from bound state

- 0 for bound state
- becomes large when deviation is large

$$U = |Z| + |X| - 1$$

→ U : uncertainty of interpretation

c.f. T. Berggren, Phys. Lett. 33B, 547 (1970)



Step 2: define new compositeness/elementariness

- interpreted as probabilities $\tilde{Z} + \tilde{X} = 1$, $\tilde{Z}, \tilde{X} \in [0, 1]$
- coincide with Z, X for bound state if $U \rightarrow 0$

$$\tilde{Z} = \frac{1 - |X| + |Z|}{2}, \quad \tilde{X} = \frac{1 - |Z| + |X|}{2}$$

compositeness when U is small

Application

Generalized weak binding relation $X \leftarrow (E_{QB}, a_0)$

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\left|\frac{R_{\text{typ}}}{R}\right|\right) + \sqrt{\frac{\mu'^3}{\mu^3}} \mathcal{O}\left(\left|\frac{l}{R}\right|^3\right) \right\}, \quad R = \frac{1}{\sqrt{-2\mu E_{QB}}}, \quad l \equiv \frac{1}{\sqrt{2\mu\nu'_0}}$$

- $\Lambda(1405)$ pole position and $\bar{K}N$ scattering length

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881 98 (2012), ...

- $E_{QB} = -10 - 26i$ MeV $\rightarrow |R| \sim 2$ fm \rightarrow small correction term

energy difference from $\pi\Sigma$

$$\left|\frac{R_{\text{typ}}}{R}\right| \lesssim 0.12, \quad \left|\frac{l}{R}\right|^3 \lesssim 0.16$$

vector meson exchange

Ref.	E_{QB} (MeV)	a_0 (fm)	$X_{\bar{K}N}$	$\tilde{X}_{\bar{K}N}$	U	$ r_e/a_0 $
[43]	$-10 - i26$	$1.39 - i0.85$	$1.2 + i0.1$	1.0	0.5	0.2
[44]	$-4 - i8$	$1.81 - i0.92$	$0.6 + i0.1$	0.6	0.0	0.7
[45]	$-13 - i20$	$1.30 - i0.85$	$0.9 - i0.2$	0.9	0.1	0.2
[46]	$2 - i10$	$1.21 - i1.47$	$0.6 + i0.0$	0.6	0.0	0.7
[46]	$-3 - i12$	$1.52 - i1.85$	$1.0 + i0.5$	0.8	0.6	0.4

systematic error

$\Lambda(1405)$ is $\bar{K}N$ composite \leftarrow observables

Summary

- Compositeness of **near-threshold** bound state can be determined only by **observables**.

S. Weinberg, Phys. Rev. 137, B672 (1965)

- Weak binding relation can be **generalized to unstable states** with effective field theory.

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O} \left(\left| \frac{R_{\text{typ}}}{R} \right| \right) + \sqrt{\frac{\mu'^3}{\mu^3}} \mathcal{O} \left(\left| \frac{l}{R} \right|^3 \right) \right\}, \quad R = \frac{1}{\sqrt{-2\mu E_{QB}}}, \quad l \equiv \frac{1}{\sqrt{2\mu\nu'_0}}$$

- Precise determination of the pole position and scattering length shows that $\Lambda(1405)$ is dominated by **$\bar{K}N$ composite component**.

Y. Kamiya, T. Hyodo, arXiv:1509.00146 [hep-ph]

