

Structure and compositeness of hadrons




Tetsuo Hyodo

Yukawa Institute for Theoretical Physics, Kyoto Univ.

2014, Feb. 20th 1

Contents



Introduction: structure of $\Lambda(1405)$

- Comparison of model and data
- **Not a simple issue!**



Compositeness of hadrons

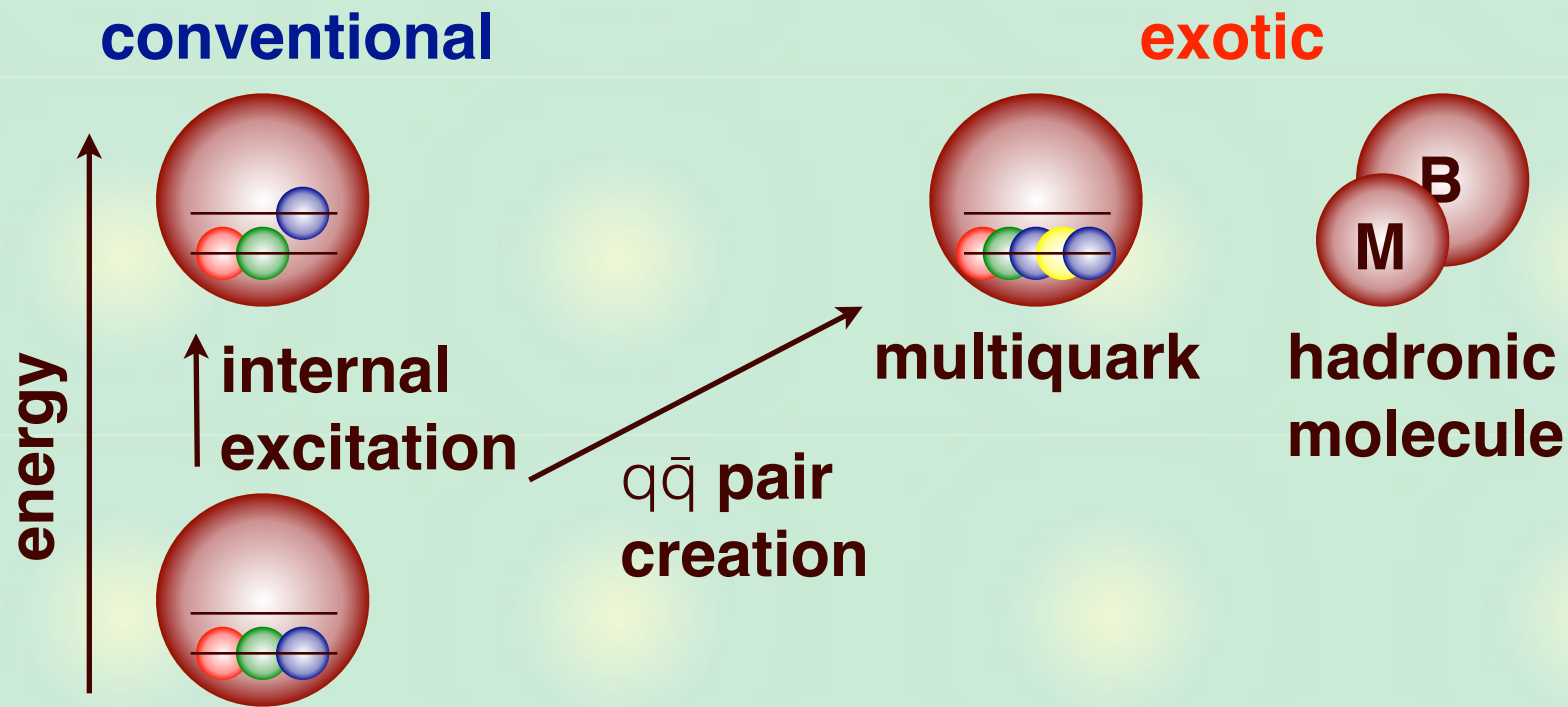
- Field renormalization constant Z
- Negative effective range r_e

[T. Hyodo, Phys. Rev. Lett. 111, 132002 \(2013\)](#)

[T. Hyodo, Int. J. Mod. Phys. A 28, 1330045 \(2013\)](#)

Exotic structure of hadrons

Various excitations of baryons



Physical state: superposition of $3q$, $5q$, MB , ...

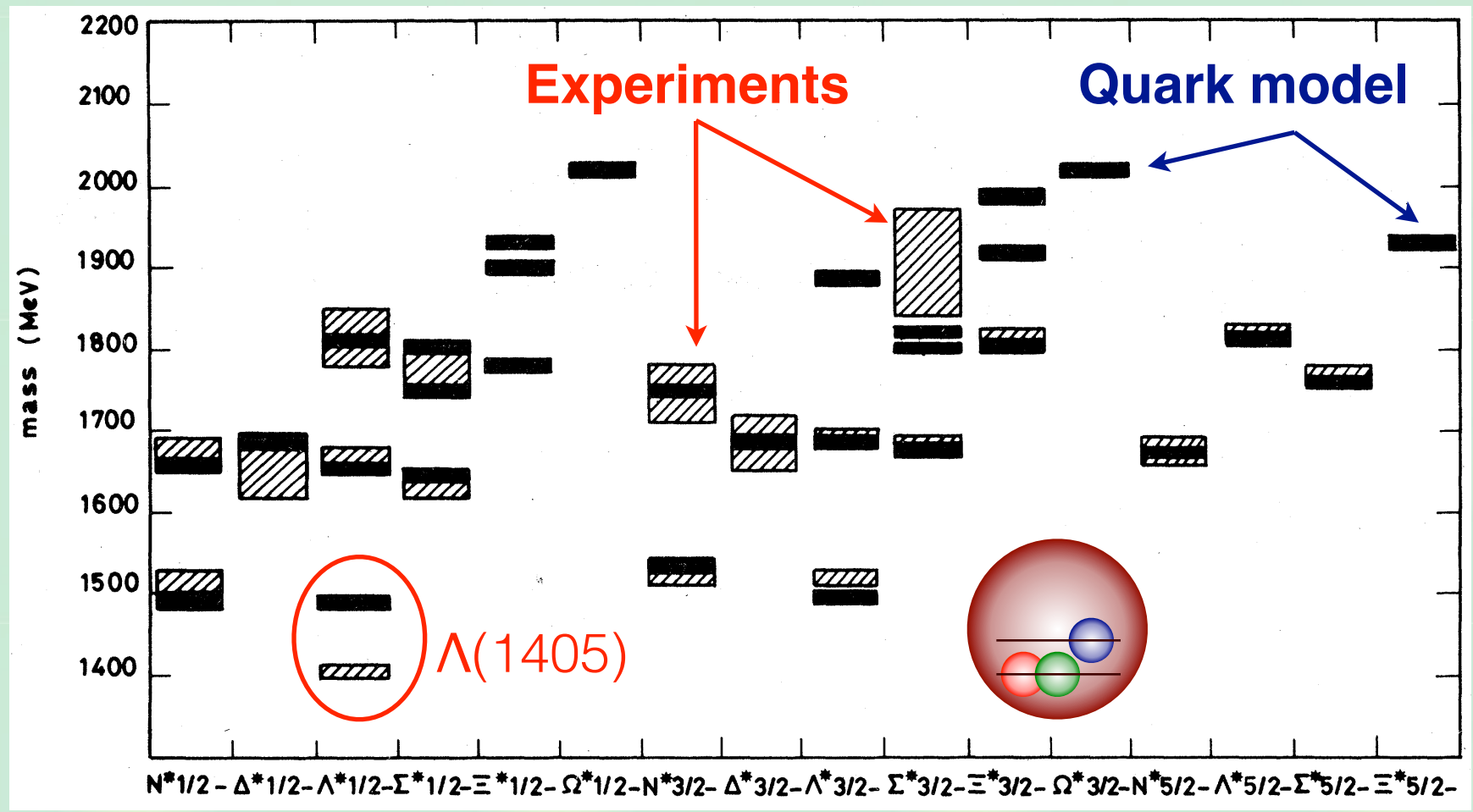
$$|\Lambda(1405)\rangle = \underline{N_{3q}}|uds\rangle + \underline{N_{5q}}|uds\ q\bar{q}\rangle + \underline{N_{\bar{K}N}}|\bar{K}N\rangle + \dots$$

How can we identify the structure of hadrons?

$\Lambda(1405)$ in quark model

Baryon excited states in a constituent quark model

N. Isgur, G. Karl, Phys. Rev. D18, 4187 (1978)



Prediction does not fit experimental data of $\Lambda(1405)$

$\Lambda(1405)$ in hadron molecule model

Dynamical coupled-channel scattering model

R.H. Dalitz, T.C. Wong, G. Rajasekaran, *Phys. Rev.* **153**, 1617 (1967)

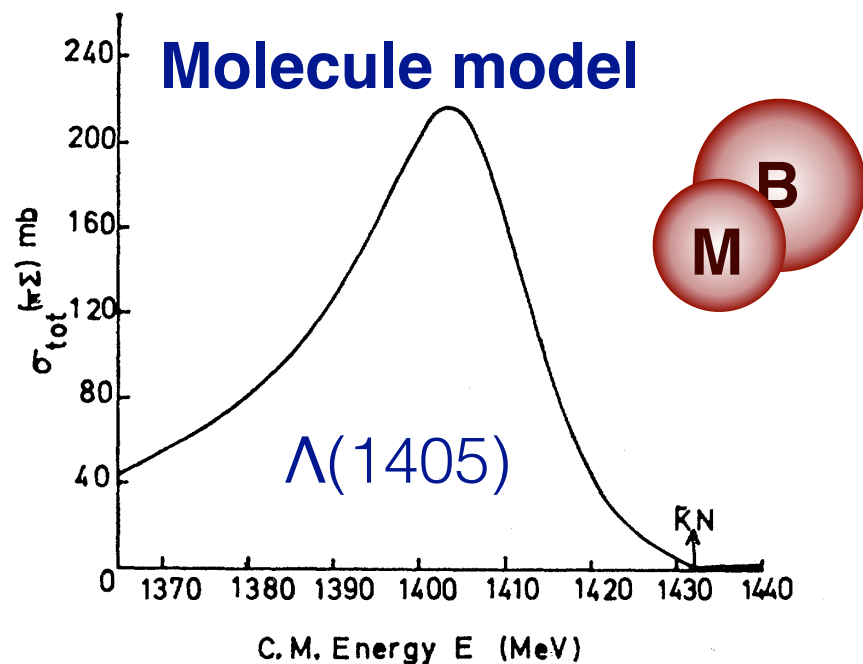
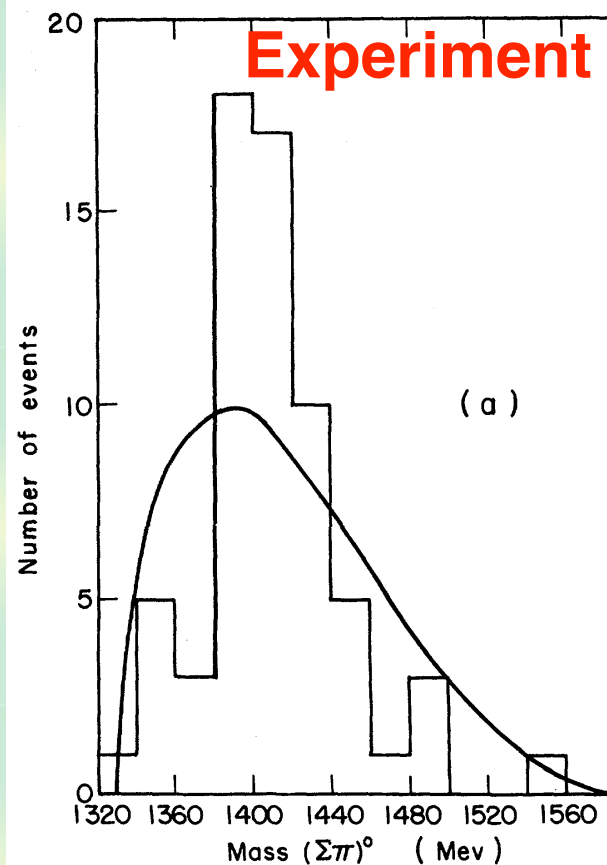


FIG. 1. The total s -wave $\pi\Sigma$ scattering cross section calculated for the multichannel potential model for $Y_0^*(1405)$ is plotted as a function of the total c.m. energy. The cross section becomes very small at the $\bar{K}N$ threshold, where only the term $\gamma(E)$ contributes to the $\pi\Sigma$ scattering.

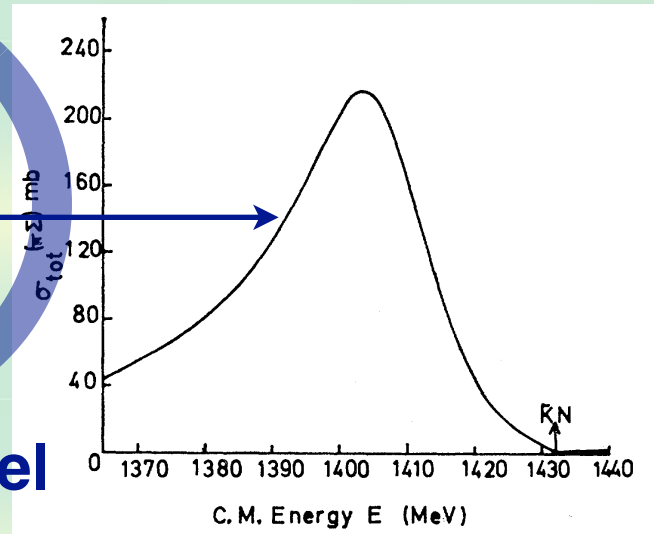
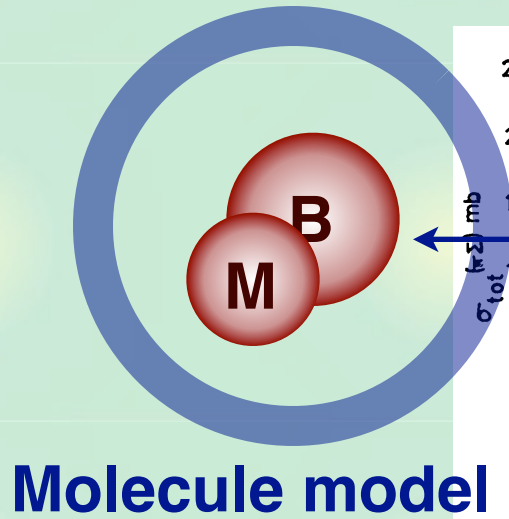
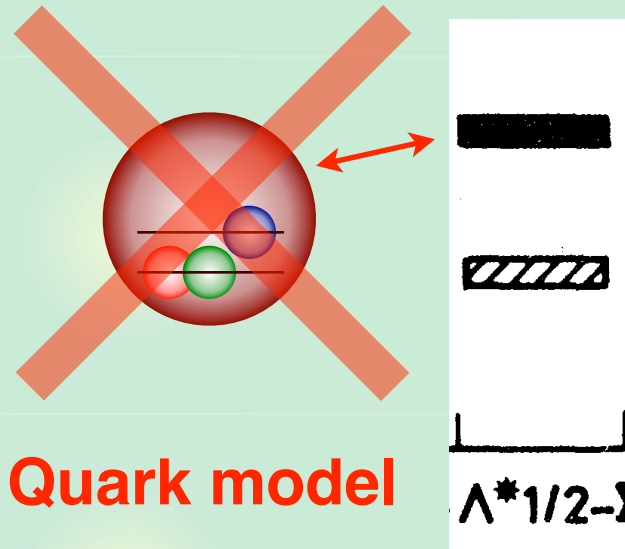


M.H. Alston *et al.*, *Phys. Rev. Lett.* **6**, 698-702 (1961)

Good description of the spectrum (mass and width)

qqq v.s. molecule

Comparison with experimental data



The model prediction contradicts/agrees with data.

(hidden) **assumption:**

- Model space \leftrightarrow structure of the predicted state

Is this so simple?

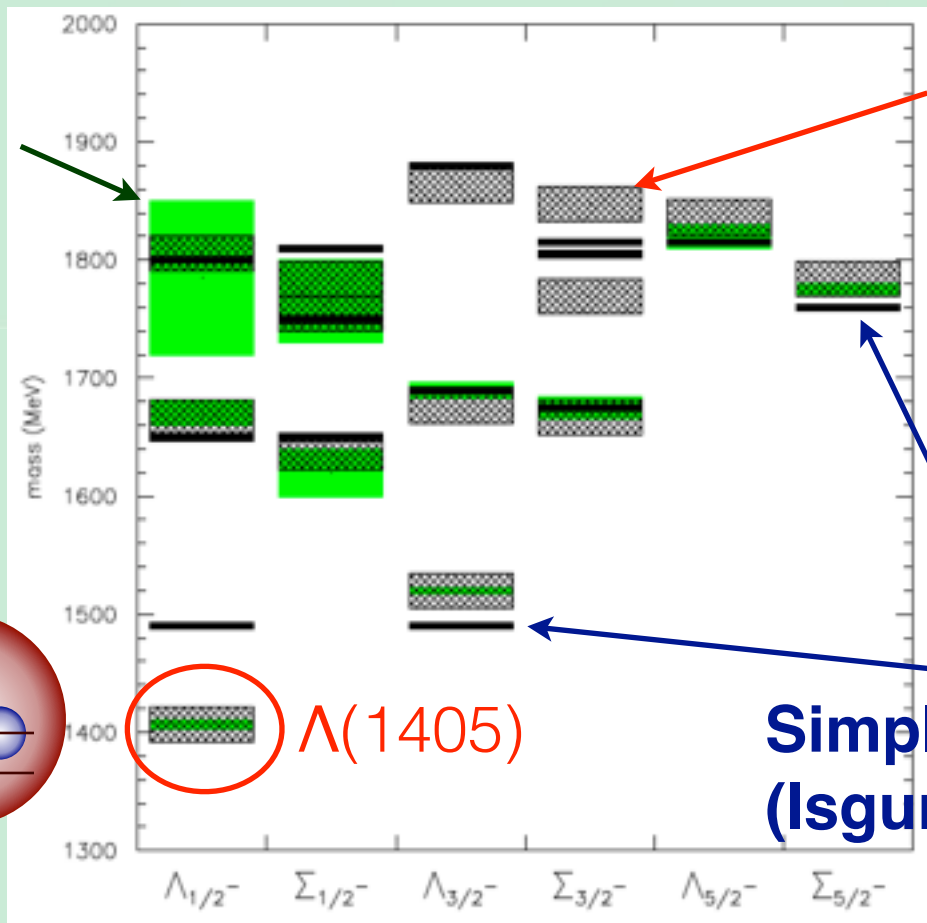
Improvement of models

Quark model with more interactions (large N_c expansion)

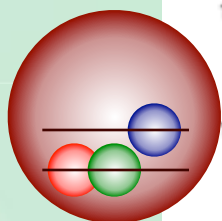
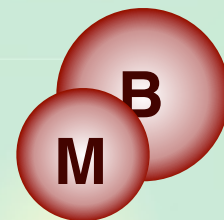
C.L. Schat, J.L. Goity, N.N. Scoccola, Phys. Rev. Lett. 88, 102002 (2002);

J.L. Goity, C.L. Schat, N.N. Scoccola, Phys. Rev. D 66, 114014 (2002)

Experiments



“improved”
quark model
(15 parameters)



Simple quark model
(Isgur-Karl)

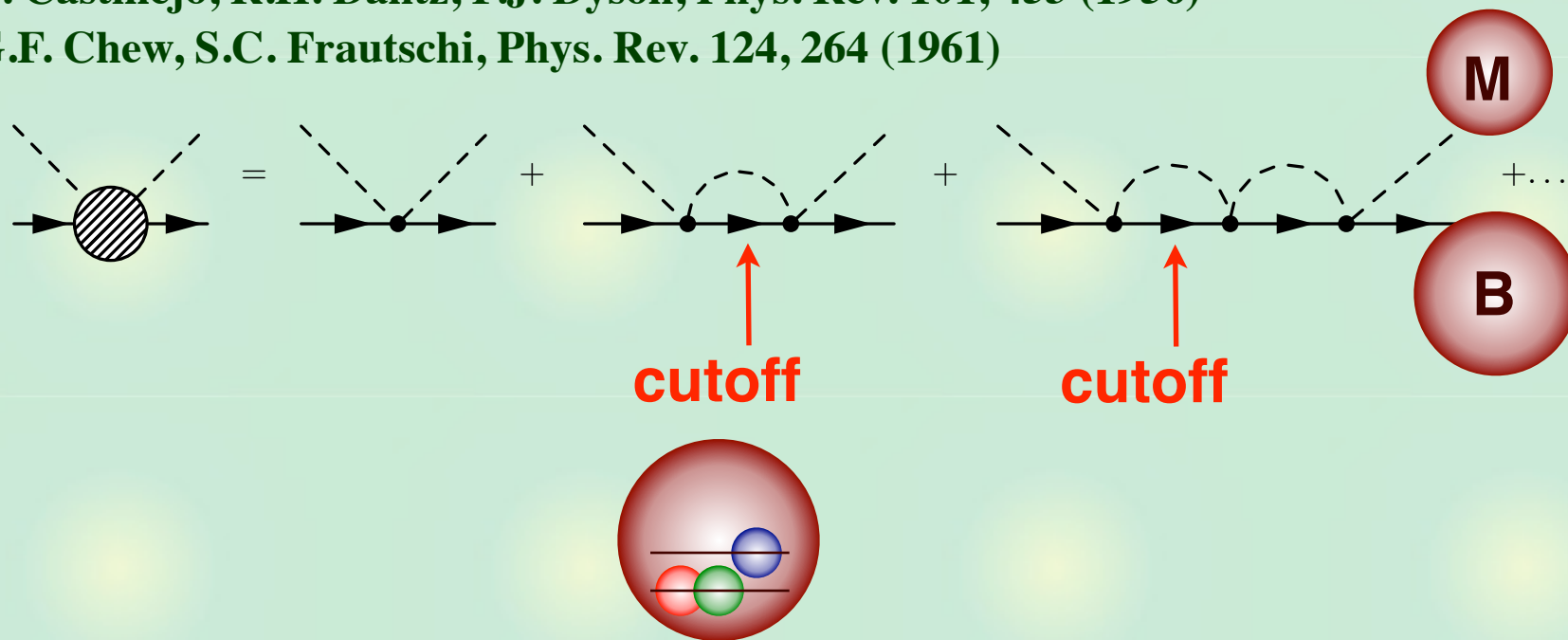
qqq model can reproduce $\Lambda(1405)$??

Ambiguity in the scattering equation

CDD pole contribution

L. Castillejo, R.H. Dalitz, F.J. Dyson, *Phys. Rev.* 101, 453 (1956)

G.F. Chew, S.C. Frautschi, *Phys. Rev.* 124, 264 (1961)

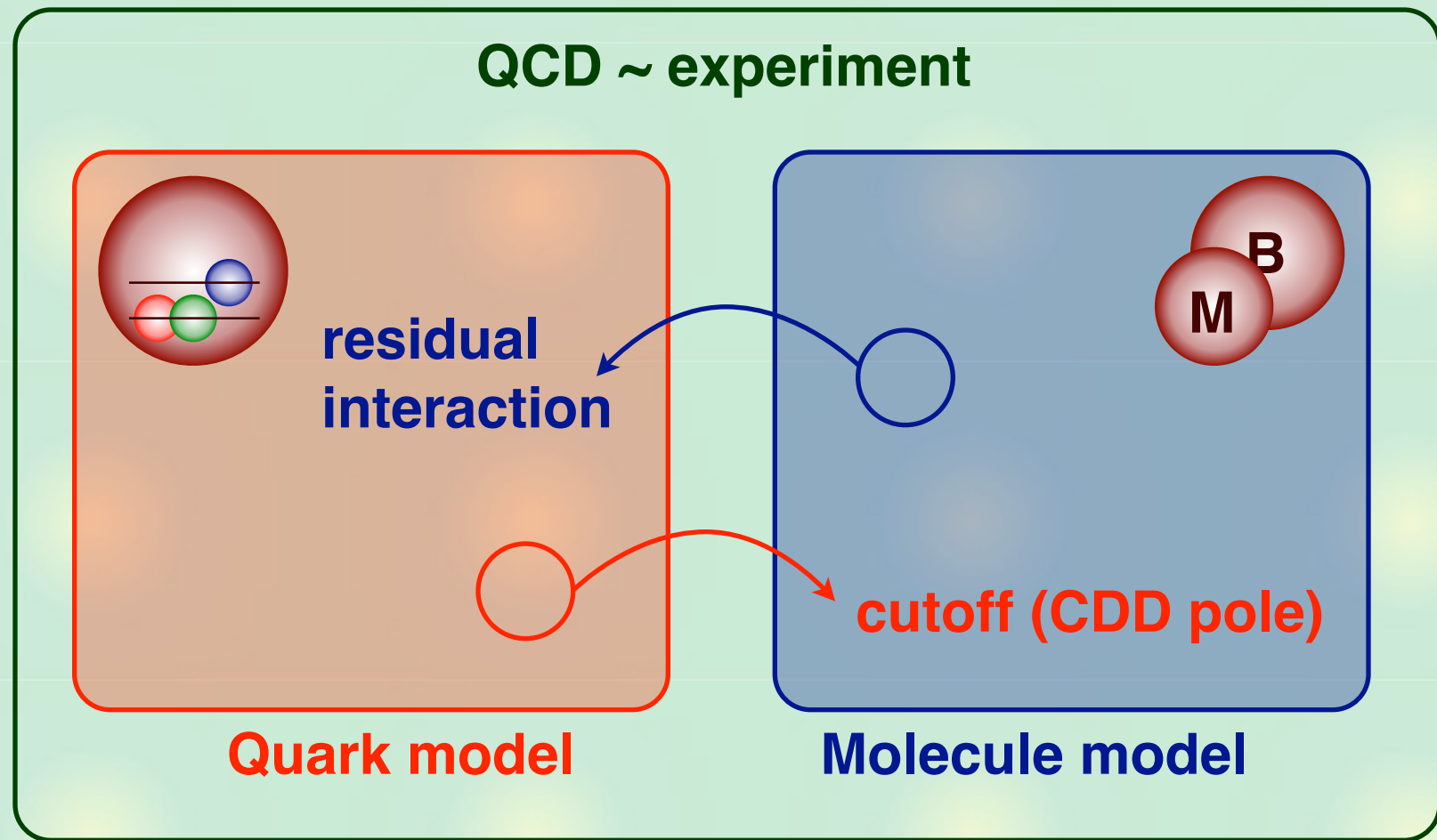


- contributions other than the model space

Parameters in the model effectively encode the effect of the configurations which are not included in the model space.


Ambiguities in the model analysis


Schematic picture:



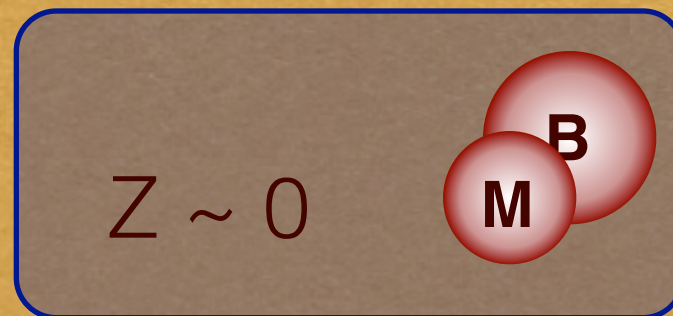
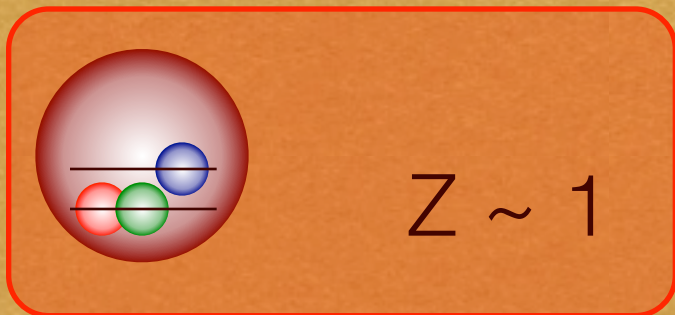
=> model space \neq structure of the predicted state

Summary of introduction

 **Model space \neq structure of hadron**

 **What we need is a **model-independent** measure for the hadron structure.**

QCD \sim experiment



Compositeness of bound states

Compositeness approach for a bound state $|B\rangle$

S. Weinberg, *Phys. Rev.* **137**, B672 (1965); T. Hyodo, *IJMPA* **28**, 1330045 (2013)

$$H = H_0 + V \quad H|B\rangle = -B|B\rangle, \quad \langle B|B\rangle = 1$$

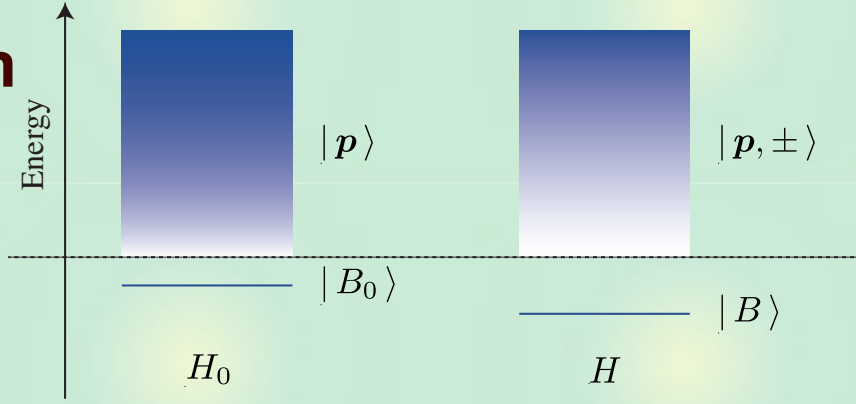
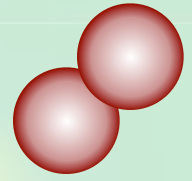
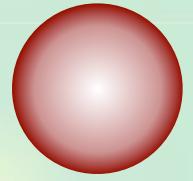
Decompose H into free part + interaction

Complete set for free Hamiltonian : bare $|B_0\rangle$ + continuum

$$1 = |B_0\rangle\langle B_0| + \int dp |p\rangle\langle p|$$

$$1 = \underbrace{\langle B|B_0\rangle\langle B_0|B\rangle}_{Z} + \underbrace{\int dp \langle B|p\rangle\langle p|B\rangle}_{X}$$

Z : elementary **X : composite**



In QCD,

H_0 : free hadrons

V : hadron interaction

Z, X : real and nonnegative --> **probabilistic interpretation**

$$\Rightarrow 0 \leq Z \leq 1, \quad 0 \leq X \leq 1$$

Weak binding limit

In general, Z depends on the choice of the potential V .

- Z : model-(scheme-)dependent quantity

$$1 - Z = \int d\mathbf{p} \frac{|\langle \mathbf{p} | V | B \rangle|^2}{(E_p + B)^2} \longleftarrow \text{V-dependent}$$

At the **weak binding** ($R \gg R_{\text{typ}}$), Z is related to observables.

S. Weinberg, Phys. Rev. 137, B672 (1965); T. Hyodo, IJMPA 28, 1330045 (2013)

$$a = \frac{2(1 - Z)}{2 - Z} R + \mathcal{O}(R_{\text{typ}}), \quad r_e = \frac{-Z}{1 - Z} R + \mathcal{O}(R_{\text{typ}}),$$

a : **scattering length**, r_e : **effective range**

$R = (2\mu B)^{-1/2}$: **radius (binding energy)**

R_{typ} : **typical length scale of the interaction**

Criterion for the structure:

$$\begin{cases} a \sim R_{\text{typ}} \ll -r_e & \text{(elementary dominance), } Z \sim 1 \\ a \sim R \gg r_e \sim R_{\text{typ}} & \text{(composite dominance). } Z \sim 0 \text{ (deuteron)} \end{cases}$$

Interpretation of negative effective range

For $Z > 0$, effective range is always **negative**.

$$a = \frac{2(1-Z)}{2-Z}R + \mathcal{O}(R_{\text{typ}}), \quad r_e = \frac{-Z}{1-Z}R + \mathcal{O}(R_{\text{typ}}),$$

$$\begin{cases} a \sim R_{\text{typ}} \ll -r_e & \text{(elementary dominance),} \\ a \sim R \gg r_e \sim R_{\text{typ}} & \text{(composite dominance).} \end{cases}$$

Simple attraction (no barrier, energy-indep., ...) : $r_e > 0$

--> only “composite dominance” is possible.

$r_e < 0$: energy- (momentum-)dependence of the potential

D. Phillips, S. Beane, T.D. Cohen, *Annals Phys.* 264, 255 (1998)

E. Braaten, M. Kusunoki, D. Zhang, *Annals Phys.* 323, 1770 (2008)

<-- pole term/Feshbach projection of coupled-channel effect

Negative r_e --> **Something other than $|p\rangle$: CDD pole**

Generalization to resonances

Compositeness approach of bound states

$$1 - Z = \int d\mathbf{p} \frac{|\langle \mathbf{p} | V | B \rangle|^2}{(E_p + B)^2}$$

Generalization to **general** resonances in chiral models

T. Hyodo, D. Jido, A. Hosaka, Phys. Rev. C85, 015201 (2012)

F. Aceti, E. Oset, Phys. Rev. D86, 014012 (2012)

$$1 - Z = -g^2 \frac{dG(W)}{dW} \Big|_{W \rightarrow M_B} \quad \rightarrow \quad 1 - Z = -g_{II}^2 \frac{dG_{II}(W)}{dW} \Big|_{\underline{W \rightarrow z_R}}$$

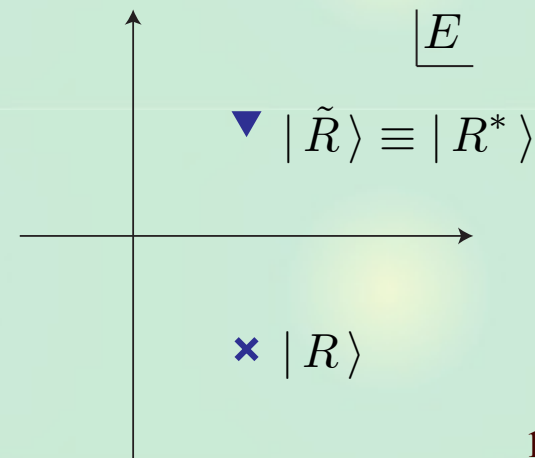
- Z is in general **complex**. Interpretation?

$$\langle R | R \rangle \rightarrow \infty, \quad \langle \tilde{R} | R \rangle = 1$$

$$1 = \underline{\langle \tilde{R} | B_0 \rangle \langle B_0 | R \rangle} + \int d\mathbf{p} \langle \tilde{R} | \mathbf{p} \rangle \langle \mathbf{p} | R \rangle$$

complex

$$\langle \tilde{R} | B_0 \rangle = \langle B_0 | R \rangle \neq \langle B_0 | R \rangle^*$$



Z of hadron resonances

Z can be calculated in chiral models

Table 1. Field renormalization constant Z of the hadron resonances evaluated on the resonance pole. The momentum cutoff q_{\max} is chosen to be 1 GeV for the $\rho(770)$ and $K^*(892)$ mesons,^{55,59} 0.5 GeV for the $\Delta(1232)$ baryon, and 0.45 GeV for the $\Sigma(1385)$, $\Xi(1535)$, Ω baryons.⁶⁰

Baryons	Z	$ Z $	Mesons	Z	$ Z $
$\Lambda(1405)$ higher pole ⁵⁸	$0.00 + 0.09i$	0.09	$f_0(500)$ or σ ⁵⁸	$1.17 - 0.34i$	1.22
$\Lambda(1405)$ lower pole ⁵⁸	$0.86 - 0.40i$	0.95	$f_0(980)$ ⁵⁸	$0.25 + 0.10i$	0.27
$\Delta(1232)$ ⁶⁰	$0.43 + 0.29i$	0.52	$a_0(980)$ ⁵⁸	$0.68 + 0.18i$	0.70
$\Sigma(1385)$ ⁶⁰	$0.74 + 0.19i$	0.77	$\rho(770)$ ⁵⁵	$0.87 + 0.21i$	0.89
$\Xi(1535)$ ⁶⁰	$0.89 + 0.99i$	1.33	$K^*(892)$ ⁵⁹	$0.88 + 0.13i$	0.89
Ω ⁶⁰	0.74	0.74			
$\Lambda_c(2595)$ ⁵⁶	$1.00 - 0.61i$	1.17			

F. Aceti, E. Oset, Phys. Rev. D86, 014012 (2012)

T. Sekihara, T. Hyodo, Phys. Rev. C87, 045202 (2013)

C.W. Xiao, F. Aceti, M. Bayar, Eur. Phys. J. A 49, 22 (2013)

F. Aceti, L. Dai, L. Geng, E. Oset, T. Zhang, arXiv:1301.2554 [hep-ph]

In some cases, Z and/or $|Z|$ exceed unity. Interpretation?

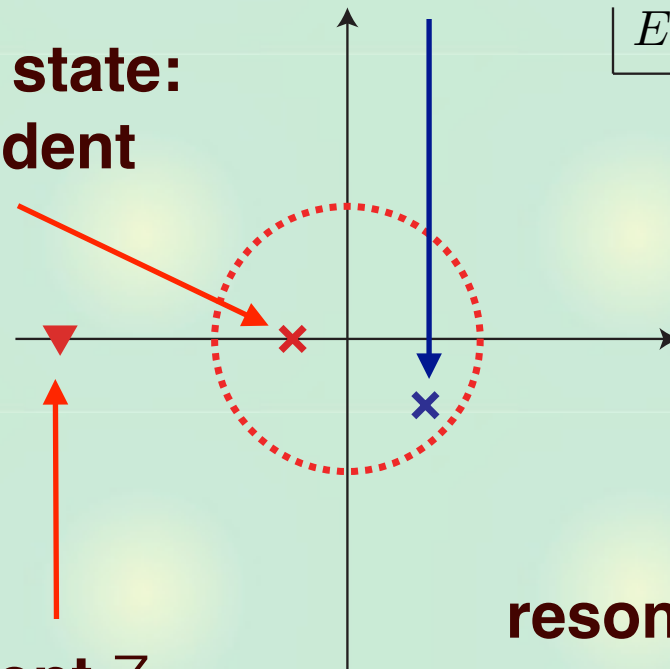
Generalization to resonances

Compositeness approach at the **weak binding**:

- Model-independent (no potential, wavefunction, ...)
- Related to experimental observables

What about **near-threshold resonances** (\sim small binding) ?

shallow bound state:
model-independent
structure



bound state:
model-dependent Z

resonance: model-
dependent complex Z

Poles of the amplitude

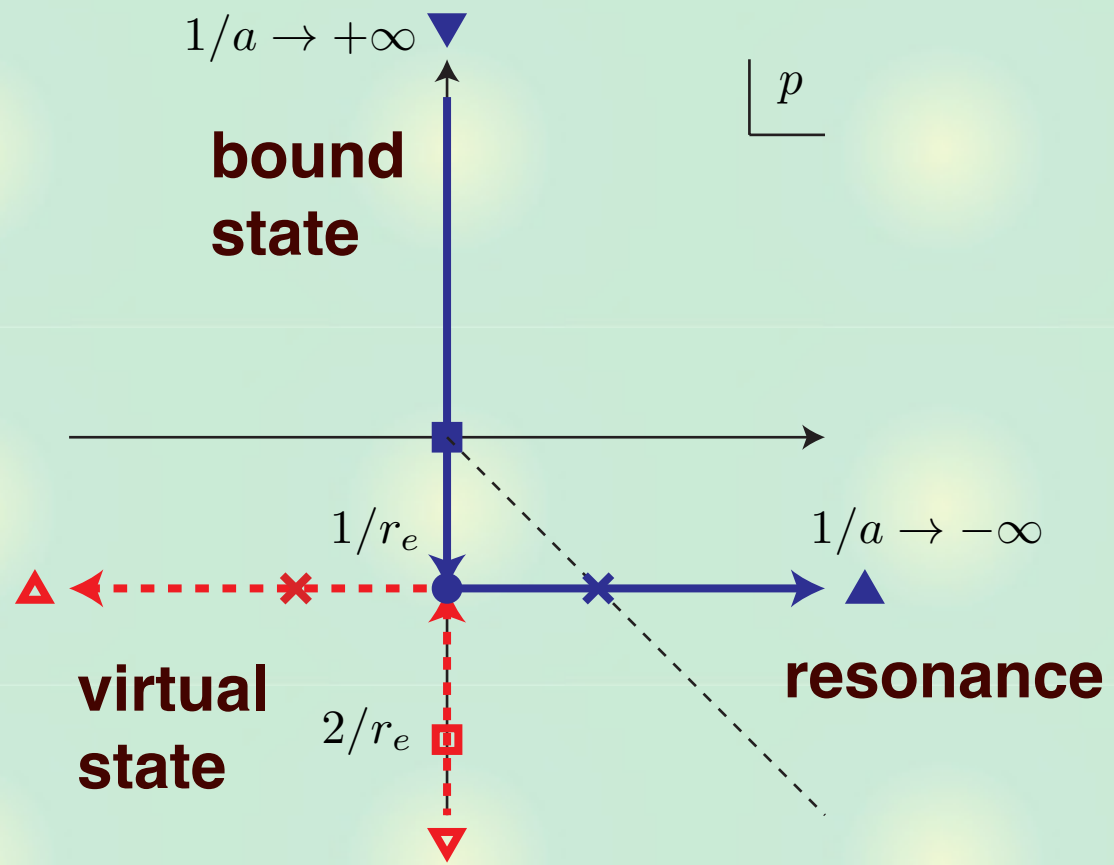
Near-threshold phenomena: effective range expansion

T. Hyodo, Phys. Rev. Lett. 111, 132002 (2013) with opposite sign of scattering length

$$f(p) = \left(-\frac{1}{a} + \frac{r_e^2}{2} p^2 - ip \right)^{-1}$$

$$p^\pm = \frac{i}{r_e} \pm \frac{1}{r_e} \sqrt{\frac{2r_e}{a} - 1}$$

Pole trajectories with a fixed $r_e < 0$



Resonance pole position $\leftrightarrow (a, r_e)$

Example of resonance: $\Lambda_c(2595)$

Pole position of $\Lambda_c(2595)$ in $\pi\Sigma_c$ scattering

- central values in PDG

$$E = 0.67 \text{ MeV}, \quad \Gamma = 2.59 \text{ MeV} \quad p^\pm = \sqrt{2\mu(E \mp i\Gamma/2)}$$

- deduced threshold parameters of $\pi\Sigma_c$ scattering

$$a = -\frac{p^+ + p^-}{ip^+p^-} = -10.5 \text{ fm}, \quad r_e = \frac{2i}{p^+ + p^-} = -19.5 \text{ fm}$$

- field renormalization constant: complex

$$Z = 1 - 0.608i$$





Large negative effective range

←- substantial elementary contribution other than $\pi\Sigma_c$
(three-quark, other meson-baryon channel, or ...)

$\Lambda_c(2595)$ is **not likely a $\pi\Sigma_c$ molecule**

Summary

Composite/elementary nature of resonances

-  Renormalization constant Z measures elementariness of a stable bound state.
-  In general, Z of a resonance is complex.
-  Negative effective range r_e : CDD pole
-  Near-threshold resonance : pole position is related to r_e --> elementariness

[T. Hyodo, Phys. Rev. Lett. 111, 132002 \(2013\)](#)

[T. Hyodo, Int. J. Mod. Phys. A 28, 1330045 \(2013\)](#)