Compositeness of hadron resonances in chiral dynamics





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composite v.s. elementary

- Field renormalization constant Z

- Negative effective range re

<u>T. Hyodo, Phy. Rev. Lett. 111, 132002 (2013)</u> <u>T. Hyodo, arXiv:1310.1176 [hep-ph]</u>



Exotic structure of hadrons

Various excitations of baryons



Physical state: superposition of 3q, 5q, MB, ...

 $|\Lambda(1405)\rangle = N_{3q}|uds\rangle + N_{5q}|uds|q\bar{q}\rangle + N_{\bar{K}N}|\bar{K}N\rangle + \cdots$

Find out the dominant component among others.

Introduction

Structure of resonances?

Excited states : finite width (unstable against strong decay)

- stable (ground) states
- unstable states

Most of hadrons are unstable!

State vector of resonance?

								-					
n	$1/2^{+}$	****	A(1232)	3/2+ ****	Σ+	$1/2^{+}$	****	=0	$1/2^{+}$	****	A+	$1/2^{+}$	****
P D	1/2+	****	A(1600)	3/2+ ***	τ ⁰	1/2+	****	-	1/2+	****	Λ (DEDE)+	1/2	***
N(1440)	1/2	****	A(1620)	1/2 ****	Σ-	1/2+	****	=(1520)	3/2+	****	$\Lambda_{C}(2355)$	2/2-	***
N(1520)	3/2-	****	$\Delta(1700)$	2/2 ****	Σ(1385)	3/2+	****	=(1500) =(1620)	3/2	*	$\Lambda_{C}(2025)^{+}$	3/2	*
N(1520)	1/2-	****	$\Delta(1760)$	1/0+ *	$\Sigma(1400)$	3/2	*	=(1020) =(1600)		***	$\Lambda_{C}(2700)^{+}$	F /o+	
N(16E0)	1/2	****	$\Delta(1750)$	1/2 **	$\Sigma(1400)$		**	=(1090)	2/0-	***	$\Lambda_{c}(2880)^{+}$	5/2 '	***
N(167E)	1/2	****	A(100E)	1/2 ···	$\Sigma(1500)$	2/2-	*	=(1020)	3/2	***	/l _c (2940)'	1.0+	****
N(1075)	5/2	****	$\Delta(1905)$	5/2 1/0+ ****	$\Sigma(1000)$	3/2	**	=(1950)	< 5 ?	***	$\Sigma_{C}(2455)$	1/2	****
N(1080)	5/21	*	∆(1910) A(1000)	2/2+ ***	$\Sigma(1020)$	1/2	444	=(2030)	< <u>2</u>		$\Sigma_{c}(2520)$	3/21	***
N(1000)	2/2-		A(1020)	5/2 ***	$\Sigma(1000)$	2/2-		=(2120)			$\frac{Z_{c}(2800)}{-1}$	1 /0+	444
N(1710)	3/2	***	$\Delta(1930)$	5/2	$\Sigma(1070)$	5/2	**	=(2250)		**	- <u>-</u>	1/2	***
N(1710)	2/2+	****	∠1(1940) A(10E0)	3/2 ***	$\Sigma(1090)$	1/0-	444	=(2370)		** *	= <u>c</u>	1/2	***
N(1720)	3/21	**	∆(1950) A(0000)	1/2 + ++	$\Sigma(1750)$	1/2	*	=(2500)		*	$=_{c}^{\prime+}$	1/2+	***
N(1075)	5/2'	**	$\Delta(2000)$	5/2 ***	$\Sigma(1770)$	1/2 '	- -	0-	2/0+	****	Ξ_c^0	$1/2^{+}$	***
N(1875)	3/2	***	$\Delta(2150)$	1/2 *	$\Sigma(1775)$	5/2	****	<u>M</u>	3/2 '	***	$\Xi_{c}(2645)$	3/2+	***
N(1880)	1/21	**	$\Delta(2200)$	1/2 *	$\Sigma(1840)$	3/2	The second secon	$\Omega(2250)^{-1}$		***	<i>Ξ_c</i> (2790)	$1/2^{-}$	***
N(1895)	1/2-	**	$\Delta(2300)$	9/2 **	Σ(1880)	1/2	**	$\Omega(2380)^{-}$		**	$\Xi_{c}(2815)$	$3/2^{-}$	***
N(1900)	3/21	***	$\Delta(2350)$	5/2 *	$\Sigma(1915)$	5/2 1	****	<u>12(2470)</u> ⁻		**	$\Xi_{c}(2930)$		*
N(1990)	1/21	**	$\Delta(2390)$	(/2 *	$\Sigma(1940)$	3/2	***				$\Xi_{c}(2980)$		***
N(2000)	5/2	**	⊿(2400)	9/2 **	$\Sigma(2000)$	1/2-	*				$\Xi_{c}(3055)$		**
N(2040)	3/2*	*	$\Delta(2420)$	11/2+ ****	$\Sigma(2030)$	1/2+	****				$\Xi_{c}(3080)$		***
N(2060)	5/2	**	$\Delta(2750)$	13/2 **	$\Sigma(2070)$	5/2	*				$\Xi_{c}(3123)$		*
N(2100)	1/2+	*	$\Delta(2950)$	15/2+ **	Σ(2080)	3/2+	**				Ω_c^0	$1/2^{+}$	***
N(2120)	3/2-	**		- (o	$\Sigma(2100)$	7/2-	*				$\Omega_{c}(2770)^{0}$	$3/2^{+}$	***
N(2190)	7/2-	****	/	1/2 ****	Σ(2250)		***						
N(2220)	9/2+	****	A(1405)	1/2 ****	Σ(2455)		**				=+		*
N(2250)	9/2-	****	/(1520)	3/2 ****	Σ(2620)		**						
N(2600)	11/2-	***	Λ(1600)	1/2+ ***	Σ(3000)		*				Λ_{b}^{0}	$1/2^{+}$	***
N(2700)	13/2+	**	Λ(1670)	1/2- ****	Σ(3170)		*				Σ_{b}	$1/2^{+}$	***
			/(1690)	3/2 ****							Σ_{b}^{*}	$3/2^{+}$	***
			A(1800)	1/2 ***							$\Xi_{bl}^{\bar{0}}$, Ξ_{b}^{-}	$1/2^{+}$	***
			A(1810)	1/2+ ***							Ω_{h}^{-}	$1/2^{+}$	***
			Λ(1820)	5/2+ ****							D		
			A(1830)	5/2 ****									
			A(1890)	3/2+ ****									
			Л(2000)	*									
			A(2020)	7/2+ *									
			Λ(2100)	7/2 ****									
			Л(2110)	5/2+ ***									
			A(2325)	3/2 *									7
			Λ(2350)	9/2+ ***									
			A(2585)	**									_

$$|\Lambda(1405)\rangle = N_{3q}|uds\rangle + N_{5q}|uds|q\bar{q}\rangle + N_{\bar{K}N}|\bar{K}N\rangle + \cdots$$

We need a classification scheme applicable to resonances.

Field renormalization constant Z and compositeness

Compositeness of bound states

Compositeness approach for a bound state |B>

- S. Weinberg, Phys. Rev. 137, B672 (1965); <u>T. Hyodo, arXiv:1310.1176 [hep-ph]</u>
 - $H = H_0 + V \qquad H | B \rangle = -B | B \rangle, \quad \langle B | B \rangle = 1$

Decompose H into free part + interaction part



Field renormalization constant Z and compositeness

Weak binding limit

In general, Z depends on the choice of the potential V.

- Z : model-(scheme-)dependent quantity

$$1 - Z = \int d\mathbf{p} \frac{|\langle \mathbf{p} | V | B \rangle|^2}{(E_p + B)^2}$$

At the weak binding ($R \gg R_{typ}$), Z is related to observables.

S. Weinberg, Phys. Rev. 137, B672 (1965); <u>T. Hyodo, arXiv:1310.1176 [hep-ph]</u>

$$a = \frac{2(1-Z)}{2-Z}R + \mathcal{O}(R_{\text{typ}}), \quad r_e = \frac{-Z}{1-Z}R + \mathcal{O}(R_{\text{typ}}),$$

a : scattering length, r_e : effective range $R = (2\mu B)^{-1/2}$: radius (binding energy) R_{typ} : typical length scale of the interaction

Criterion for the structure:

 $\begin{cases} a \sim R_{\text{typ}} \ll -r_e & \text{(elementary dominance)}, \ \mathsf{Z} \sim \mathsf{1} \\ a \sim R \gg r_e \sim R_{\text{typ}} & \text{(composite dominance)}, \ \mathsf{Z} \sim \mathsf{0} \text{ (deuteron)} \end{cases}$

Field renormalization constant Z and compositeness

Interpretation of negative effective range

For Z > 0, **effective range is always negative.**

$$a = \frac{2(1-Z)}{2-Z}R + \mathcal{O}(R_{\text{typ}}), \quad r_e = \frac{-Z}{1-Z}R + \mathcal{O}(R_{\text{typ}}),$$

 $\begin{cases} a \sim R_{\text{typ}} \ll -r_e & \text{(elementary dominance),} \\ a \sim R \gg r_e \sim R_{\text{typ}} & \text{(composite dominance).} \end{cases}$

Simple attractive potential: r_e > 0 --> only "composite dominance" is possible.

$r_e < 0$: energy- (momentum-)dependence of the potential

D. Phillips, S. Beane, T.D. Cohen, Annals Phys. 264, 255 (1998) E. Braaten, M. Kusunoki, D. Zhang, Annals Phys. 323, 1770 (2008)

<-- pole term/Feshbach projection of coupled-channel effect

Negative r_e --> Something other than |p> : CDD pole

Generalization to resonances

Compositeness approach of bound states

$$1 - Z = \int d\mathbf{p} \frac{|\langle \mathbf{p} | V | B \rangle|^2}{(E_p + B)^2}$$

Generalization to general resonances in chiral models

<u>T. Hyodo, D. Jido, A. Hosaka, Phys. Rev. C85, 015201 (2012)</u> F. Aceti, E. Oset, Phys. Rev. D86, 014012 (2012)

$$1 - Z = -g^2 \frac{dG(W)}{dW} \Big|_{W \to M_B} \quad \to \quad 1 - Z = -g_{II}^2 \frac{dG_{II}(W)}{dW} \Big|_{W \to z_R}$$

- Z and X = 1 - Z are in general complex. Interpretation?

*

$$\langle R \,|\, R \rangle \to \infty, \quad \langle \tilde{R} \,|\, R \rangle = 1$$

$$1 = \frac{\langle \tilde{R} \,|\, B_0 \rangle \langle B_0 \,|\, R \rangle}{\mathsf{complex}} + \int d\mathbf{p} \langle \tilde{R} \,|\, \mathbf{p} \rangle \langle \, \mathbf{p} \,|\, R \rangle$$

$$\langle \tilde{R} \,|\, B_0 \rangle = \langle B_0 \,|\, R \rangle \neq \langle B_0 \,|\, R \rangle$$

Z of hadron resonances

Z can be calculated in chiral models

Table 1. Field renormalization constant Z of the hadron resonances evaluated on the resonance pole. The momentum cutoff q_{max} is chosen to be 1 GeV for the $\rho(770)$ and $K^*(892)$ mesons,^{55,59} 0.5 GeV for the $\Delta(1232)$ baryon, and 0.45 GeV for the $\Sigma(1385), \Xi(1535), \Omega$ baryons.⁶⁰

Baryons	Ζ	Z	Mesons	Ζ	Z
	$\begin{array}{c} 0.00+0.09i\\ 0.86-0.40i\\ 0.43+0.29i\\ 0.74+0.19i\\ 0.89+0.99i\\ 0.74\\ 1.00-0.61i\end{array}$	$\begin{array}{c} 0.09 \\ 0.95 \\ 0.52 \\ 0.77 \\ 1.33 \\ 0.74 \\ 1.17 \end{array}$	$egin{aligned} f_0(500) & ext{or} & \sigma^{58} \ f_0(980)^{58} \ a_0(980)^{58} \ ho(770)^{55} \ K^*(892)^{59} \end{aligned}$	$\begin{array}{c} 1.17-0.34i\\ 0.25+0.10i\\ 0.68+0.18i\\ 0.87+0.21i\\ 0.88+0.13i\end{array}$	$ \begin{array}{c} 1.22 \\ 0.27 \\ 0.70 \\ 0.89 \\ 0.89 \end{array} $

F. Aceti, E. Oset, Phys. Rev. D86, 014012 (2012) <u>T. Sekihara, T. Hyodo, Phys. Rev. C87, 045202 (2013)</u> C.W. Xiao, F. Aceti, M. Bayar, Eur. Phys. J. A 49, 22 (2013) F. Aceti, L. Dai, L. Cong, F. Oset, T. Zhang, arXiv:1301 2554 [box

F. Aceti, L. Dai, L. Geng, E. Oset, T, Zhang, arXiv:1301.2554 [hep-ph]

In some cases, Z and/or |Z| exceed unity. Interpretation?

Generalization to resonances

Compositeness approach at the weak binding:

- Model-independent (no potential, wavefunction, ...)
- Related to experimental observables
- Only for bound states with small binding

What about near-threshold resonances (~ small binding) ?



Poles of the amplitude

Near-threshold phenomena: effective range expansion

T. Hyodo, Phy. Rev. Lett. 111, 132002 (2013) with opposite sign of scattering length



Resonance pole position <--> (a, r_e**)**

Example of resonance: $\Lambda_c(2595)$

- Pole position of $\Lambda_c(2595)$ in $\pi\Sigma_c$ scattering
 - central values in PDG

 $E = 0.67 \text{ MeV}, \quad \Gamma = 2.59 \text{ MeV}$

- deduced threshold parameters

$$a = -\frac{p^+ + p^-}{ip^+p^-} = -10.5 \text{ fm}, \quad r_e = \frac{2i}{p^+ + p^-} = -19.5 \text{ fm}$$

- field renormalization constant: complex
 - Z = 1 0.608i

Large negative effective range

<-- substantial elementary contribution other than $\pi\Sigma_c$ (three-quark, other meson-baryon channel, or ...)

 $\Lambda_c(2595)$ is not likely a $\pi\Sigma_c$ molecule

Summary



Composite/elementary nature of resonances

Renormalization constant Z measures elementariness of a stable bound state.

 \checkmark In general, Z of a resonance is complex.

Negative effective range r_e: CDD pole

Near-threshold resonance : pole position is related to r_e --> elementariness

> <u>T. Hyodo, Phy. Rev. Lett. 111, 132002 (2013)</u> <u>T. Hyodo, arXiv:1310.1176 [hep-ph]</u>