## D-meson-nucleon interaction

## and DNN systems



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## Contents

## Introduction <br> $+$

- Variational calculation with DN potential
- FCA to Faddeev equation


## Summary <br>  .

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## DN interaction and $\Lambda_{c}(2595)$

## DNN quasi-bound state <br> DNN quasi-bound state



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Introduction

## Conventions for heavy mesons

Convention of quantum number of quarks

| strange | charm | bottom |
| :--- | :--- | :--- |
| $\mathrm{S}=-1$ | $\mathrm{C}=+1$ | $\mathrm{~B}=-1$ |

Heavy-light mesons: bar for negative flavor-ness (q~u,d)

| with $\overline{\mathrm{q}}$ | $\overline{\mathrm{K}}(\mathrm{s} \overline{\mathrm{q}})$ | $\mathrm{D}(\mathrm{c} \overline{\mathrm{q}})$ | $\overline{\mathrm{B}}(\mathrm{b} \overline{\mathrm{q}})$ |
| :--- | :--- | :--- | :--- |
| with q | $\mathrm{K}(\overline{\mathrm{s} q})$ | $\overline{\mathrm{D}}(\overline{\mathrm{q} q})$ | $\mathrm{B}(\overline{\mathrm{b} q})$ |

DN <--> K̄N : non-exotic light quark annihilation

$\overline{\mathrm{D}} \mathrm{N}$ <--> KN : exotic
$\mathbf{O}^{+}$, Ikeda's talk

## Why DN and DNN?

Comparison with $\bar{K} N$ system in $\mathrm{I}=0$ channel


- large mass splitting between DN and $\boldsymbol{\pi} \boldsymbol{\Sigma}_{\mathrm{c}}$
- narrow negative parity $\Lambda_{c}{ }^{*}$, analogous to $\Lambda(1405)$ ?
$\Lambda^{*}$ : a $\bar{K} N$ bound state in the $\Pi \Sigma$ continuum $-->\bar{K}$ nuclei $\Lambda_{c}{ }^{*}$ : a DN bound state in the $\boldsymbol{\pi} \boldsymbol{\Sigma}_{\mathrm{c}}$ continuum $-->\mathrm{D}$ nuclei?

DN interaction and $\Lambda_{c}(2595)$

## DN bound state picture?

Can $\Lambda_{c}{ }^{*}$ (with large binding) be a DN quasi-bound state?

- $D(1867 \mathrm{MeV})$ is heavier than $\overline{\mathrm{K}}$ ( 496 MeV ). Kinetic energy is suppressed. If the DN interaction were the same with $\bar{K} N$, system would develop a deeper quasi-bound state.
- Vector meson exchange picture leads to a stronger DN interaction than $\overline{\mathrm{K}} \mathrm{N}$ at threshold

$$
\frac{V_{D}}{V_{K}}=\frac{m_{D}}{m_{K}} \sim 3.8 \quad \text { (next slide) }
$$

DN system can generate a strongly bound state: $\Lambda_{c}{ }^{*}$.

## Vector meson exchange for DN

DN ( $\bar{K} N$ ) interaction in vector meson exchange (low energy)

$$
\begin{aligned}
V & \sim g \bar{u} \gamma^{\mu} u \times \frac{1}{k^{2}-m_{v}^{2}}\left[g_{\mu \nu}-\frac{k_{\mu} k_{\nu}}{m_{v}^{2}}\right] \times g\left(q+q^{\prime}\right)^{\nu} \\
& \rightarrow-\bar{u} \gamma^{\mu} u \frac{g^{2}}{m_{v}^{2}} g_{\mu \nu}\left(q+q^{\prime}\right)^{\nu} \quad\left(k \ll m_{v}\right) \\
& \rightarrow-\frac{1}{2 f^{2}} \bar{u}\left(q+\not q^{\prime}\right) u \quad \text { (KSRF relation) (Weinl } \\
& \rightarrow-\frac{1}{2 f^{2}}\left(q^{0}+q^{0 \prime}\right) \quad \text { (nonrel. leading) } \\
& \rightarrow-\frac{m}{f^{2}} \quad(\text { at threshold) }
\end{aligned}
$$

Interaction in DN- $\boldsymbol{\Pi} \boldsymbol{\Sigma}_{\mathrm{c}}$ system

$$
V \sim\left(\begin{array}{cc}
-3 m_{D} & \sqrt{\frac{3}{\frac{3}{2} \kappa_{c}} \frac{m_{D}+m_{\pi}}{2}} \\
\sqrt{\frac{\sqrt[3]{2} \hbar_{c} m^{m}+m_{\pi}}{2}} & -4 m_{\pi}
\end{array}\right)
$$

$$
\kappa_{c} \sim \frac{m_{K^{*}}^{2}}{m_{D^{*}}^{2}} \sim \frac{1}{4}
$$

- strong DN interaction --> large binding energy
- suppressed off-diagonal coupling --> narrow width of $\Lambda_{c}{ }^{*}$

DN interaction and $\Lambda_{c}(2595)$

## DN scattering amplitude

Coupled-channel DN ( $\left.\boldsymbol{\Pi} \Sigma_{c}, \eta \Lambda_{c}, K \Xi_{c}, K \Xi_{c}{ }^{\prime}, D_{s} \Lambda, \eta{ }^{\prime} \Lambda_{c}\right)$ scattering see T. Mizutani, A. Ramos, Phys. Rev. C74, 065201 (2006)

Subtraction constants (cutoff parameters) are chosen to reproduce $\Lambda_{c}{ }^{*}$ in $\mathrm{I}=0$. Apply the same constants to $\mathrm{I}=1$.



A resonance at $\sim 2760 \mathrm{MeV}$ is generated in l=1 channel. c.f. PDG $1^{*}$ : $\Lambda_{c}{ }^{*}(2765)$ or $\Sigma_{c}{ }^{*}(2765)$ ??

DN interaction and $\Lambda_{c}(2595)$

## DN local potential

Equivalent single-channel local potential see T. Hyodo, W. Weise, Phys. Rev. C77, 035204 (2008)

$$
v_{D N}(r ; W)=\frac{M_{N}}{2 \pi^{3 / 2} a_{s}^{3} \tilde{\omega}(W)}\left[v^{\mathrm{eff}}(W)+\Delta v(W)\right] \exp \left[-\left(r / a_{s}\right)^{2}\right]
$$

- reproduces the coupled channel amplitude




This potential reproduces the DN amplitude in CC model. Larger (smaller) real (imaginary) part than $\bar{K} N$

## Strategy for DNN bound state

Coupled-channel model DN amplitude, $\Lambda_{c}(2595)$

DN singlechannel potential
real part

Three-body variational calculation

- Structure from wave function
- NN dynamics is dynamically solved. Coupled-channel $\left(\pi Y_{c} N\right)$ effect is partly included.
- Two-body absorption
- Imaginary part of the amplitude is treated.


## Fixed-center approximation to Faddeev equation

## Variational calculation: setup

Quantum number: $\mathrm{I}=1 / 2, \mathrm{JP}^{\mathrm{P}}=\mathrm{O}^{-}, 1^{-}$

- JP=0-"D+nn"

$$
\begin{aligned}
& S_{\mathrm{NN}}=0 \\
& \mathrm{I}_{\mathrm{NN}}=1 \text { (s-wave) }- \text {-> DN(I=0):DN(I=1) = 3:1 }
\end{aligned}
$$

- JP=1-"D+d"

$$
\begin{aligned}
& S_{\text {NN }}=1 \\
& I_{\text {NN }}=0 \text { (s-wave) }-->\text { DN(I=0):DN(I=1) = } 1: 3
\end{aligned}
$$

Two-body interactions

- DN imaginary part is neglected
- energy dependence is fixed at $\Lambda_{c}{ }^{*}$ ( $l=1$ QBS disappears)
- three kinds of NN forces (Av18, HN1R, Minnesota)


## DNN quasi-bound state

## Variational calculation: results

## Results of the DNN system

- J=0 bound, J=1 unbound w.r.t. [DN]N
- mesonic decay width is small
- softer the core, larger the binding

|  | HN1R |  | Minnesota | Av18 |
| :--- | ---: | ---: | ---: | ---: |
|  | $J=1$ | $J=0$ | $J=0$ | $J=0$ |
|  | unbound | bound | bound | bound |
| $B$ | 208 | 225 | 251 | 209 |
| $M_{B}$ | 3537 | 3520 | 3494 | 3536 |
| $\Gamma_{\pi Y_{C} N}$ | - | 26 | 38 | 22 |
| $E_{\text {kin }}$ | 338 | 352 | 438 | 335 |
| $V(N N)$ | 0 | -2 | 19 | -5 |
| $V(D N)$ | -546 | -575 | -708 | -540 |
| $T_{\text {nuc }}$ | 113 | 126 | 162 | 117 |
| $E_{N N}$ | 113 | 124 | 181 | 113 |
| $P(\mathrm{Odd})$ | $75.0 \%$ | $14.4 \%$ | $7.4 \%$ | $18.9 \%$ |



## Variational calculation: DN correlation

Isospin decomposition of DN two-body correlation

$$
\rho_{D N}(x)=\langle\Psi| \sum_{i=1,2} \delta^{3}\left(\left|\boldsymbol{r}_{D}-\boldsymbol{r}_{i}\right|-x\right)|\Psi\rangle
$$

DN $(\mathrm{l}=0)$ correlation is similar to $\Lambda_{c}{ }^{*}$

DNN quasi-bound state

## FCA calculation

Fixed-center approximation to Faddeev equation


- Complex DN amplitude
- all two-body pairs are in s-wave
- NN distribution is assumed
(chosen to be smaller than the deuteron)


## DNN quasi-bound state

## FCA calculation: two-body absorption

Two-body absorption --> imaginary part of DN amplitude


$$
g_{D N} \rightarrow g_{D N}+i \operatorname{Im} \delta \tilde{g}
$$

DN loop
two-body absorption contribution


## FCA calculation: result

Magnitude of the three-body amplitude square



J=0 channel: M ~ 3500 MeV

- strong signal, consistent with the variational calculation

J=1 channel: M ~ 3500 MeV and M ~ 3700 MeV ?

- week signal, not found in the variational calculation??
- I=1 DN interaction is important for this channel.


## Possible experiments

Antiproton beam

$$
\bar{p}+{ }^{3} \mathrm{He} \rightarrow \bar{D}^{0} D^{0} p n \rightarrow \bar{D}^{0}[D N N]
$$

- PANDA?

Pion beam

$$
\begin{aligned}
& \pi^{-}+d \rightarrow D^{-} D^{+} n p \rightarrow D^{-}[D N N] \\
& \pi^{-}+d \rightarrow D^{-} \Lambda_{c}^{+} n \rightarrow D^{-}[D N N]
\end{aligned}
$$

- J-PARC high momentum beamline?

Heavy lon collision
Coalescence DNN, $\Lambda_{c}{ }^{*} \mathbf{N}$

- RHIC, LHC,...
S. Cho, et al, Phys. Rev. Lett. 106, 212001 (2011); C 84, 064910 (2011)


## Summary <br> Summary <br> Sumimary

## We study DN interaction and DNN system

DN interaction is constructed by
regarding $\Lambda_{c}{ }^{*}$ as "DN quasi-bound state".
DN interaction is constructed by
regarding $\Lambda_{c}{ }^{*}$ as "DN quasi-bound state". A narrow DNN quasi-bound state in spin
$\mathrm{J}=0$ channel. A narrow DNN quasi-bound state in spin
$\mathrm{J}=0$ channel.

DN interaction in I=1 channel (negative parity $\Sigma_{c}{ }^{*}$ ) is important for $\mathrm{J}=1$ result. M. Bayar et al., arXiv:1205.2275 [hep-ph]

$$
\begin{aligned}
& \text { BDNN } \sim 250 \mathrm{MeV}, \quad \mathrm{~B}_{\wedge c^{*} \mathrm{~N}} \sim 40 \mathrm{MeV} \\
& \Gamma \sim 20-40 \mathrm{MeV}
\end{aligned}
$$

DNN system



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*M. Bayar et al., arXiv:1205.2275 [hep-ph] —

