

Determination of the $\pi\Sigma$ scattering lengths from the weak decays of Λ_c



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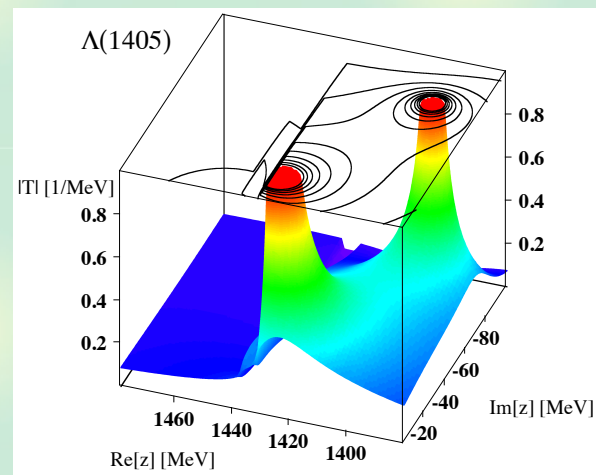
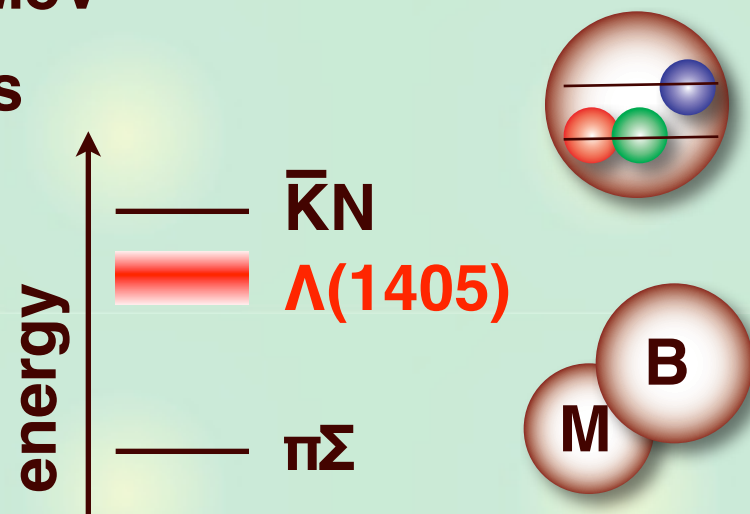
\bar{K} meson and $\bar{K}N$ interaction

Two aspects of \bar{K} meson

- NG boson of chiral $SU(3) \otimes SU(3) \rightarrow SU(3)$
- relatively heavy mass: $M_K \sim 495$ MeV
- > peculiar role in hadron physics

$\bar{K}N$ interaction is ...

- coupled with $\pi\Sigma$ channel
- strongly **attractive**
- > quasi-bound state $\Lambda(1405)$
meson-baryon v.s. qqq state,
double pole, ...
- fundamental building block
for \bar{K} -nuclei, \bar{K} in medium, ...



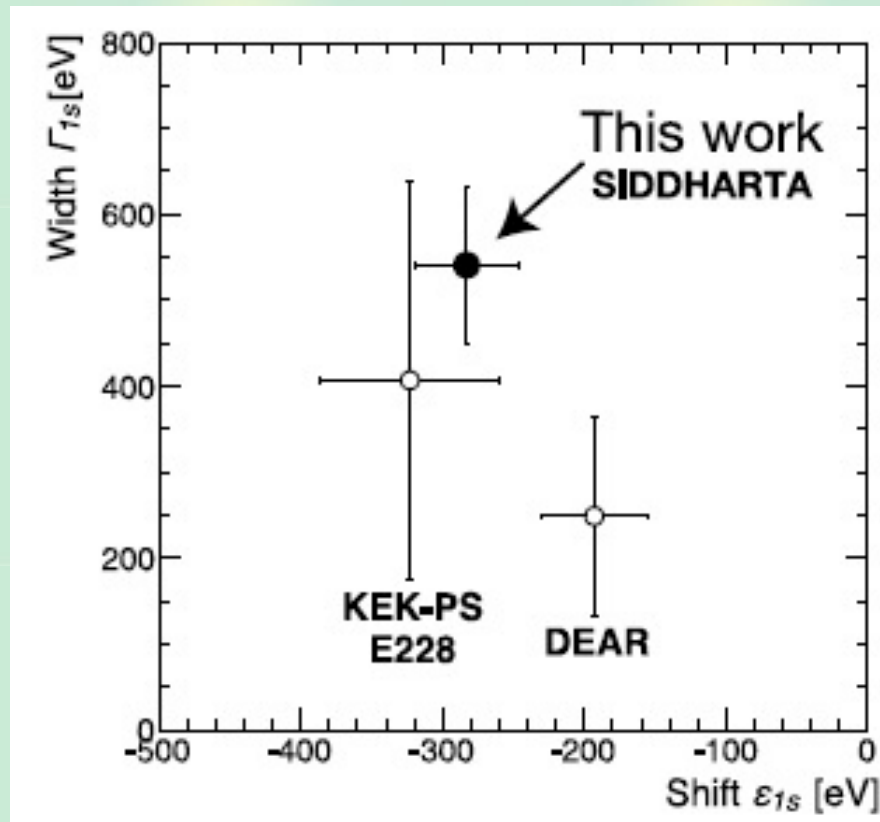
Recent progress

Accurate measurement of Kaonic hydrogen by SIDDHARTA

M. Bazzi, *et al.*, Phys. Lett. B704, 113 (2011)

- smallest uncertainties

$$\Delta E = -283 \pm 36 \pm 6 \text{ eV}, \quad \Gamma = 541 \pm 89 \pm 22 \text{ eV}$$



--> New constraint on the meson-baryon amplitude

Experimental constraints for $S=-1$ MB scattering

K-p total cross sections

- old bubble chamber data, large errors

$\bar{K}N$ threshold observables

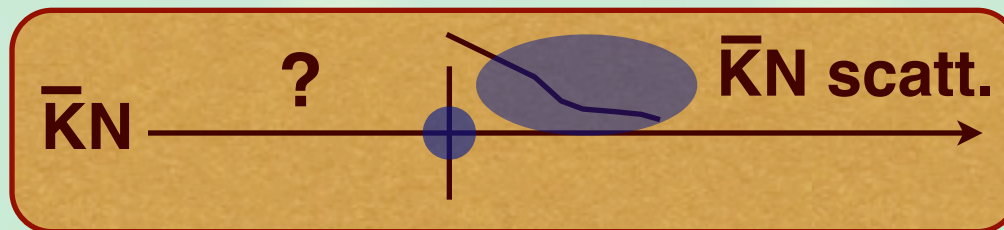
- threshold branching ratios (old but accurate)

$$\gamma = \frac{\Gamma(K^-p \rightarrow \pi^+\Sigma^-)}{\Gamma(K^-p \rightarrow \pi^-\Sigma^+)} = 2.36 \pm 0.04$$

$$R_c = \frac{\Gamma(K^-p \rightarrow \pi^+\Sigma^-, \pi^-\Sigma^+)}{\Gamma(K^-p \rightarrow \text{all inelastic channels})} = 0.664 \pm 0.011$$

$$R_n = \frac{\Gamma(K^-p \rightarrow \pi^0\Lambda)}{\Gamma(K^-p \rightarrow \text{neutral states})} = 0.189 \pm 0.015$$

- K-p scattering length \leftarrow SIDDHARTA



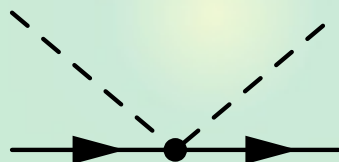
Construction of the realistic amplitude

Systematic analysis by chiral dynamics with SIDDHARTA

Y. Ikeda, T. Hyodo, W. Weise, Phys. Lett. B 706, 63 (2011); Nucl. Phys. A881, 98 (2012)

Interaction kernel: NLO ChPT

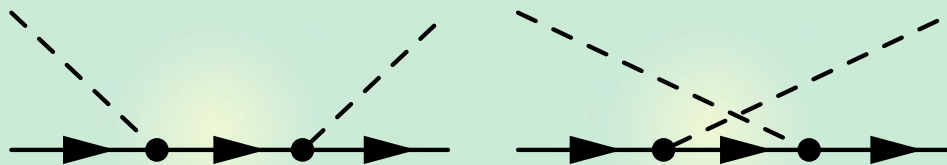
1) TW term



$\mathcal{O}(p)$

TW model

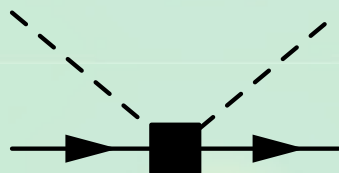
2) Born terms



$\mathcal{O}(p)$

TWB model

3) NLO terms



$\mathcal{O}(p^2)$

NLO model

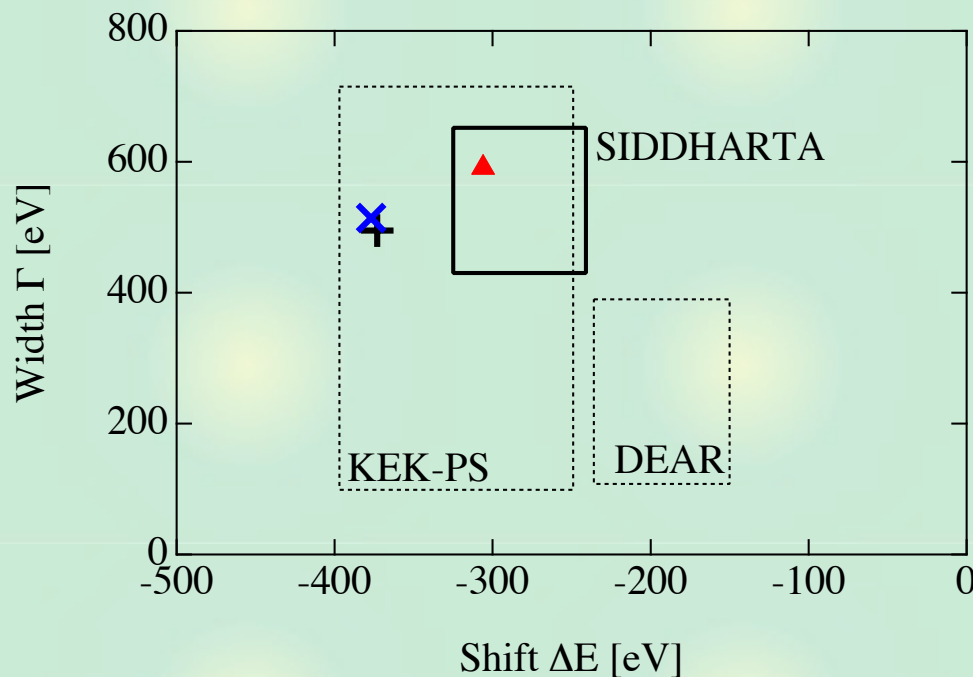
Parameters: 6 cutoffs (+ 7 low energy constants in NLO)

--> fitted to cross section, branching ratio, and SIDDHARTA

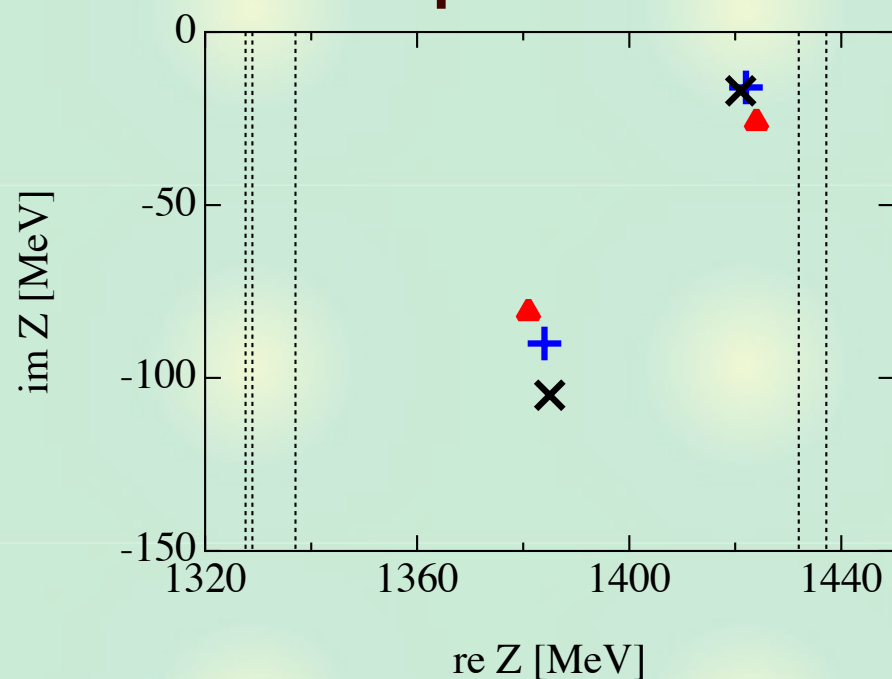
Shift, width, and pole positions

| | TW | TWB | NLO |
|---------------------|------|------|-------|
| χ^2/dof | 1.12 | 1.15 | 0.957 |

Shift and width



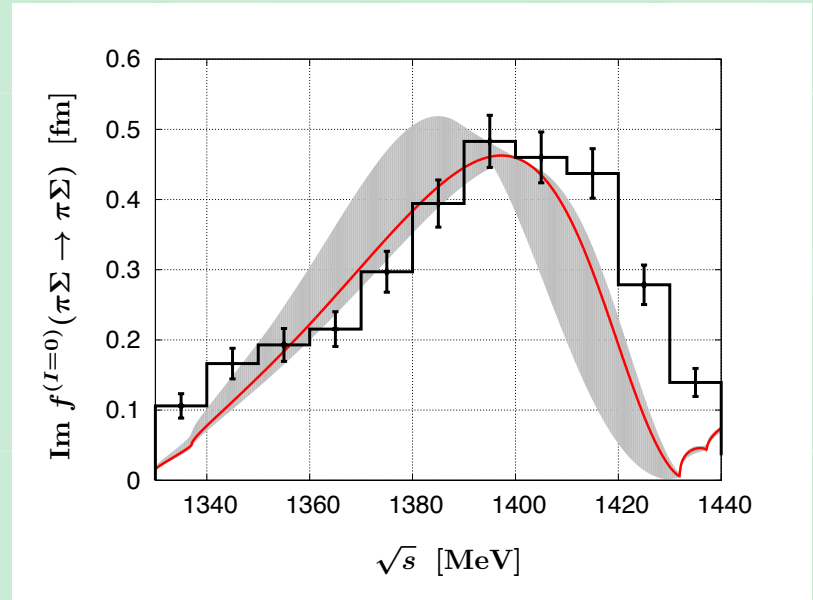
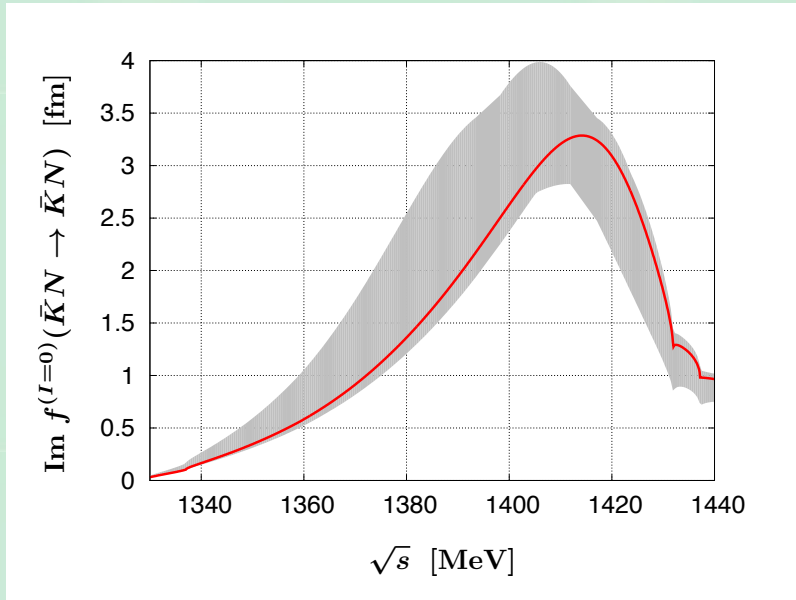
Pole positions



TW and **TWB** are reasonable, while best-fit requires **NLO**. Pole positions are now converging.

Subthreshold extrapolation

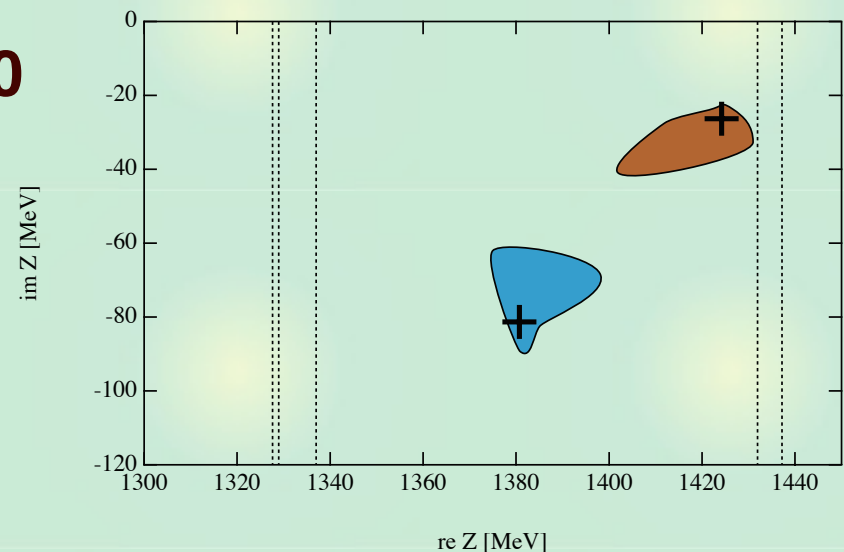
Predicted $\pi\Sigma$ spectrum in comparison with $\bar{K}N$



Note: Hemingway data is not $I=0$

**Shift of the peak position
←-- two poles**

Uncertainty is reduced.

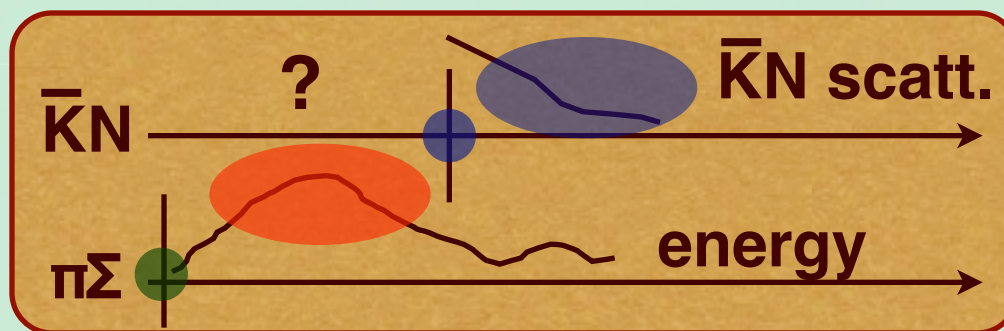


Experimental constraints for $S=-1$ MB scattering

K-p total cross sections

$\bar{K}N$ threshold observables

- threshold branching ratios
- K-p scattering length \leftarrow SIDDHARTA



$\pi\Sigma$ mass spectra

- new data is becoming available (LEPS, CLAS, HADES,...)

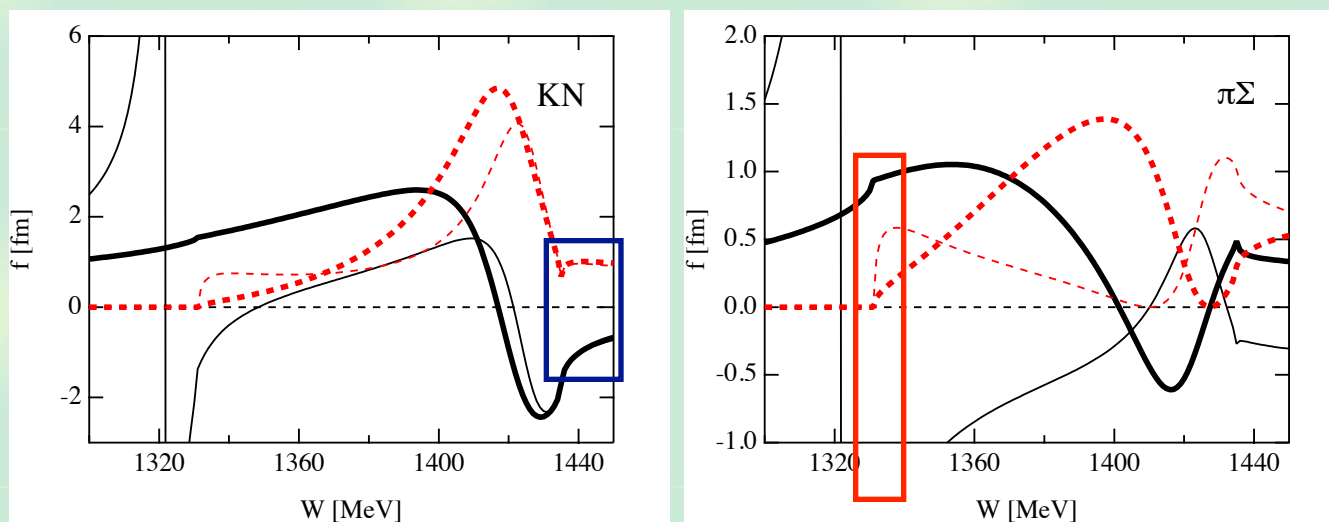
$\pi\Sigma$ threshold observables (so far no data) ?

Importance of $\pi\Sigma$ scattering length

Why threshold behavior of $\pi\Sigma$ channel?

Y. Ikeda, T. Hyodo, D. Jido, H. Kamano, T. Sato, K. Yazaki, PTP 125, 1205 (2011)

Simple model extrapolation with $\bar{K}N(I=0)$ being fixed
 --> large uncertainty at $\pi\Sigma$ threshold



Determination of $\pi\Sigma$ threshold observables

--> understanding of $\Lambda(1405)$, K nuclei, DISTO result,...

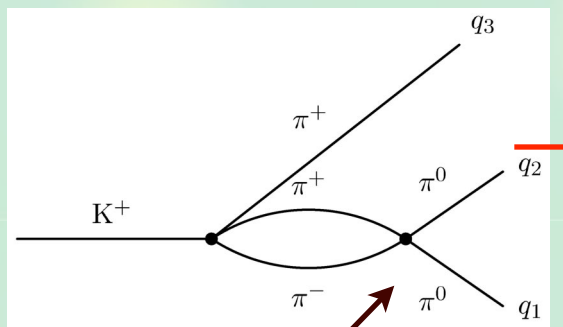
Determination of hadron scattering length

Extraction of hadron scattering length

- shift and width of atomic state (c.f. Kaonic hydrogen)
- extrapolation of low energy phase shift
- final state interaction from heavy particle's decay

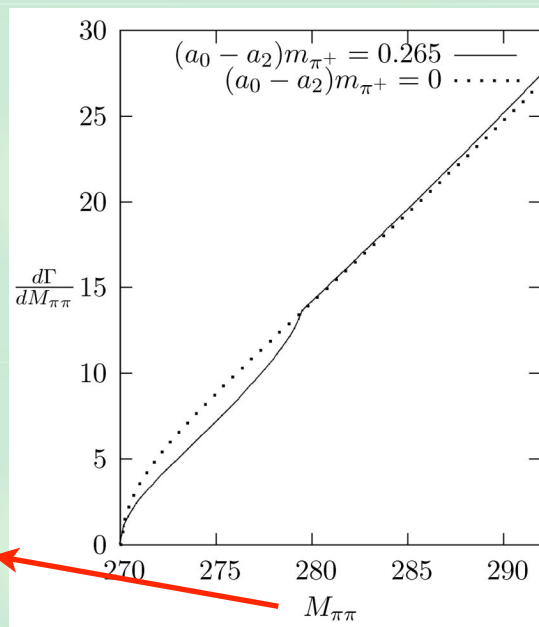
Isospin violation + threshold cusp + amplitude interference

N. Cabibbo, Phys. Rev. Lett. 93, 121801 (2004)

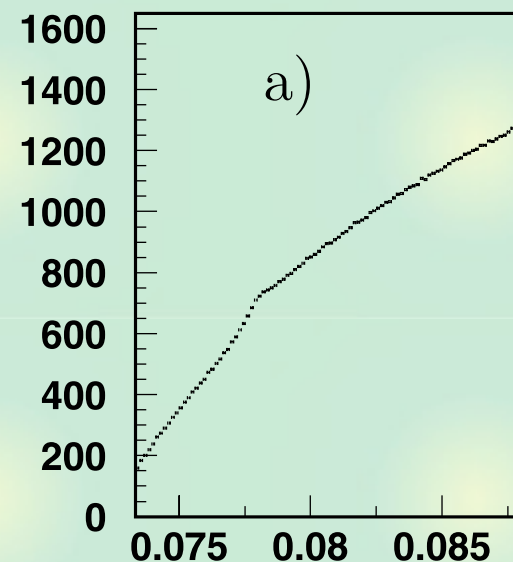


$a_{I=0} - a_{I=2}$

$M_{\pi\pi\pi}$

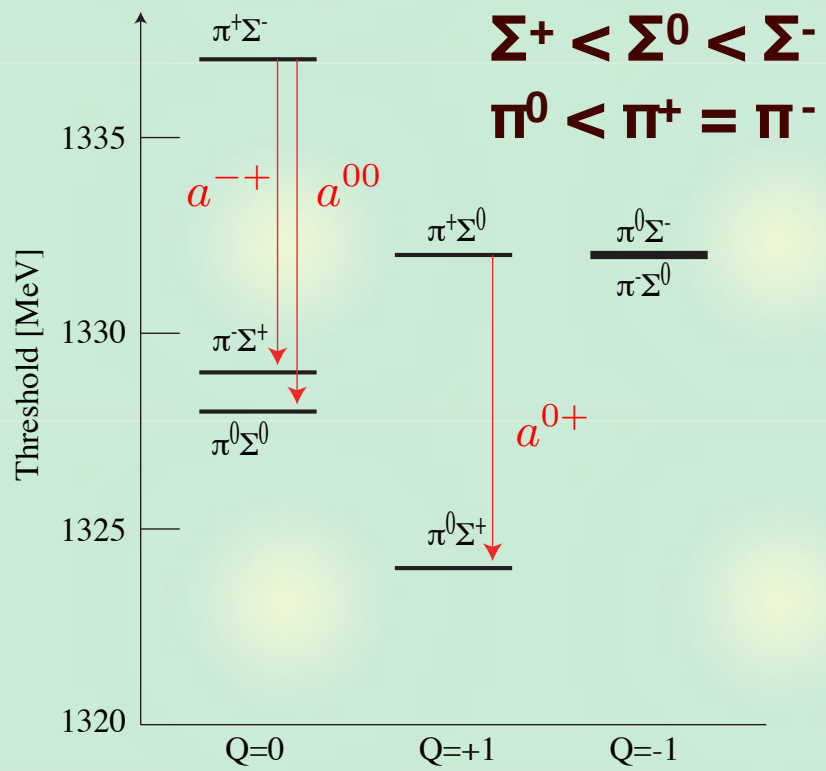
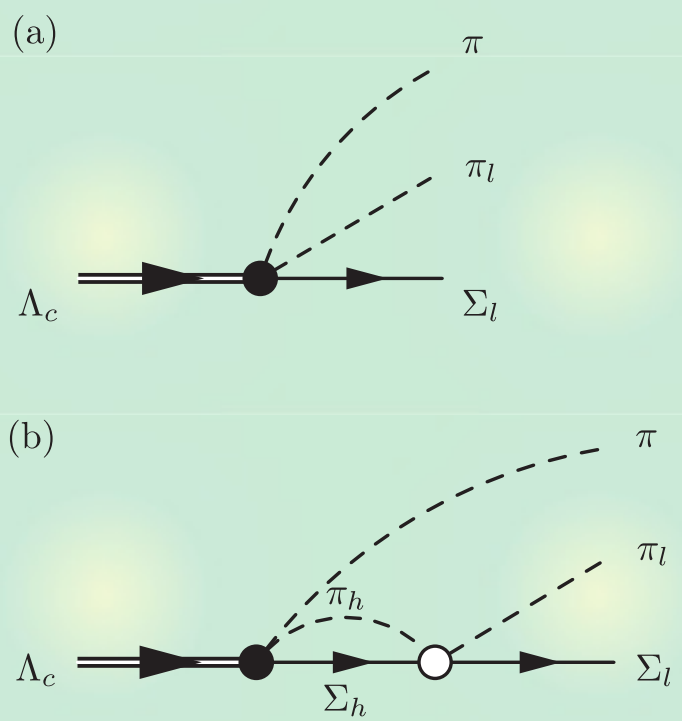


$\times 10^2$



Threshold difference of $\pi\Sigma$ channels

Isospin violation in $\pi\Sigma$ channels



To utilize threshold cusp, appreciable mass difference between $(\pi\Sigma)_h$ and $(\pi\Sigma)_l$ is necessary.

$$a^{--+} : \pi^+\Sigma^- \rightarrow \pi^-\Sigma^+, \quad a^{00} : \pi^+\Sigma^- \rightarrow \pi^0\Sigma^0, \quad a^{0+} : \pi^+\Sigma^0 \rightarrow \pi^0\Sigma^+$$

Determination of $\pi\Sigma$ scattering length

Structure around the cusp in $(\pi\Sigma)_l$ + spectrum in $(\pi\Sigma)_h$
 --> extraction of the scattering length

| mode | $\Lambda_c \rightarrow \pi(\pi\Sigma)_h$ | $\Lambda_c \rightarrow \pi(\pi\Sigma)_l$ |
|----------|--|--|
| a^{-+} | $1.7 \pm 0.5 \%$ | $3.6 \pm 1.0 \%$ |
| a^{00} | $1.7 \pm 0.5 \%$ | $1.8 \pm 0.8 \%$ |
| a^{0+} | $1.8 \pm 0.8 \%$ | not known |

A lot of Λ_c (Belle, Babar, LHC, ...) --> feasible?

Isospin decomposition of three channels

$$a^{-+} = \frac{1}{3}a^0 - \frac{1}{2}a^1 + \frac{1}{6}a^2 + \dots,$$

$$a^{00} = \frac{1}{3}a^0 - \frac{1}{3}a^2 + \dots,$$

$$a^{0+} = -\frac{1}{2}a^1 + \frac{1}{2}a^2 + \dots,$$

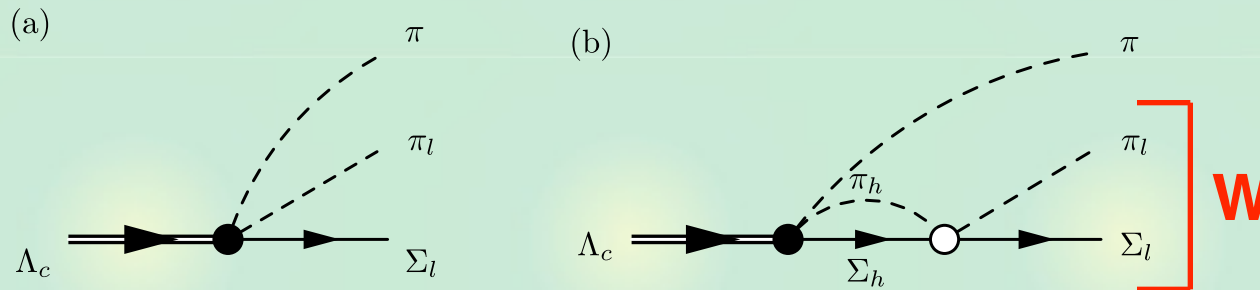
Three unknown scattering lengths, two constraints

$$a^{-+} - a^{00} = a^{0+} + \dots$$

I=2 scattering length: lattice QCD (HAL QCD, NPLQCD,...)

Decay process and intermediate loop

Decay diagrams for $\Lambda_c \rightarrow \pi \pi_l \Sigma_l$ process



Spectral representation of the loop function

$$G(W) = \frac{1}{2\pi} \int_{W_{th}}^{\infty} dW' \frac{\rho(W')}{W - W' + i\epsilon} + (\text{subtractions}) \quad \rho(W) = 2M_h \frac{q(W)}{4\pi W}$$

ρ : phase space, q : three-momentum

Imaginary part of the loop function (on-shell part):

$$\text{Im } G(W) = -\frac{\rho(W)}{2} \Theta(W - W_{th})$$

amplitude (a) : real

amplitude (b) : real ($W < W_{th}$), **complex ($W > W_{th}$)**

Threshold cusp in the spectrum

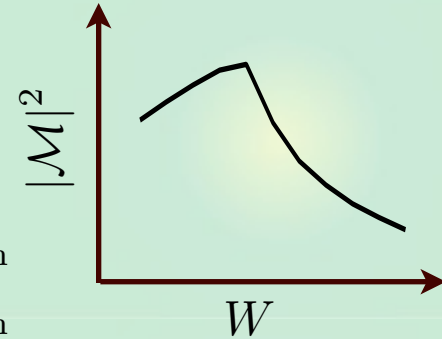
Decomposition of the amplitude

$$\mathcal{M}(W) = \mathcal{M}_0(W) + \tilde{\mathcal{M}}_1(W)m_h\delta$$

$\delta \sim$ real ($W < W_{\text{th}}$), imaginary ($W > W_{\text{th}}$)

$\pi_i \Sigma_i$ invariant mass spectrum (M_0, M_1 : real)

$$|\mathcal{M}|^2 = \begin{cases} (\mathcal{M}_0)^2 + (\tilde{\mathcal{M}}_1 m_h)^2 |\delta|^2 & \text{for } W > W_{\text{th}} \\ (\mathcal{M}_0)^2 + \underline{2\mathcal{M}_0 \tilde{\mathcal{M}}_1 m_h \delta} + (\tilde{\mathcal{M}}_1 m_h)^2 \delta^2 & \text{for } W < W_{\text{th}} \end{cases}$$



- Spectrum is continuous (δ vanishes at threshold)
- **Derivative** of the spectrum is **discontinuous**

$$\left. \frac{d|\mathcal{M}|^2}{dW} \right|_{W \rightarrow W_{\text{th}} - 0} - \left. \frac{d|\mathcal{M}|^2}{dW} \right|_{W \rightarrow W_{\text{th}} + 0} \propto - \frac{2\mathcal{M}_0 \tilde{\mathcal{M}}_1 m_h M_h}{M_h + m_h} \frac{1}{\delta} + \mathcal{O}(\delta)$$

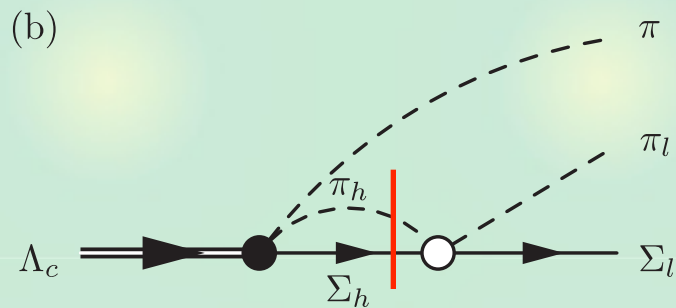
--> threshold cusp

It is purely **kinematical** effect. General phenomena.

Relation to scattering length

The term which produces the cusp

- Energy is fixed at $W = W_{\text{th}}$
- On-shell kinematics for $\pi_h \Sigma_h$ channel



--> amplitude of $\pi_h \Sigma_h \rightarrow \pi_l \Sigma_l$ at threshold: **scattering length**

General decomposition of the amplitude

$$\mathcal{M}(W) = \mathcal{M}_0(W) + \tilde{\mathcal{M}}_1(W) \boxed{e^{i\theta}} n_h \delta$$

- Cusp appears, but **relative phase** affect to the structure.
- Relative phase can be calculated by the dynamical model of final state interactions.

Determination of $\pi\Sigma$ scattering length

Expansion of the decay spectrum (M_0, M_1 : real)

$$|\mathcal{M}|^2 = \begin{cases} A + C'|\delta|^2 + \mathcal{O}(|\delta|^4) & \text{for } W > W_{\text{th}} \\ A + B\delta + C\delta^2 + \mathcal{O}(\delta^3) & \text{for } W < W_{\text{th}} \end{cases}$$

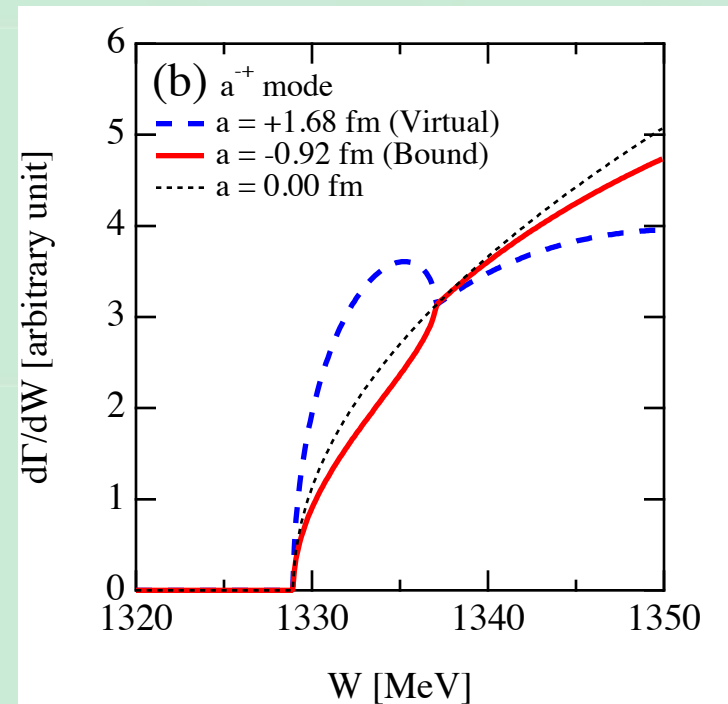
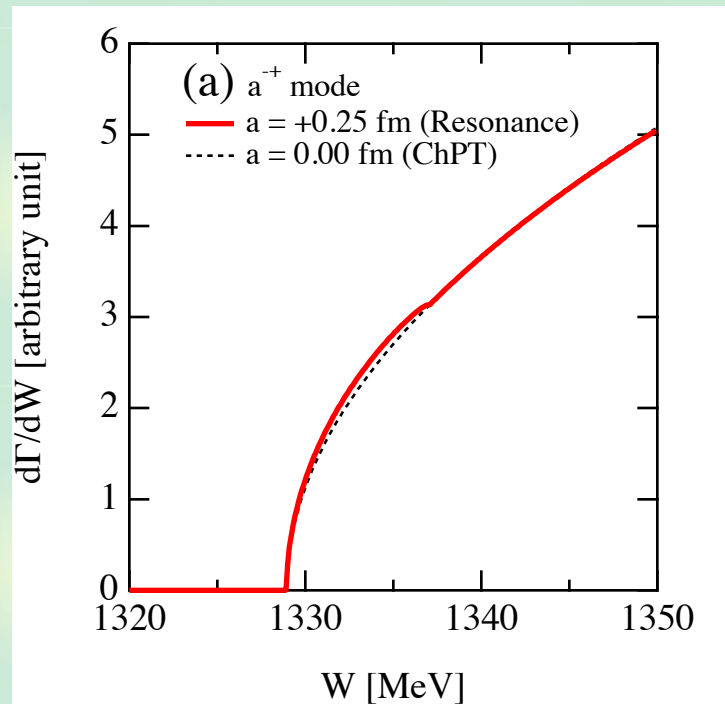
weak process

$$M_0^{h(0)}, M_0^{(0)}$$

$$|a_{h \rightarrow l}| = \frac{|B|}{2m_h \sqrt{A} |\mathcal{M}_0^{h(0)}|}, \quad \frac{a_{h \rightarrow l}}{|a_{h \rightarrow l}|} = \frac{\mathcal{M}_0^{(0)}}{|\mathcal{M}_0^{(0)}|} \cdot \frac{\mathcal{M}_0^{h(0)}}{|\mathcal{M}_0^{h(0)}|} \cdot \frac{B}{|B|}$$

spectrum

$$A, B, C$$

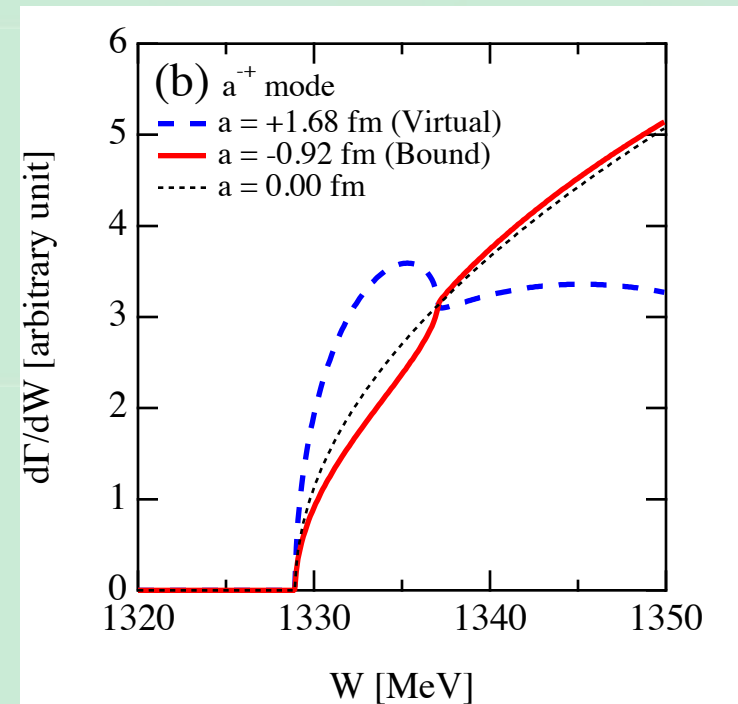
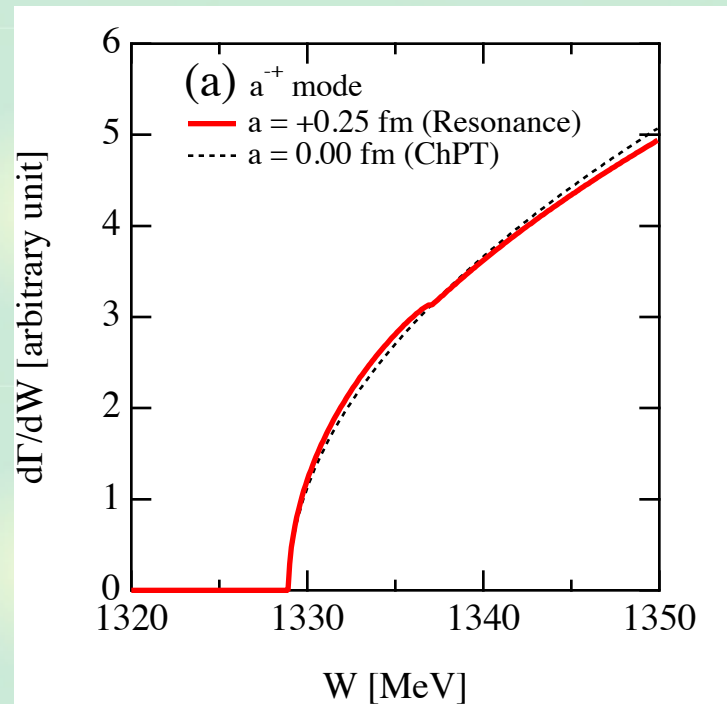


Determination of $\pi\Sigma$ scattering length

Expansion of the decay spectrum (relative phase $\theta = -12$ deg)


$$|\mathcal{M}|^2 = \begin{cases} A + B'|\delta| + C'|\delta|^2 + \mathcal{O}(|\delta|^4) & \text{for } W > W_{\text{th}} \\ A + B\delta + C\delta^2 + \mathcal{O}(\delta^3) & \text{for } W < W_{\text{th}} \end{cases},$$

$$|a_{h \rightarrow l}| = \frac{\sqrt{B^2 + (B')^2}}{2m_h \sqrt{A} |\mathcal{M}_0^{h(0)}|}, \quad \frac{a_{h \rightarrow l}}{|a_{h \rightarrow l}|} = \frac{\mathcal{M}_0^{(0)}}{|\mathcal{M}_0^{(0)}|} \cdot \frac{\mathcal{M}_0^{h(0)}}{|\mathcal{M}_0^{h(0)}|} \cdot \frac{B / \cos \theta}{|B / \cos \theta|}$$




Summary

$\pi\Sigma$ scattering length from Λ_c decay

-  $\pi\Sigma$ scattering length is important for low energy $\bar{K}N$ - $\pi\Sigma$ amplitude.
--> \bar{K} nuclei and $\Lambda(1405)$ physics

Y. Ikeda, T. Hyodo, D. Jido, H. Kamano, T. Sato, K. Yazaki,
Prog. Theor. Phys. 125, 1205 (2011)

-  **Threshold cusp** : kinematical effect

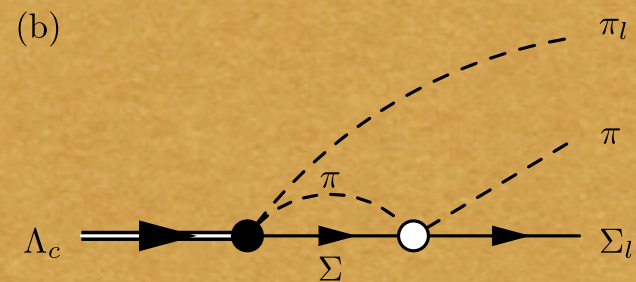
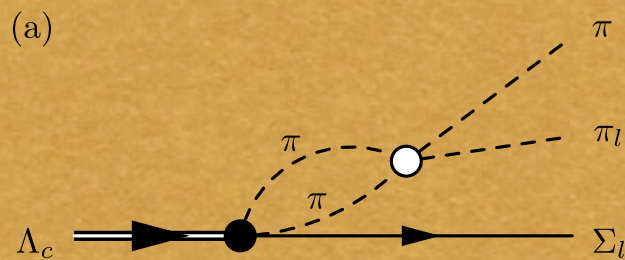
-  Cusp in Λ_c --> $\pi\Sigma$ decay is related with the $\pi\Sigma$ scattering length.

T. Hyodo, M. Oka, Phys. Rev. C 83, 055202 (2011)

Summary

Future plans: estimate of amplitude

$\pi\Sigma$ scattering length final state interaction



- relative phase between M_0 and M_1

Cusp in $\pi\bar{K}N$ channel

$$a(\bar{K}^0 n \rightarrow K^- p) = \frac{1}{2}(a^{I=0} - a^{I=1})$$

c.f. Kaonic hydrogen: $a(K^- p) = \frac{1}{2}(a^{I=0} + a^{I=1})$