## On the compositeness of dynamically generated hadrons



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Introduction

## Compositeness of deuteron

What is the structure of the deuteron?
S. Weinberg, Phys. Rev. 137, B672 (1965)
 $\notin$ NN model space ~ elementary particle

$$
Z \equiv\left|\left\langle B_{0} \mid B\right\rangle\right|^{2} \quad 1=\left|B_{0}\right\rangle\left\langle B_{0}\right|+\int d \boldsymbol{k}|\boldsymbol{k}\rangle\langle\boldsymbol{k}|
$$

model independent relation for weakly bound state

$$
\begin{aligned}
& a_{s}= {\left[\frac{2(1-Z)}{2-Z}\right] \sqrt{R}+\mathcal{O}\left(m_{\pi}^{-1}\right), \quad, r_{e}=\left[\frac{-Z}{1-Z}\right]\left[R+\mathcal{O}\left(m_{\pi}^{-1}\right)<=-\right. \text { Experiments }} \\
& a_{s}=+5.41[\mathrm{fm}], \quad r_{e}=+1.75[\mathrm{fm}], \quad R \equiv(2 \mu B)^{-1 / 2}=4.31[\mathrm{fm}] \\
& \Rightarrow Z \lesssim 0.2 \quad \text {--> deuteron is almost composite }
\end{aligned}
$$

Introduction

## Outline of this talk

Definition of compositeness (Yukawa model)


Example of "physical" state (dynamical model with attractive contact interaction)


## Z in Yukawa model

Field theory with Yukawa coupling ( $\Psi, \Phi, \mathrm{B}_{0}$ )
see D. Lurie and A. J. Macfarlane, Phys. Rev. 136, B816 (1963)

$$
\begin{aligned}
& \mathcal{L}_{0}=\bar{\psi}(i \not \partial-M) \psi+\frac{1}{2}\left(\partial_{\mu} \phi \partial^{\mu} \phi-m^{2} \phi^{2}\right)+\bar{B}_{0}\left(i \not \partial-M_{B_{0}}\right) B_{0} \\
& \mathcal{L}_{\text {int }}=g_{0} \bar{\psi} \phi B_{0}+(\text { h.c. })
\end{aligned}
$$



Physical bound state $B$ at total energy $W=M_{B}$
Free (full) propagator of $\mathrm{B}_{0}(\mathrm{~B})$ field (positive energy part)

$$
\Delta_{0}(W)=\frac{1}{W-M_{B_{0}}}, \quad \Delta(W)=\frac{Z}{W-M_{B}}
$$

Z: field renormalization constant
Dyson equation: relation between full and free propagators

$$
\Delta(W)=\Delta_{0}(W)+\Delta_{0}(W) g_{0} G(W) g_{0} \Delta(W)
$$



## Master formula of compositeness

Solution of Dyson equation and renormalization

$$
\Delta(W)=\frac{1}{W-M_{B_{0}}-g_{0}^{2} G(W)} \rightarrow \frac{1}{W-g_{0}^{2} G(W ; a)}
$$

Renormalization condition, pole at $\mathbf{W}=\mathbf{M}_{\mathbf{B}}: M_{B}=g_{0}^{2} G\left(M_{B} ; a\right)$

The field renormalization constant: residue of the propagator

$$
Z=\lim _{W \rightarrow M_{B}} \frac{W-M_{B}}{W-g_{0}^{2} G(W ; a)}=\frac{1}{1-g_{0}^{2} G^{\prime}\left(M_{B}\right)}
$$

Physical coupling constant: residue of T-matrix

$$
g^{2}=g_{0}^{2} Z
$$

Compositeness in Yukawa theory

$$
1-Z=-g^{2} G^{\prime}\left(M_{B}\right)
$$

## Physical coupling constant

## Single-channel scattering of meson $m$ and baryon $M$

$$
T(W)=\frac{1}{1-V(W) G(W ; a)} V(W)
$$

V: 4-point interaction, attractive

$$
V(W)= \begin{cases}V^{(\text {const })}=C m & \text { constant interaction } \\ V^{(\mathrm{WT})}(W)=C(W-M) & \text { WT interaction }\end{cases}
$$

Bound state condition: pole at $\mathrm{W}=\mathrm{M}_{\mathrm{B}}$

$$
1-V\left(M_{B}\right) G\left(M_{B} ; a\right)=0
$$

Coupling constant: residue of the pole

$$
g^{2}=\lim _{W \rightarrow M_{B}}\left(W-M_{B}\right) T(W)= \begin{cases}-\left[G^{\prime}\left(M_{B}\right)\right]^{-1} & \text { constant interaction } \\ -\left[G^{\prime}\left(M_{B}\right)+\frac{G\left(M_{B} ; a\right)}{M_{B}-M}\right]^{-1} & \text { WT interaction }\end{cases}
$$

Coupling g <-- mass (and cutoff)

Example of "physical" bound state

## Compositeness of bound states

Coupling g --> master formula

$$
1-Z=-g^{2} G^{\prime}\left(M_{B}\right)= \begin{cases}1 & \text { constant interaction } \\ {\left[1+\frac{G\left(M_{B} ; a\right)}{\left(M_{B}-M\right) G^{\prime}\left(M_{B}\right)}\right]^{-1}} & \text { WT interaction }\end{cases}
$$

Constant interaction : purely composite bound state

- equivalence between $\phi^{4}$ and Yukawa for $M_{0}->\infty$

WT interaction : mixture of composite and elementary

- Compositeness is normalized

$$
0 \leq 1-Z \leq 1
$$

- pole term through renormalization
T. Hyodo, D. Jido, A. Hosaka, Phys. Rev. C78, 025203 (2008).

$$
\tilde{V}(W)=C(W-M)-C \frac{(W-M)^{2}}{\left(W-M_{\mathrm{eff}}\right)}
$$

## Check of the natural renormalization scheme

## Natural renormalization condition

## <-- to exclude elementary contribution from the loop function

T. Hyodo, D. Jido, A. Hosaka, Phys. Rev. C78, 025203 (2008)

$$
G\left(W=M ; a_{\text {natural }}\right)=0
$$

1) $\mathbf{a}=a_{\text {natural, }}$ vary $B=M_{B}-M-m$
2) $B=5 \mathrm{MeV}$, vary a

natural scheme --> Z ~ 0
large deviation --> Z ~ 1

We study compositeness of bound states

## constant Z: compositeness

Field renormalization

Master formula
coupling: $g^{2}=g_{0}^{2} Z$
compositeness: $1-Z=-g^{2} G^{\prime}\left(M_{B}\right)$
Example of dynamical bound state constant int. --> purely composite WT int. --> mixture
T. Hyodo, D. Jido, A. Hosaka, Phys. Rev. C 85, 015201 (2012)


# Summary 

Field renormalization
constant Z: compositeness

| Master formula |
| :--- |
| coupling: $g^{2}=g_{0}^{2} Z$ |

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