Softening of the dynamical sigma meson





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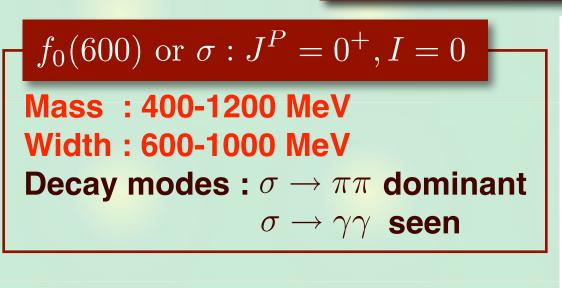
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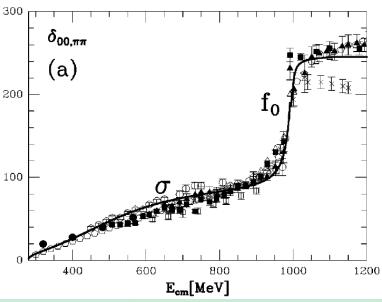
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Introduction Structure of the sigma meson Softening phenomena **Dynamical chiral models** Chiral symmetry and low energy interaction Unitarity and π-π scattering amplitude **Chiral symmetry restoration** Prescription for symmetry restoration Analysis in the restoration limit **Numerical analysis** Softening of the sigma meson Summary

The sigma meson





σ meson

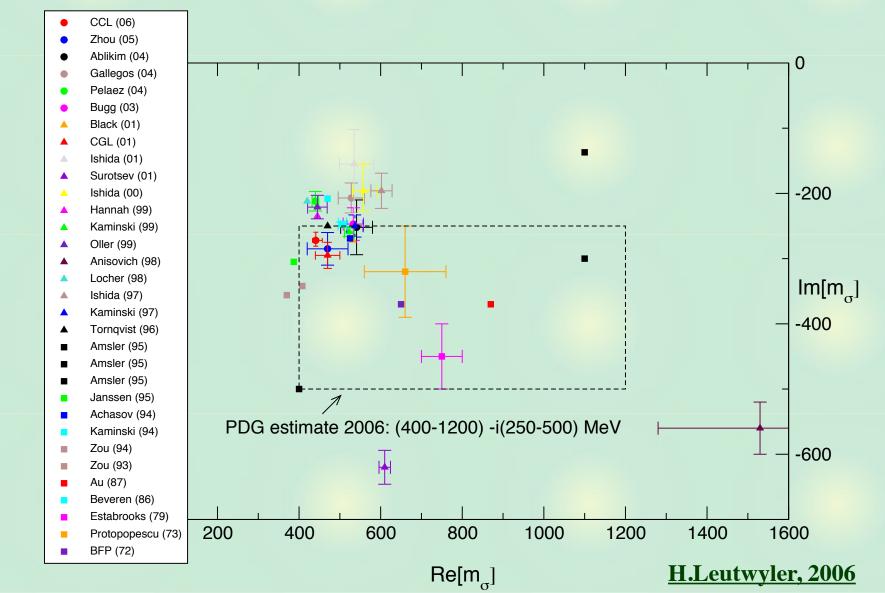
- is the lowest resonance in QCD
- plays an important role in hadron mass generation due to spontaneous chiral symmetry breaking
- provides attraction in phenomenological nuclear force

Recent development

: precise pole position is now available.

Existence of the sigma pole

Development of scattering theory + experimental data



Structure of the sigma meson

- Sigma meson in naive constituent quark model ($\sim \overline{q}q$) has some difficulties: light mass (v.s. p-wave excitation), mass ordering of scalar nonet (v.s. $\sigma > \kappa > f_0 \sim a_0$)
- Alternative descriptions of the sigma meson
 - Chiral sigma (e.g. linear sigma model)

M. Gell-Mann, M. Levy, Nuovo Cim. 16, 705 (1960), ...

- Dynamical sigma (e.g. mesonic molecule generated by π-π attraction) J.A. Oller, E. Oset, J.R. Pelaez, Phys. Rev. D59, 074001 (1999), ...
- CDD pole contribution (e.g. constituent four-quark model, glueball, ...)

We want to clarify the structure <-- softening





Softening of the sigma meson

Softening of chiral sigma

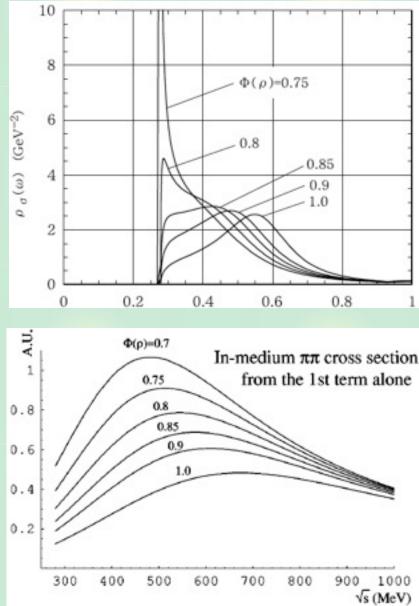
T. Hatsuda, T. Kunihiro, H. Shimizu Phys. Rev. Lett. 82, 2840 (1999)

Spectral enhancement in I=J=0 channel near threshold, when the chiral sym. is partially restored.

sigma: fluctuation of the order parameter of chiral phase transition

Threshold enhancement of π-π cross section, also for the dynamical sigma meson

D. Jido, T. Hatsuda, T. Kunihiro, Phys. Rev. D63, 011901 (2001)



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Softening of the sigma meson

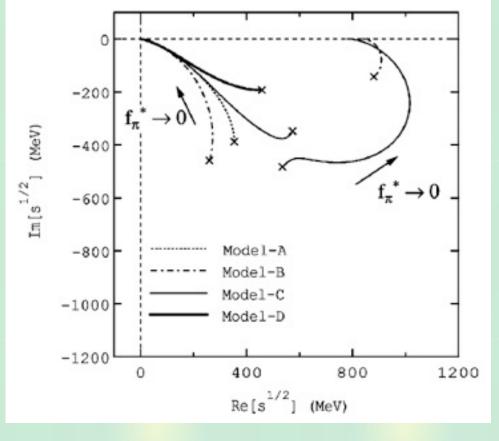
Systematic study up to restoration limit.

K. Yokokawa, T. Hatsuda, A. Hayashigaki, T. Kunihiro, Phys. Rev. C66, 022201 (2002)

 dynamical model with chiral symmetry, unitarity, analyticity, (crossing)

K. Igi, K. Hikasa, Phys. Rev. D59, 034005 (1999)

- 4 cases: σ pole on/off
 ⊗ ρ pole on/off
- roughly corresponds to dynamical sigma and/or CDD pole



- "universal softening" at $f_{\pi}^*/f_{\pi} \ll 1$

Mechanism of the softening

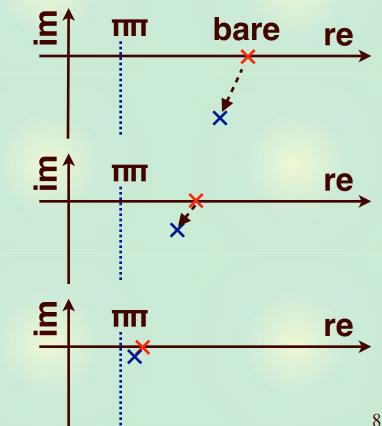
In the previous studies, it seems that the softening takes place, irrespective to the structure of the sigma meson.

Mechanism of the softening?

Softening of the chiral sigma (linear sigma model)

Sigma meson: bare sigma pole acquires finite width through the coupling to π-π

Chiral symmetry restoration: --> lowering bare sigma mass --> reduction of the phase space --> narrow spectrum



Mechanism of the softening

- Softening of the dynamical sigma (ChPT + unitarization)
 - Sigma meson: dynamically generated by π - π attraction
 - **Chiral symmetry restoration:**
 - $--> f_{\pi} \sim <\sigma>$ decreases
 - --> (attractive) interaction ~ $(f_{\pi})^{-2}$ increases
 - --> resonance turns into bound state, spectrum gets narrow
- Special nature of the s-wave resonance:



state (I) resonance Ξ virtual state Ш re pole on the 2nd Riemann resonance (II) sheet below threshold. ex) spin singlet deuteron virtual state (II) bound state --> novel softening pattern?

bound

The aim of this study

We want to study the structure of the "sigma meson" through the behavior in the softening phenomena.

For this, we use a schematic model of chiral dynamics.

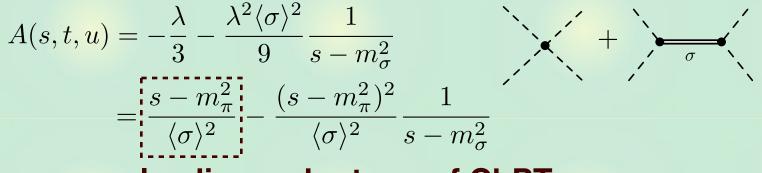
Comparison with previous studies

	mπ	chiral restoration	sigma meson
Jido, Hatsuda, Kunihiro	finite	phi -> 0.7	chiral, dynamical
Yokokawa, et al.	0	phi -> 0	dynamical, CDD, mixture
This work	finite	phi -> 0	chiral, dynamical, CDD, mixture

It is important to keep m_{π} finite and to take restoration limit.

Tree level interaction

- Lagrangian of 2-flavor linear sigma model
 - $\mathcal{L} = \frac{1}{4} \text{Tr} \left[\partial M \partial M^{\dagger} \mu^2 M M^{\dagger} \frac{2\lambda}{4!} (M M^{\dagger})^2 + h(M + M^{\dagger}) \right], M = \sigma + i \boldsymbol{\tau} \cdot \boldsymbol{\pi}$
- 3 parameters <-- m_{π} , m_{σ} , < σ > at mean field level.
- **π-π scattering amplitude in general (crossing symmetry)** $T_{\text{tree}}(s,t,u) = A(s,t,u)\delta_{ab}\delta_{cd} + A(t,s,u)\delta_{ac}\delta_{bd} + A(u,t,s)\delta_{ad}\delta_{bc}$
- Tree-level π-π scattering amplitude



leading order term of ChPT

- low energy expansion
- 1st and 2nd terms are chiral invariant

Tree level interaction

Introduce a parameter using chiral invariant decomposition

$$A(s;x) = \frac{s - m_{\pi}^2}{\langle \sigma \rangle^2} - x \frac{(s - m_{\pi}^2)^2}{\langle \sigma \rangle^2} \frac{1}{s - m_{\sigma}^2}$$

- $x \rightarrow 1$: linear sigma model
- $x \rightarrow 0$: leading order term in ChPT

 $x \rightarrow 1/2$: model C in Yokokawa et al. (σ - ρ degeneracy, KSRF relation, and duality)

Parameter *x* is useful to extrapolate models. The origin of the resonance can be investigated (later).

Projecting the amplitude onto I=J=0, we obtain

$$T_{\text{tree}}(s;x) = \frac{m_{\sigma}^2 - m_{\pi}^2}{\langle \sigma \rangle^2} \left[\frac{2s - m_{\pi}^2}{m_{\sigma}^2 - m_{\pi}^2} (1 - x) - 5x - 3x \frac{m_{\sigma}^2 - m_{\pi}^2}{s - m_{\sigma}^2} - 2x \frac{m_{\sigma}^2 - m_{\pi}^2}{s - 4m_{\pi}^2} \ln\left(\frac{m_{\sigma}^2}{m_{\sigma}^2 + s - 4m_{\pi}^2}\right) \right]$$

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Unitarization

Unitarity of S-matrix : conservation of probability. Tree-level amplitude violates unitarity at certain energy.

Optical theorem :

Im
$$T^{-1}(s) = -\frac{\Theta(s)}{2}$$
 for $s > 4m_{\pi}^2$ $\Theta(s) = (16\pi)^{-1}\sqrt{1 - \frac{4m_{\pi}^2}{s}}$

Scattering amplitude (N/D method + matching with T_{tree})

J.A. Oller, E. Oset, Phys. Rev. D60, 074023 (1999)

$$T(s;x) = \frac{1}{T_{\text{tree}}^{-1}(s;x) + G(s)}$$
$$G(s) = \frac{1}{2} \frac{1}{(4\pi)^2} \left\{ a(\mu) + \ln \frac{m_{\pi}^2}{\mu^2} + \sqrt{1 - \frac{4m_{\pi}^2}{s}} \left[\ln \frac{\sqrt{1 - \frac{4m_{\pi}^2}{s}} + 1}{\sqrt{1 - \frac{4m_{\pi}^2}{s}} - 1} \right] \right\}$$

Left hand cut (crossed diagrams) is neglected.
 zeroth N/D iteration (N=1); c.f. single iteration (N=T_{tree}).

K. Igi, K. Hikasa, Phys. Rev. D59, 034005 (1999)

T

Renormalization

Dynamical model with chiral symmetry.

With sufficient attraction, a resonance can be generated.

Single-subtraction <=> log divergence of loop function We determine the cutoff degree of freedom as

$$G(s) = 0$$
 at $s = m_{\pi}^2$ $a(m_{\pi}) = -\frac{\pi}{\sqrt{3}}$

Exclude the CDD pole contribution from the loop function <u>T. Hyodo, D. Jido, A. Hosaka, Phys. Rev. C78, 025203 (2008)</u>

Consistency of the amplitude with chiral low energy theorem K. Igi, K. Hikasa, Phys. Rev. D59, 034005 (1999)

Crossing symmetry (matching with u-channel amplitude)

M.F.M. Lutz, E. Kolomeitsev, Nucl. Phys. A700, 193 (2002)

Prescription for symmetry restoration

We introduce the effect of chiral symmetry restoration from the outside of the model, by modifying m_{π} , m_{σ} , $<\sigma>$.

1) chiral condensate (pion decay constant) : decreases

 $\langle \sigma \rangle = \Phi \langle \sigma \rangle_0, \quad 0 \le \Phi \le 1$

2) mass of pion : no change

$$\frac{\partial m_{\pi}}{\partial \Phi} = 0$$

3) mass of sigma -- two possibilities

- case I (chiral sigma : decreases)

$$m_{\sigma}|_{\Phi \to 0} = m_{\pi}$$
 (case I) $m_{\sigma} = \sqrt{\lambda \frac{\langle \sigma \rangle^2}{3} + m_{\pi}^2}$

- case II (CDD pole : no change)

 $\frac{\partial m_{\sigma}}{\partial \Phi} = 0 \quad \text{(case II)}$

Restoration limit and chiral partner

Properties of the chiral partner in the restoration limit

- mass degeneracy with pion
- coupling to π - π scattering state vanishes

Model for chiral sigma (x=1, case I ~ linear sigma model)

- pole term in the tree-level interaction

$$T_{\text{tree}}(s;1) = -\frac{\lambda^2 \langle \sigma \rangle^2}{3} \frac{1}{s - m_{\pi}^2 - \frac{\lambda}{3} \langle \sigma \rangle^2} + \dots$$

- renormalization condition

 $G(s) = 0 \quad \text{at} \quad s = m_{\pi}^{2}$ $T(m_{\pi}^{2}; 1)|_{\Phi \to 0} = T_{\text{tree}}(m_{\pi}^{2}; 1)|_{\Phi \to 0} \equiv -\frac{g^{2}}{s - M_{\text{pole}}^{2}}$

Thus, the pole in the π - π amplitude behaves as $g \to 0$, $M_{\text{pole}} \to m_{\pi}$ for $\Phi \to 0$ like the chiral partner

Restoration limit and chiral partner

Model for dynamical sigma and/or CDD pole (case II) $T_{\rm tree}(s;x) \propto rac{1}{\langle \sigma \rangle^2}$

The amplitude in the restoration limit is solely determined by the loop function G, irrespective to x:

$$T(s;x) = \frac{1}{T_{\text{tree}}^{-1}(s,x) + G(s)} \to \frac{1}{G(s)} \quad \text{for} \quad \Phi \to 0$$

- renormalization condition requires a pole at m_{π} G(s) = 0 at $s = m_{\pi}^2$

- coupling can be calculated : proportional to m_{π}

$$g^{2}|_{\Phi \to 0} = -(s - m_{\pi}^{2})T(s)|_{s \to m_{\pi}^{2}, \Phi \to 0}$$
$$= -\frac{s - m_{\pi}^{2}}{G(s)}\Big|_{s \to m_{\pi}^{2}} = (4\pi)^{2} \left(\frac{\pi}{3\sqrt{3}} - \frac{1}{2}\right)^{-1} m_{\pi}^{2}$$

How to interpret this result?

Restoration limit and chiral partner

Properties of chiral sigma

 $g \to 0, \quad M_{\text{pole}} \to m_{\pi} \quad \text{for} \quad \Phi \to 0$

Properties of dynamical sigma

$$g^2 \to (4\pi)^2 \left(\frac{\pi}{3\sqrt{3}} - \frac{1}{2}\right)^{-1} m_\pi^2, \quad M_{\text{pole}} \to m_\pi \quad \text{for} \quad \Phi \to 0$$

- mass degenerates with pion

- coupling to π - π : vanishes in the chiral limit
- **Dynamical sigma as chiral partner of pion?**
 - Renormalization condition plays an important role.
 <-- consistency with chiral theorem. Generally,

 $G(\mu^2) = 0$ at $0 \le \mu \le 2m_{\pi}$ --> deviation ~ m_{π}

- dynamical resonance as chiral partner (mass degeneracy)

J.A. Oller, hep-ph/0007349 S. Leupold, M.F.M. Lutz, M. Wagner, 0811.2398 [nucl-th] Model setup

We numerically analyze four models:

	x	mσ	sigma origin
model A	1	case I	chiral
model B	0	-	dynamical
model C	1	case II	CDD
model D	1/2	case II	CDD + dynamical

model C : same with model A, but m_o unchanged. pole term + repulsion (c.f. linear sigma model)

$$T_{\text{tree}}(s;x) \equiv T_{\text{tree}}^{(\text{contact})}(s;x) + T_{\text{tree}}^{(\text{pole})}(s;x)$$

model D : pole term + attraction

existence of dynamical state <--> sign of the contact term

Results in vacuum

Input: $m_{\pi} = 140 \text{ MeV}, m_{\sigma} = 550 \text{ MeV}, <\sigma > = 93 \text{ MeV}$

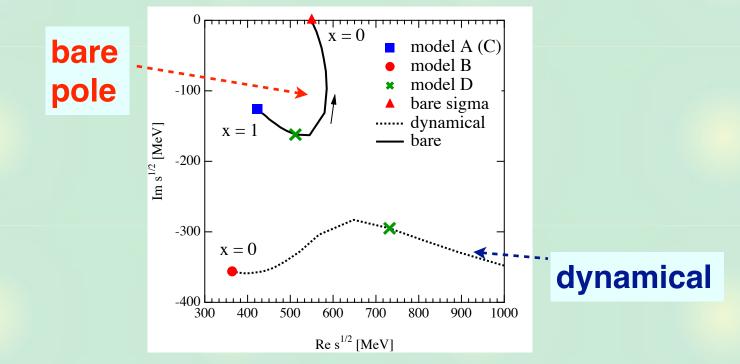
	scattering length (m _π)-1	pole position [MeV]
model A, C	0.244	423 - 126 i
model B	0.174	364 - 356 i
model D	0.208	512 - 162 i, 732 - 295 i
(experiment)	0.216 [1]	441 - 272 i [2]

[1] S. Pislak *et al.*, Phys. Rev. D67, 072004 (2003)

[2] I. Caprini, G. Colangelo, H. Leutwyler, Phys. Rev. Lett. 96, 132001 (2006)

Model extrapolations and origin of the pole

Poles in the complex energy plane



Trace pole position of model A (x=1) with x -> 0 : approaches to bare pole, through one pole in model D

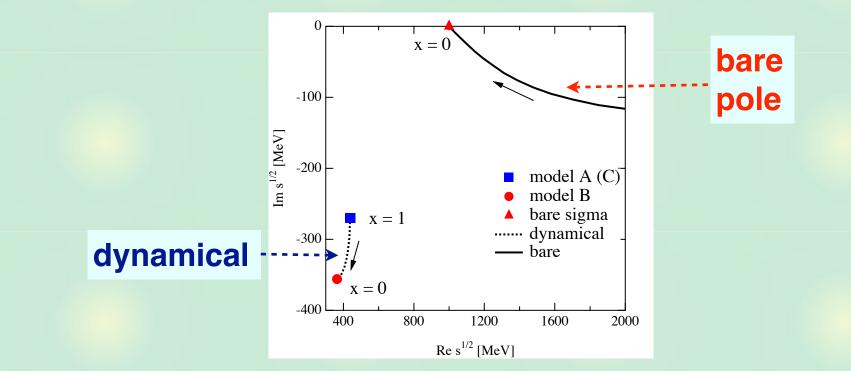
Trace pole position of model B (x=0) with x -> 1

: dissolves into continuum, through one pole in model D

Poles in model D: one bare pole origin, one dynamical.

Model extrapolations and origin of the pole

If $m_{\sigma} = 1$ GeV, then

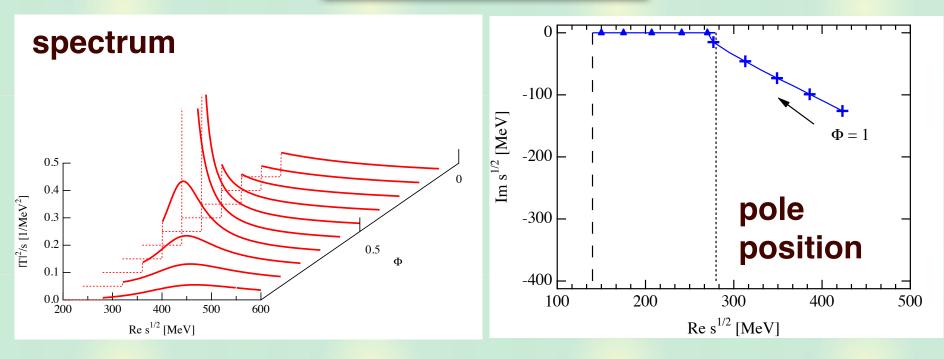


A bare pole at sufficiently high energy than the energy region under consideration ==> effective attraction

$$-G^2 \frac{1}{s - m_\sigma^2} = G^2 \frac{1}{m_\sigma^2} \left(1 + \frac{s}{m_\sigma^2} + \cdots \right) \quad \text{for} \quad s \ll m_\sigma^2.$$

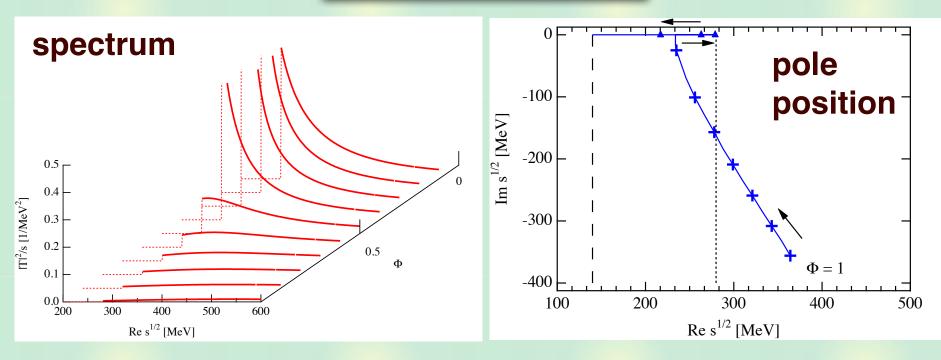
In this case the origin of the pole in model A(C) is dynamical₂₂

Results in model A



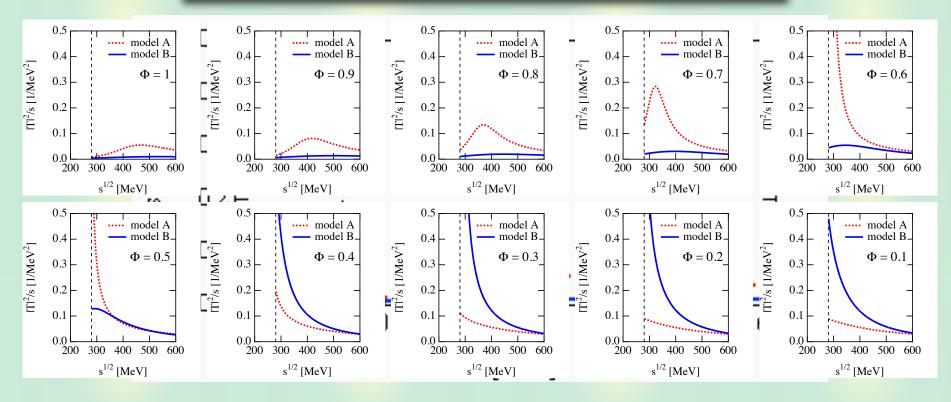
- Linear sigma model + unitarization : chiral sigma
- Softening takes place, as expected.
- $M_{pole} -> m_{\pi} as \Phi -> 0$

Results in model B



- ChPT + unitarization : dynamical sigma
- Softening takes place, but virtual state appears.
- at Re[M_{pole}] = $2m_{\pi}$ ($\Phi \sim 0.6$), due to finite width, spectrum does not show the peak structure
- peak at threshold : $\Phi \sim 0.3 \ll$ formation of bound state
- $M_{pole} -> m_{\pi} as \Phi -> 0$

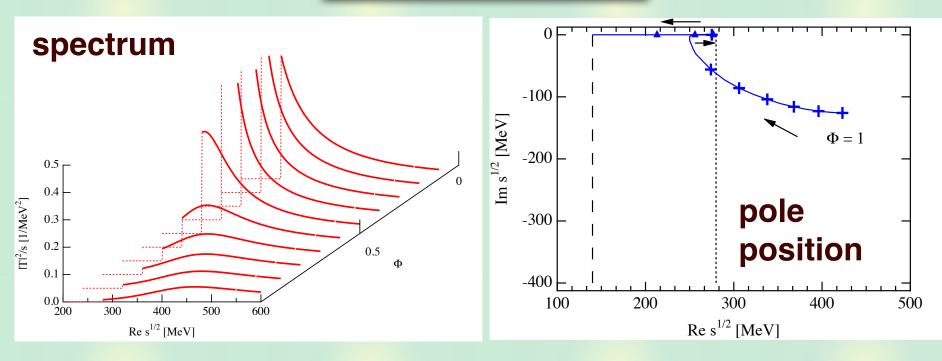
Comparison of model A and model B



- Strong threshold enhancement : different from each other.
- Shape of the spectrum?



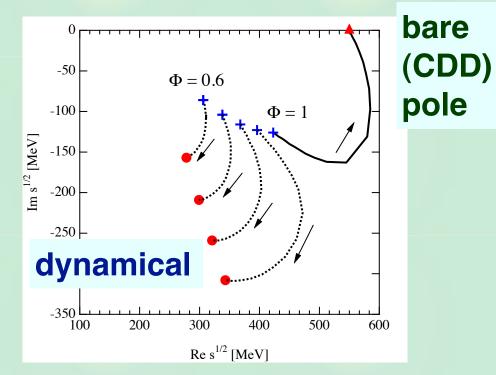
Results in model C



- Bare pole + unitarization : CDD pole
- Softening : similar to dynamical sigma (virtual state).
 <--> bare pole origin ??
- peak at threshold : $\Phi \sim 0.3 \ll$ formation of bound state
- $M_{pole} -> m_{\pi} as \Phi -> 0$

Property change of the pole

Extrapolation to x=0 for 1 \ge $\Phi \ge$ 0.6

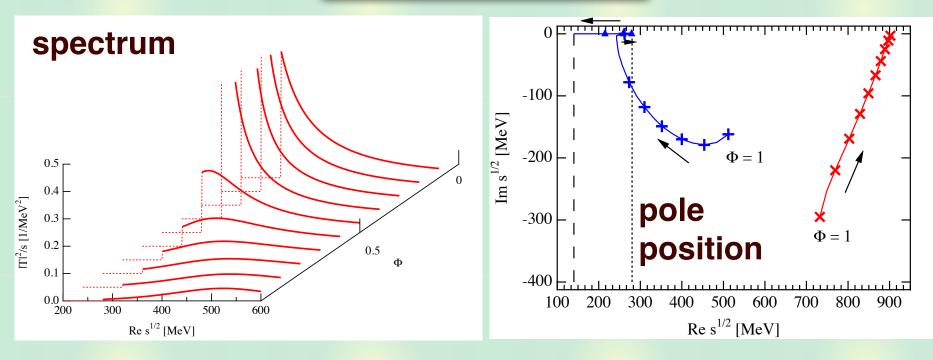


Symmetry restoration --> pole moves to lower energy

Bare pole unchanged, so it behaves as an effective attraction.

Origin of the pole : CDD pole --> dynamical sigma.

Results in model D



- Pole + attraction + unitarization : CDD pole + dynamical
- Softening : similar to dynamical sigma
- the other pole : goes to a kinematical singularity, reducing residue (irrelevant for spectrum).
- M_{pole} -> m_π as Φ -> 0

Summary

Summary : chiral dynamics and sigma meson We study the structure of the sigma meson with chiral symmetry restoration.

- We classify the possible structure of the sigma meson into three classes: (i) chiral sigma (ii) dynamical sigma (iii) CDD pole contribution We construct two-flavor dynamical chiral models which account for them.
 - **Dynamical sigma** (s-wave resonance) is expected to behave differently.

Summary : results

In the chiral restoration limit: Mass and coupling of the dynamical sigma behave similarly with chiral sigma, in the chiral limit. Chiral partner? Softening phenomena: **Dynamical sigma softens qualitatively** differently from chiral sigma. <-- virtual state (s-wave resonance) **CDD** pole case reduces to dynamical **case** around threshold. Universality?