

Chiral dynamics, structure of $\Lambda(1405)$, and $\bar{K}N$ phenomenology



Tetsuo Hyodo^a

Tokyo Institute of Technology^a

2009, Feb. 27th 1

$\Lambda(1405)$ and $\bar{K}N$ dynamics

$\Lambda(1405) : J^P = 1/2^-, I = 0$

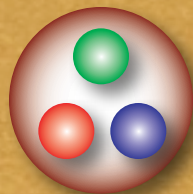
PDG

Mass : 1406.5 ± 4.0 MeV

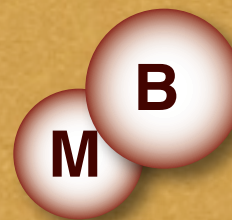
Width : 50 ± 2 MeV

Decay mode : $\Lambda(1405) \rightarrow (\pi\Sigma)_{I=0}$ **100%**

“naive” quark model
: p-wave
~1600 MeV?



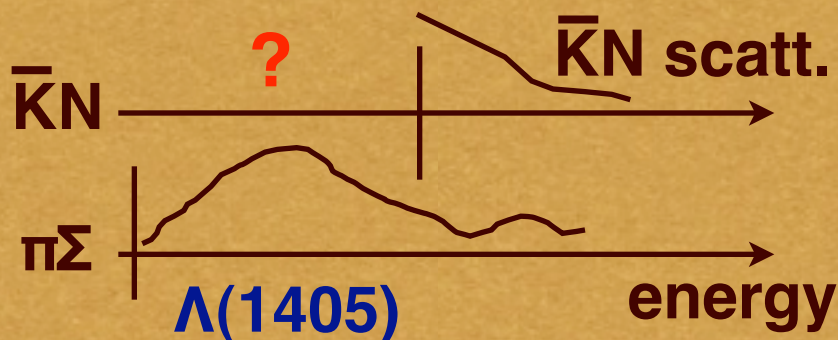
N. Isgur, G. Karl, PRD18, 4187 (1978)



Coupled channel
multi-scattering
←-- strong $\bar{K}N$ int.

R.H. Dalitz, T.C. Wong,
G. Rajasekaran, PR153, 1617 (1967)

$\bar{K}N$ int.
below
threshold



kaonic nuclei,
 $\Lambda(1405)$, ...

-->

exp. @ J-PARC

Chiral dynamics

Description of $S = -1$, $\bar{K}N$ s-wave scattering : $\Lambda(1405)$ in $l=0$

- Interaction \leftarrow chiral symmetry \leftarrow kaon as NG boson

Y. Tomozawa, *Nuovo Cim.* 46A, 707 (1966); S. Weinberg, *Phys. Rev. Lett.* 17, 616 (1966)

- Amplitude \leftarrow unitarity (coupled channel) \leftarrow strong int.

R.H. Dalitz, T.C. Wong, G. Rajasekaran, *PR*153, 1617 (1967)

$$T = \frac{1}{1 - VG} V$$

chiral

N. Kaiser, P. B. Siegel, W. Weise, *Nucl. Phys.* A594, 325 (1995),

E. Oset, A. Ramos, *Nucl. Phys.* A635, 99 (1998),

J. A. Oller, U. G. Meissner, *Phys. Lett.* B500, 263 (2001),

M.F.M. Lutz, E. E. Kolomeitsev, *Nucl. Phys.* A700, 193 (2002), many others

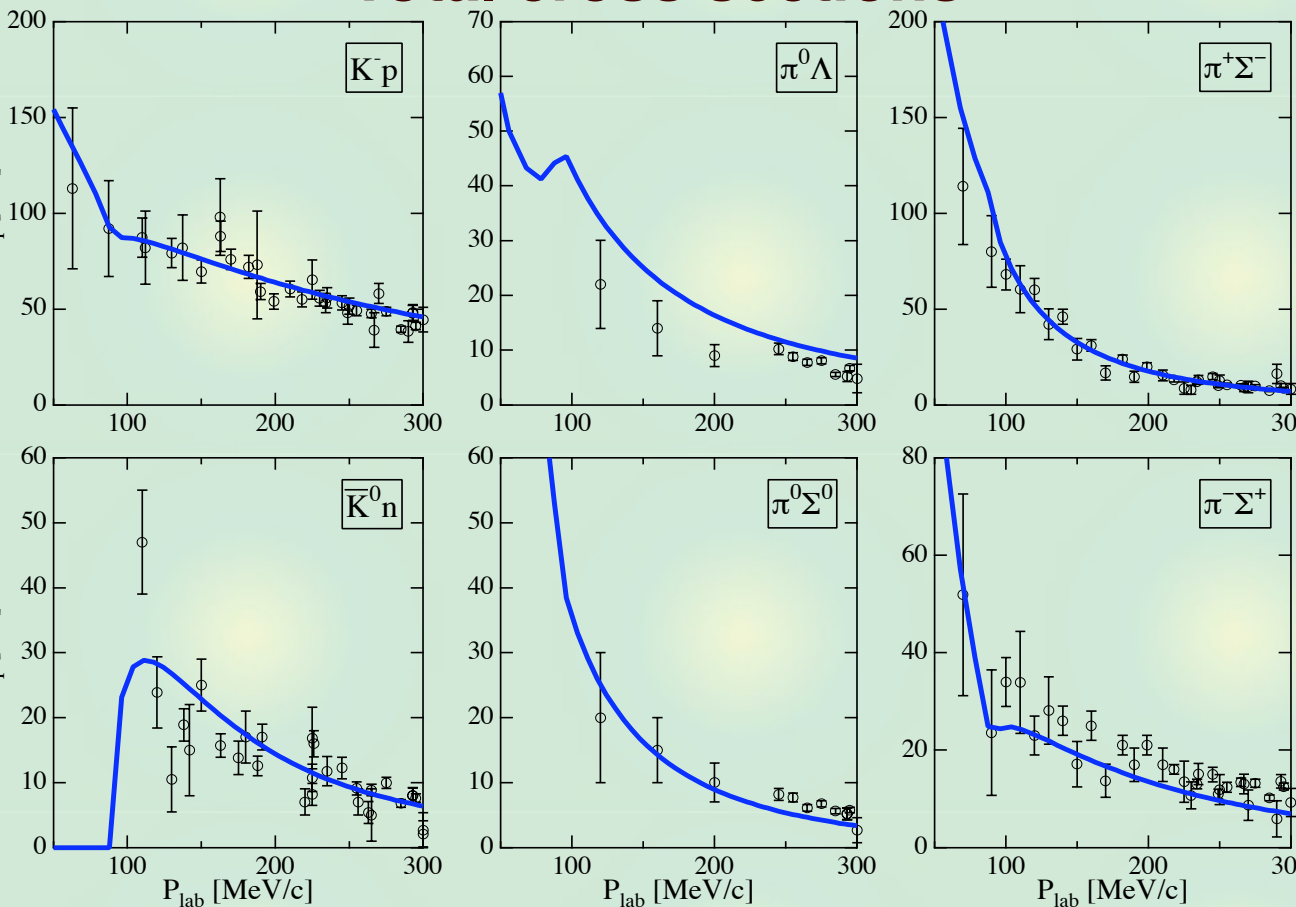
works successfully, also in $S=0$ sector, meson-meson scattering sectors, systems including heavy quarks, ...

T. Hyodo, D. Jido, A. Hosaka, *Phys. Rev. Lett.* 97, 192002 (2006)

T. Hyodo, D. Jido, A. Hosaka, *Phys. Rev. D* 75, 034002 (2007)

How it works? vs experimental data

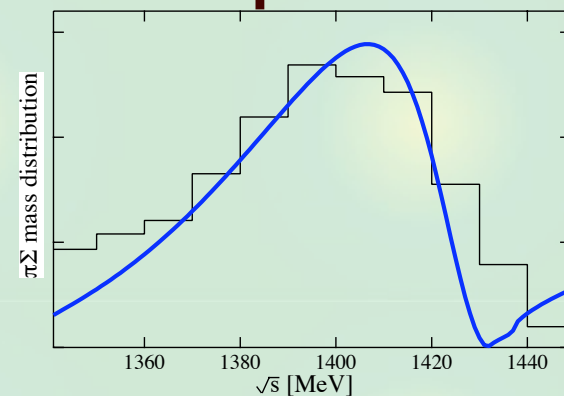
Total cross sections



threshold ratios

	γ	R_c	R_n
exp.	2.36	0.664	0.189
theo.	1.80	0.624	0.225

$\pi\Sigma$ spectrum



T. Hyodo, S.I. Nam, D. Jido, A. Hosaka, Phys. Rev. C68, 018201 (2003),

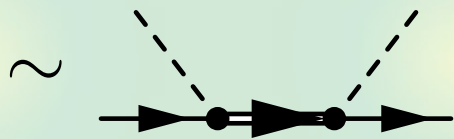
T. Hyodo, S.I. Nam, D. Jido, A. Hosaka, Prog. Theor. Phys. 112, 73 (2004)

Good agreement in wide energy region ($E >, =, <$ threshold). 4

Two poles for one resonance

Poles of the amplitude in the complex plane : resonance

$$T_{ij}(\sqrt{s}) \sim \frac{g_i g_j}{\sqrt{s} - M_R + i\Gamma_R/2}$$

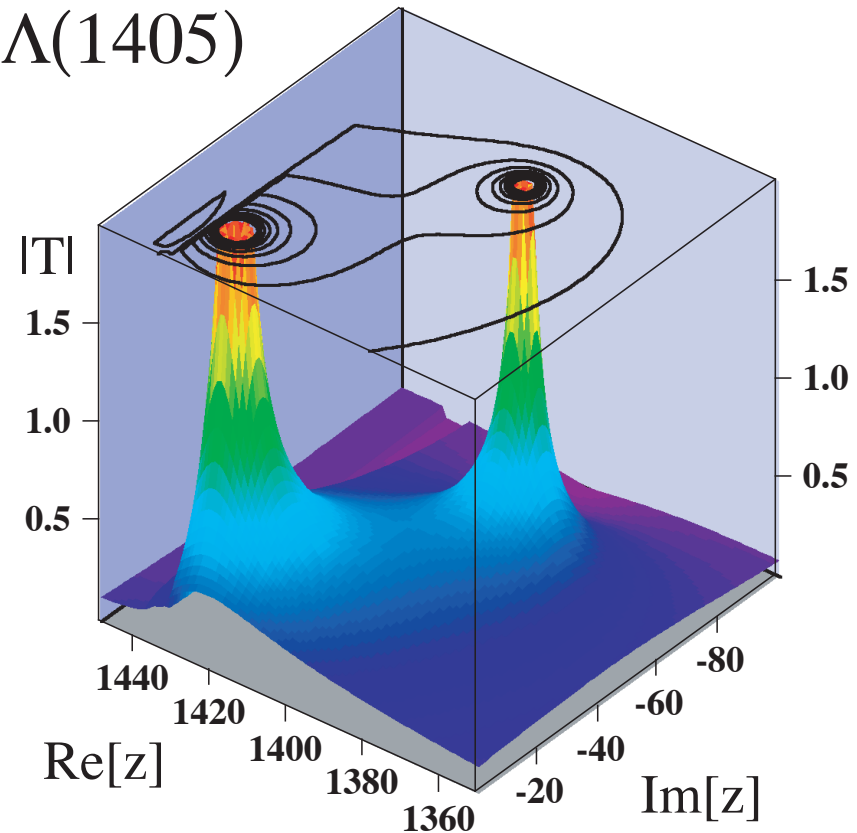


Real part	Mass
Imaginary part	Width/2
Residues	Couplings

Physical state: superposition

$$|\Lambda(1405)\rangle = a|\Lambda_1^*\rangle + b|\Lambda_2^*\rangle$$

$\Lambda(1405)$



D. Jido, J.A. Oller, E. Oset, A. Ramos, U.G. Meissner, Nucl. Phys. A 723, 205 (2003);
 T. Hyodo, W. Weise, Phys. Rev. C 77, 035204 (2008)

Λ(1405) in PDG

PDG

Λ(1405) MASS

PRODUCTION EXPERIMENTS

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	COMMENT
1406.5 ± 4.0				

R.H. Dalitz, A. Deloff, J. Phys G17, 289 (1991)

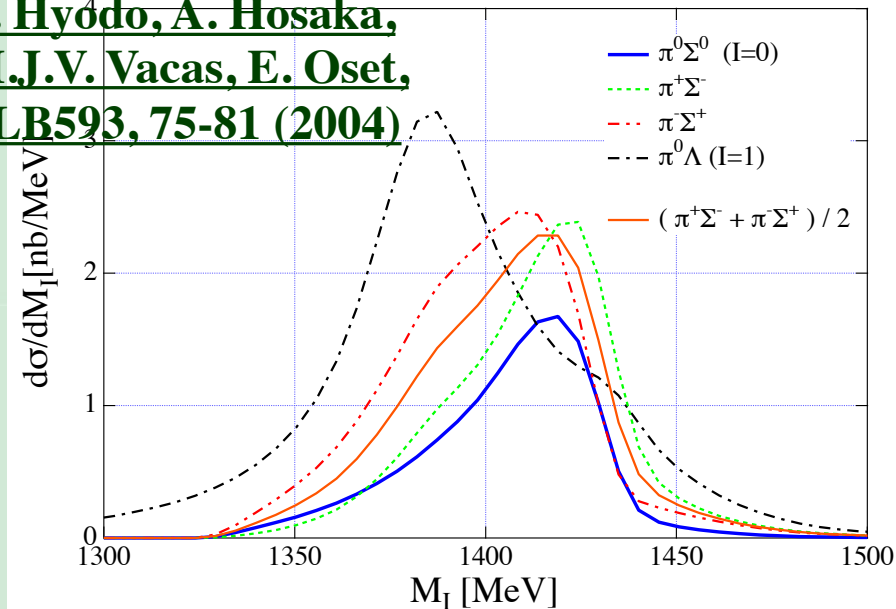
Analysis of **Hemingway data** by **I=0** model. Spectrum ($\pi^- \Sigma^+$) is **not in I=0**.

$$\sigma(\pi^- \Sigma^+) \propto \frac{1}{3} |T^{I=0}|^2$$

$$+ \frac{1}{2} |T^{I=1}|^2 - \frac{2}{\sqrt{6}} \text{Re}(T^{I=0} \cdot T^{I=1})$$

I=1
Σ(1385) **interference**

T. Hyōdo, A. Hosaka,
M.J.V. Vacas, E. Oset,
PLB593, 75-81 (2004)

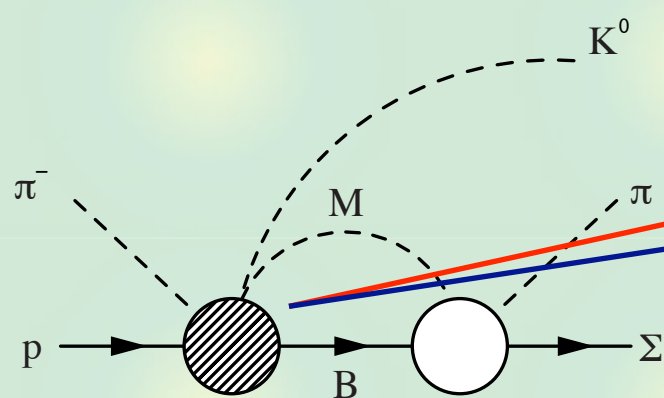


A note on the $\pi\Sigma$ spectrum

アイソスピンの正しい状態 ($I=0$) を選ぶには $\pi\Sigma$ の 3 つの荷電状態 ($\pi^0\Sigma^0, \pi^\pm\Sigma^\mp$) を **全て同時に** 測定する必要がある (未達成)。

(現実的には $\pi^0\Sigma^0$ は $I=1$ がないので理想的?)

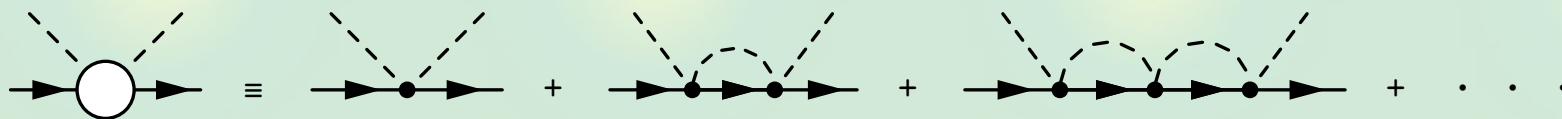
ポールが 2 つある効果は、スペクトルが **反応によって** 変化することを調べる必要がある。=> **1 つの実験で検証 / 排除は不可能**




$$|\Lambda(1405)\rangle = a|\Lambda_1^*\rangle + b|\Lambda_2^*\rangle$$

ただし反応計算は模型依存。

1 ポールでも干渉でピーク位置が変化。



Contents



Structure of $\Lambda(1405)$ resonance (Bグループ)

- Dynamical or CDD (genuine quark state) ?

T. Hyodo, D. Jido, A. Hosaka, Phys. Rev. C78, 025203 (2008).

- Nc Behavior and quark structure

T. Hyodo, D. Jido, L. Roca, Phys. Rev. D77, 056010 (2008).

L. Roca, T. Hyodo, D. Jido, Nucl. Phys. A809, 65 (2008).

- Electromagnetic properties

T. Sekihara, T. Hyodo, D. Jido, Phys. Lett. B669, 133-138 (2008).



Phenomenology of $\bar{K}N$ interaction

- Construction of local $\bar{K}N$ potential (Cグループ)

T. Hyodo, W. Weise, Phys. Rev. C77, 035204 (2008).

- Application to three-body $\bar{K}NN$ system

A. Doté, T. Hyodo, W. Weise, Nucl. Phys. A804, 197 (2008)

A. Doté, T. Hyodo, W. Weise, Phys. Rev. C 79, 014003 (2009)

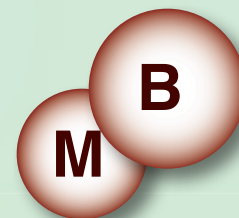
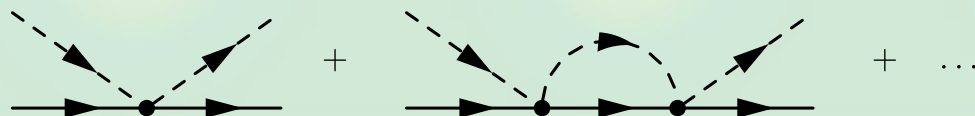
$+\alpha$

Dynamical state and CDD pole

Resonances in two-body scattering

- Knowledge of interaction (potential)
- Experimental data (cross section, phase shift,...)

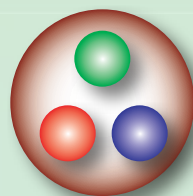
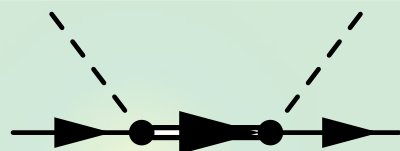
(a) **dynamical** state: molecule, quasi-bound, ...



e.g.) Deuteron in NN , positronium in e^+e^- , (σ in $\pi\pi$), ...

(b) **CDD** pole: elementary, independent, ...

L. Castillejo, R.H. Dalitz, F.J. Dyson, *Phys. Rev.* 101, 453 (1956)



e.g.) J/ψ in e^+e^- , (ρ in $\pi\pi$), ...

Resonances in chiral unitary approach \rightarrow (a) dynamical?

CDD pole contribution in chiral unitary approach

Amplitude in chiral unitary model

$$T = \frac{1}{\boxed{V^{-1}} - \boxed{G}}$$

V : interaction kernel (potential)
G : loop integral (Green's function)

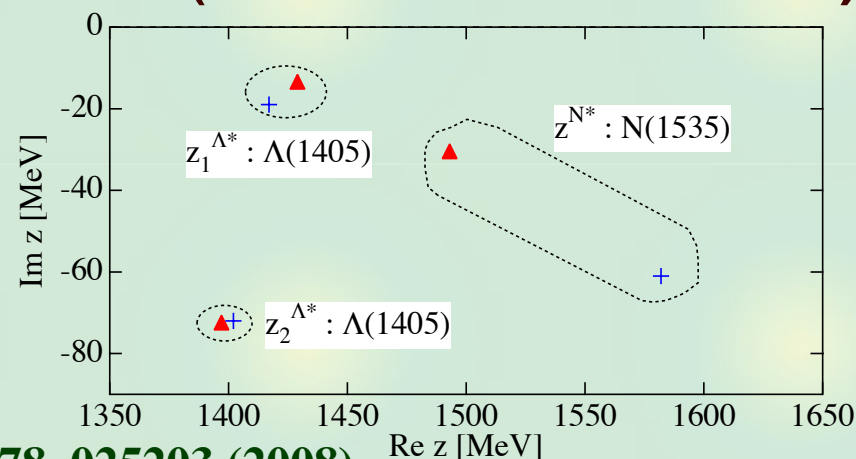
Known CDD pole contribution

- (1) Explicit resonance field in **V**
- (2) Contracted resonance propagator in **V**

Defining “natural renormalization scheme”,
we find **CDD pole contribution in G** (subtraction constant).

N(1535) in πN scattering
--> dynamical + CDD pole

$\Lambda(1405)$ in $\bar{K} N$ scattering
--> **mostly dynamical**



Nc scaling in the model

Nc : number of color in QCD

Hadron effective theory / quark structure

The Nc behavior is known from the general argument.

← introducing Nc dependence in the model,
analyze the resonance properties with respect to Nc

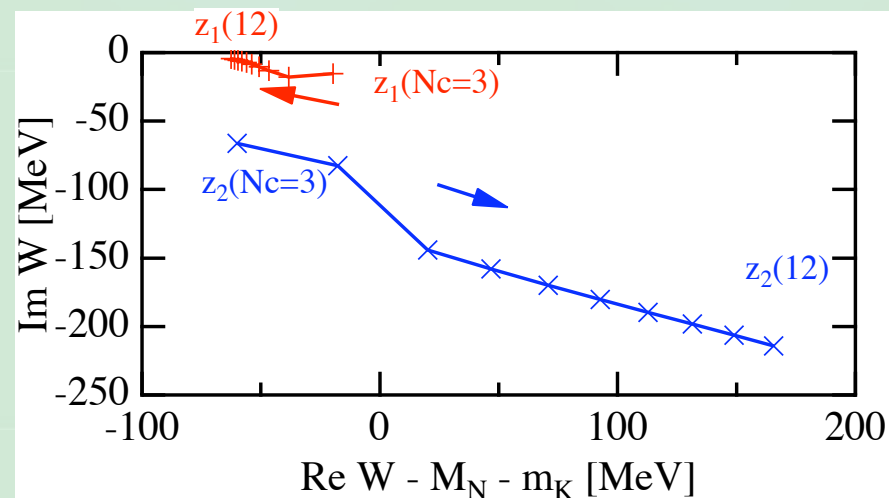
J.R. Pelaez, *Phys. Rev. Lett.* **92**, 102001 (2004)

**Nc scaling of (excited)
qqq baryon**

$$M_R \sim \mathcal{O}(N_c), \quad \Gamma_R \sim \mathcal{O}(1)$$

Result : $\Gamma_R \neq \mathcal{O}(1)$

~ non-qqq (i.e. dynamical) structure



詳細は学会で

29pSJ-2

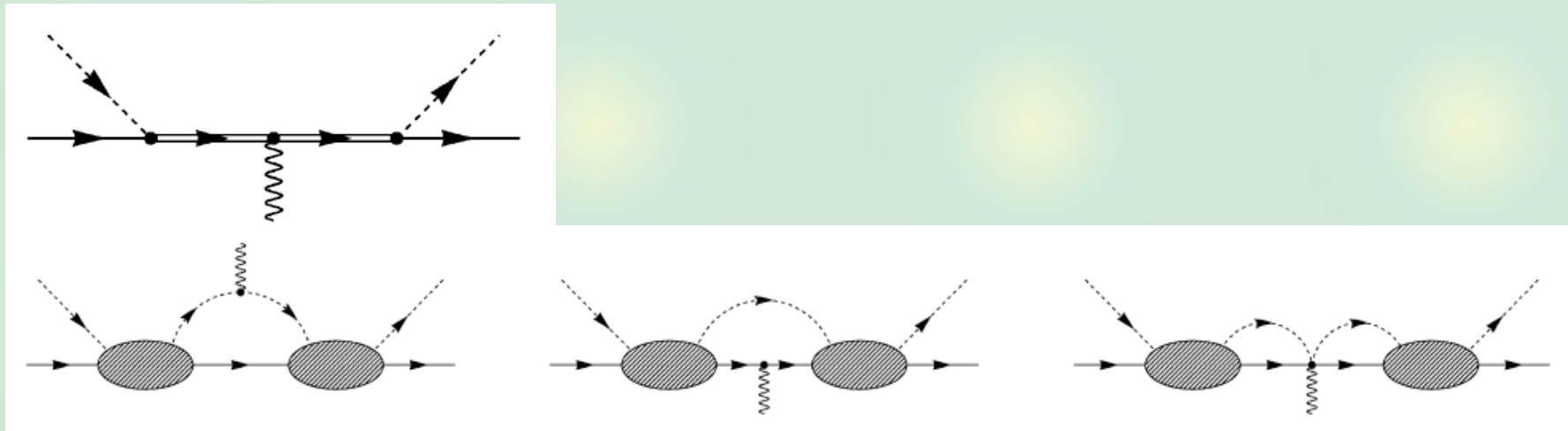
T. Hyodo, D. Jido, L. Roca, *Phys. Rev. D* **77**, 056010 (2008).

L. Roca, T. Hyodo, D. Jido, *Nucl. Phys. A* **809**, 65 (2008).

Electromagnetic properties

Attaching photon to resonance

--> em properties : rms, form factors,...



result of mean squared radii :

$$\langle r^2 \rangle_E = -2.193 \text{ fm}^2$$

large (em) size of the $\Lambda(1405)$

--> meson-baryon picture

Summary 1 : Structure of $\Lambda(1405)$

We study the structure of the $\Lambda(1405)$



Dynamical or CDD?

=> dominance of the MB components



Analysis of N_c scaling

=> non-qqq structure







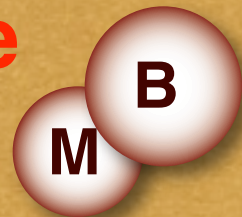
Electromagnetic properties

=> large e.m. size

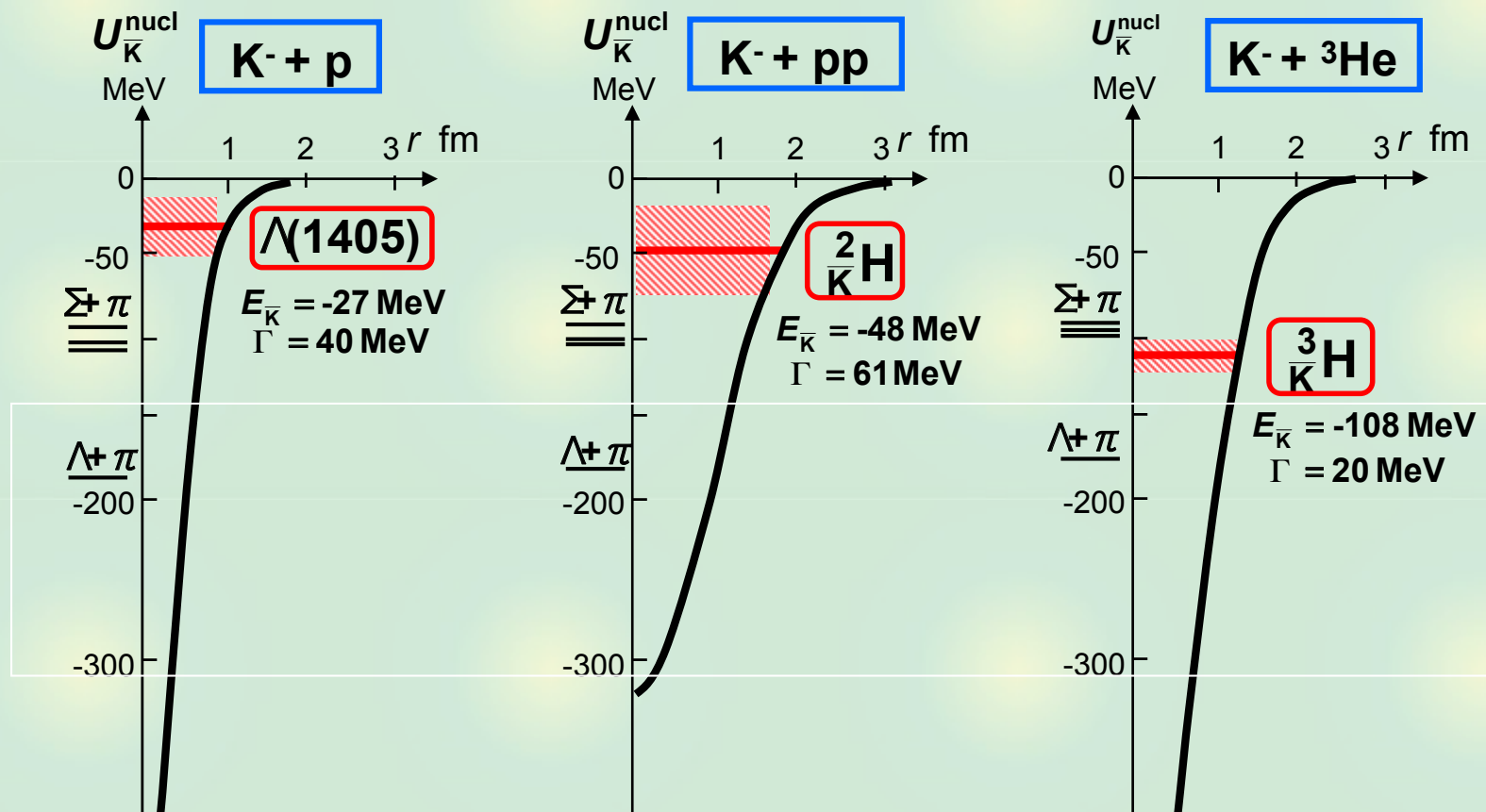
Summary 1 : Structure of $\Lambda(1405)$

We study the structure of the $\Lambda(1405)$

-  Dynamical or CDD?
=> dominance of the MB components
-  Analysis of N_c scaling
=> non-qqq structure
-  Electromagnetic properties
=> large e.m. size
-  Independent analyses consistently support the **meson-baryon molecule picture** of the $\Lambda(1405)$



Deeply bound (few-body) kaonic nuclei?



Potential is purely phenomenological.
What does chiral dynamics tell us about it?

Effective interaction based on chiral SU(3) dynamics

Few-body kaonic nuclei in chiral dynamics

- single-channel $\bar{K}N$ potential

Construction of effective single-channel potential

T. Hyodo, W. Weise, Phys. Rev. C 77, 035204 (2008)

1) Coupled-channel \rightarrow single $\bar{K}N$ channel BS equation
incorporation of $\pi\Sigma$ channel (exact)

2) Local potential in Schrödinger equation (approximate)

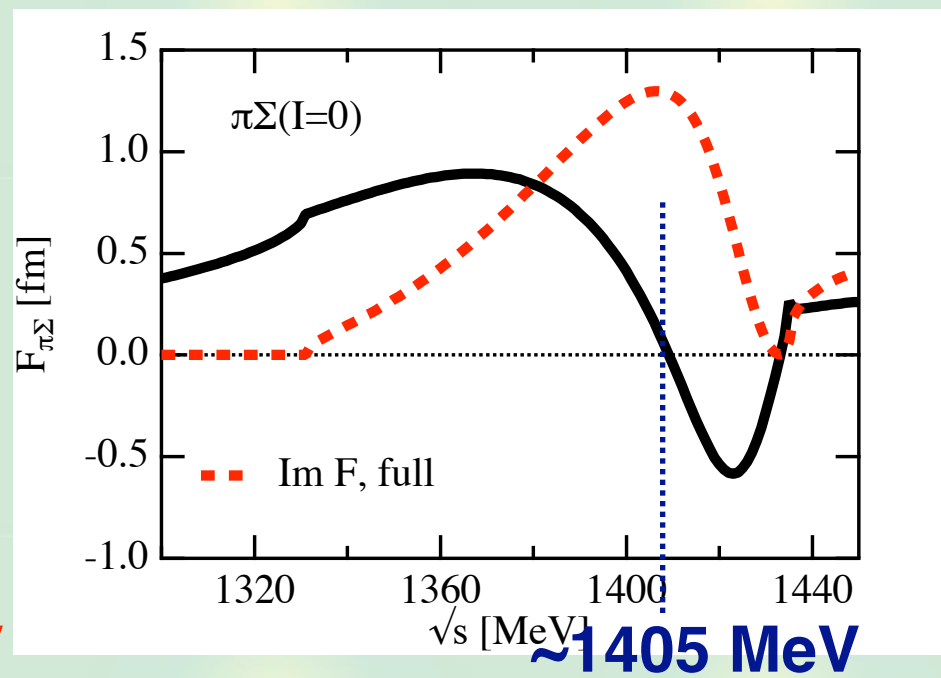
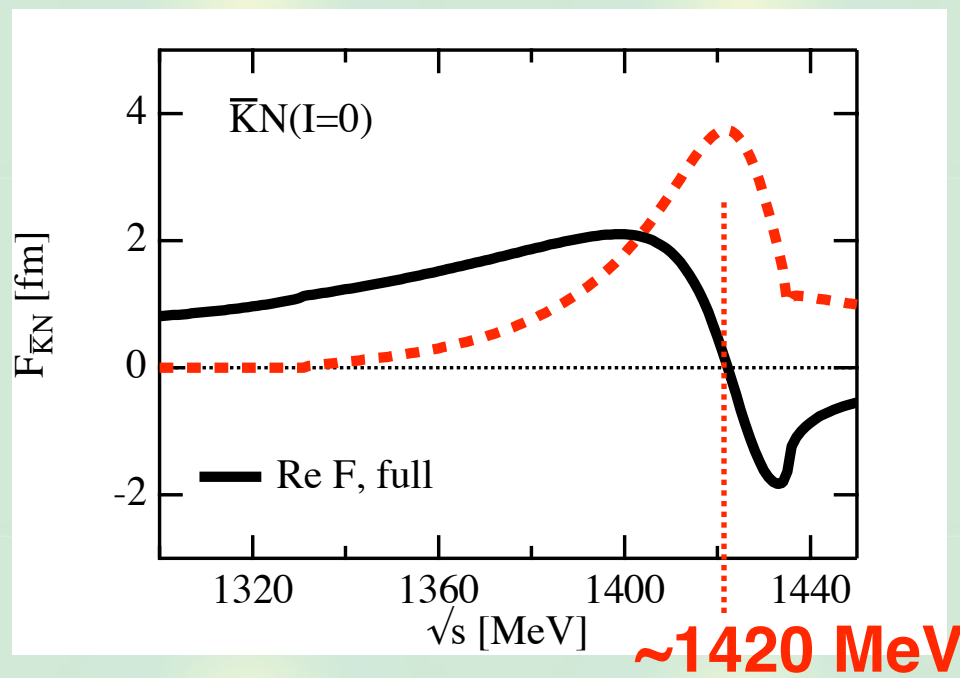
\rightarrow $\bar{K}N$ interaction : attractive, but **weaker than the phenomenological potential.**

Application to K-pp system : bound, but $B \sim 20$ MeV

A. Doté, T. Hyodo, W. Weise, Nucl. Phys. A 804, 197 (2008); Phys. Rev. C 79, 014003 (2009)

Why the interaction is weaker? \rightarrow structure of the $\Lambda(1405)$

Scattering amplitude in $\bar{K}N$ and $\pi\Sigma$



Resonance in $\bar{K}N$: around 1420 MeV
 <-- strong $\pi\Sigma$ dynamics (coupled-channel)

Binding energy : $B = 15 \text{ MeV} \leftrightarrow 30 \text{ MeV}$

Two poles with same quantum numbers
 Different weights of the pole residues --> different spectra

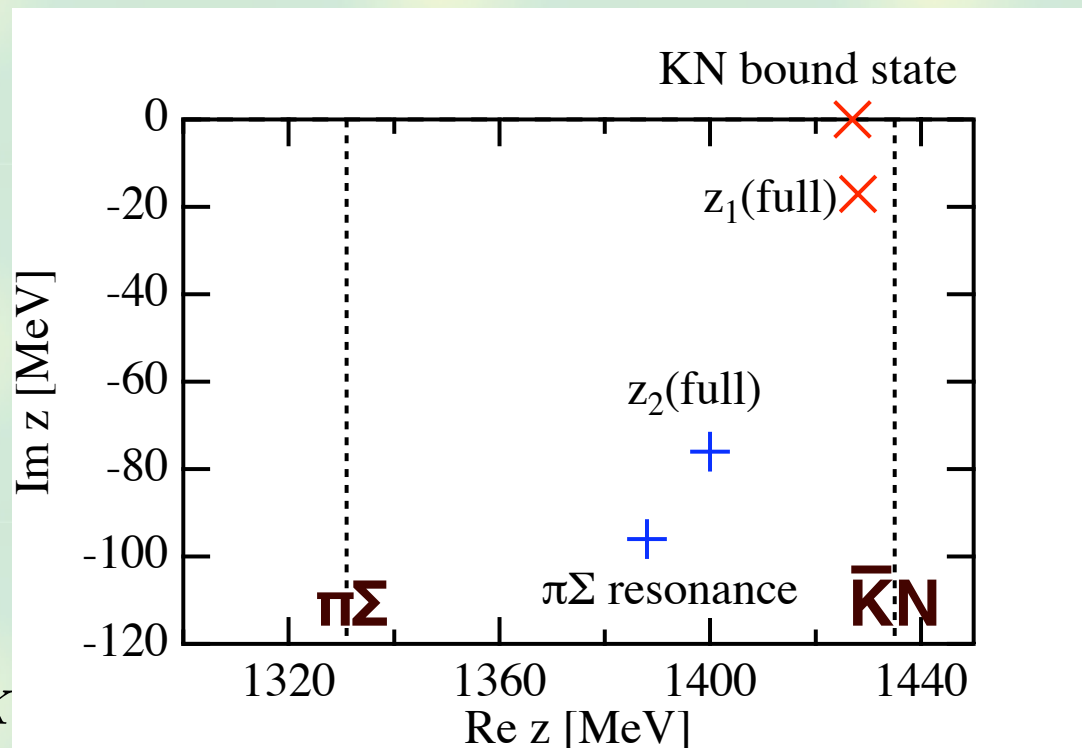
Origin of the two-pole structure

Chiral interaction

$$V_{ij} = -C_{ij} \frac{\omega_i + \omega_j}{4f^2}$$

$$C_{ij} = \begin{pmatrix} \bar{K}N & \pi\Sigma \\ 3 & -\sqrt{\frac{3}{2}} \\ -\sqrt{\frac{3}{2}} & 4 \end{pmatrix}$$

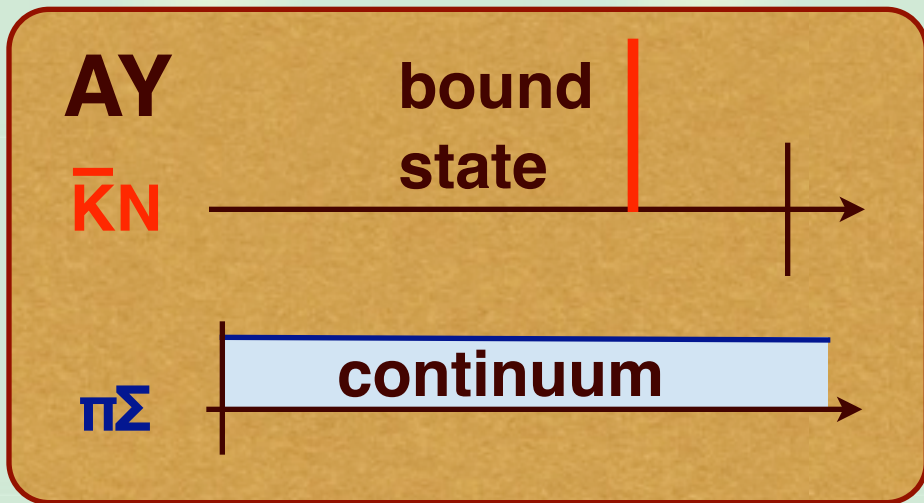
$$\omega_i \sim m_i, \quad 3.3m_\pi \sim m_K$$



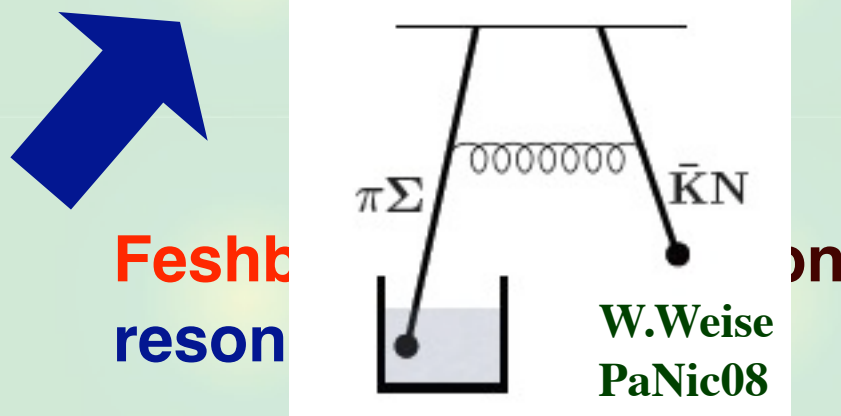
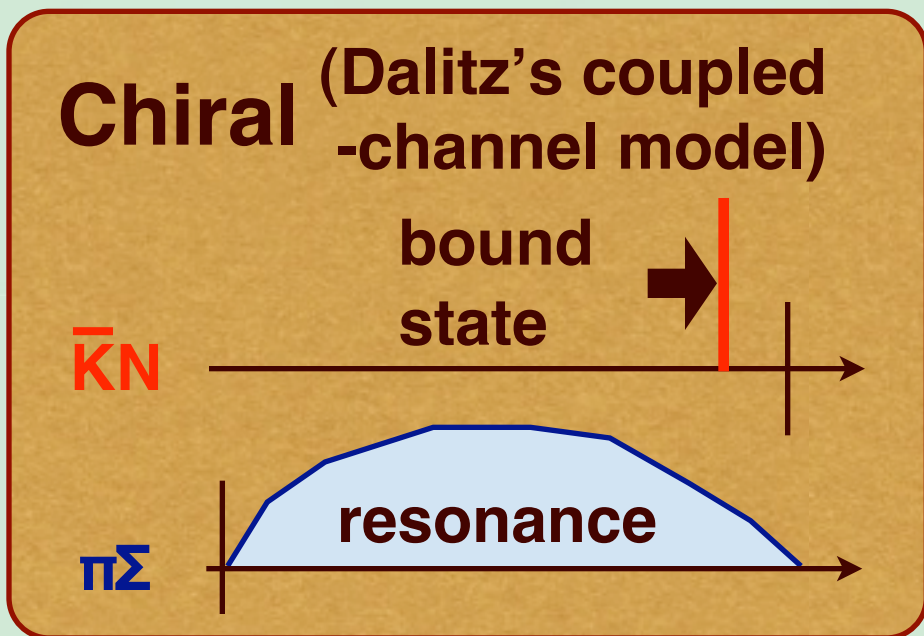
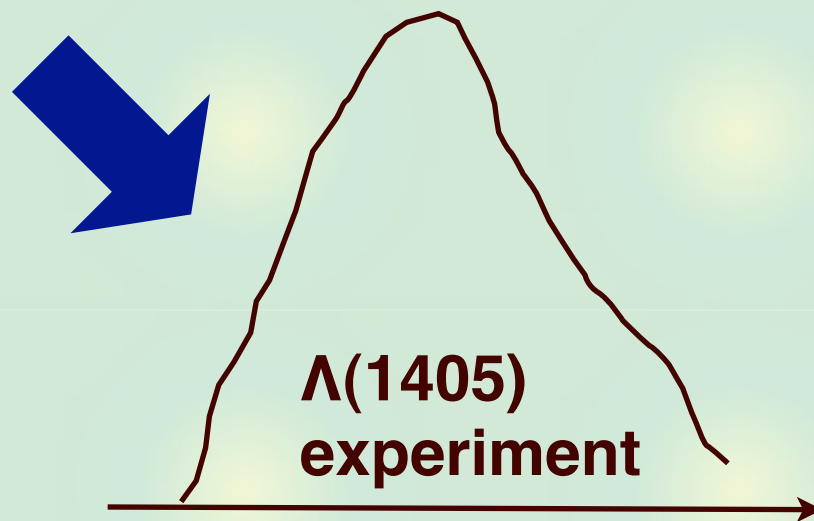
Very strong attraction in $\bar{K}N$ (higher energy) --> bound state
Strong attraction in $\pi\Sigma$ (lower energy) --> resonance

Two attractive interactions --> Two states
 $\pi\Sigma \rightarrow \pi\Sigma$ attraction : chiral SU(3) symmetry

Schematic illustration : AY vs Chiral



Feshbach resonance



Summary 2 : $\bar{K}N$ interaction

We study the consequence of chiral SU(3) dynamics in $\bar{K}N$ phenomenology.

- Single-channel effective $\bar{K}N$ interaction is **attractive** and forms K-pp bound system, about 20 MeV binding.
- Resonance structure in $\bar{K}N$ appears at around **1420 MeV** <-- **strong $\pi\Sigma$ dynamics**
- Two attractive interactions in $\bar{K}N$ and $\pi\Sigma$ --> **weaker** effective $\bar{K}N$ interaction.

[T. Hyodo and W. Weise, Phys. Rev. C 77, 035204 \(2008\)](#)

[A. Doté, T. Hyodo, W. Weise, Nucl. Phys. A804, 197 \(2008\)](#)

[A. Doté, T. Hyodo, W. Weise, Phys. Rev. C 79, 014003 \(2009\)](#)

Three-body calculations for $\bar{K}NN$

Single-channel variational calculation (DHW)

$$\text{B.E.} = 20 \pm 3 \text{ MeV}, \quad \Gamma(\pi YN) = 40 \sim 70 \text{ MeV}$$

Coupled-channel Faddeev calculation (IS)
using chiral interaction

Y. Ikeda and T. Sato, Phys. Rev. C 77, 035204 (2008)

$$\text{B.E.} \sim 79 \text{ MeV}, \quad \Gamma(\pi YN) \sim 74 \text{ MeV}$$

Why are they so different?

Inconsistency in theoretical calculations?

- $\pi\Sigma N$ dynamics?

(existence of another N when eliminating $\pi\Sigma$ channel)

Y. Ikeda and T. Sato, arXiv:0809.1285 [nucl-th]

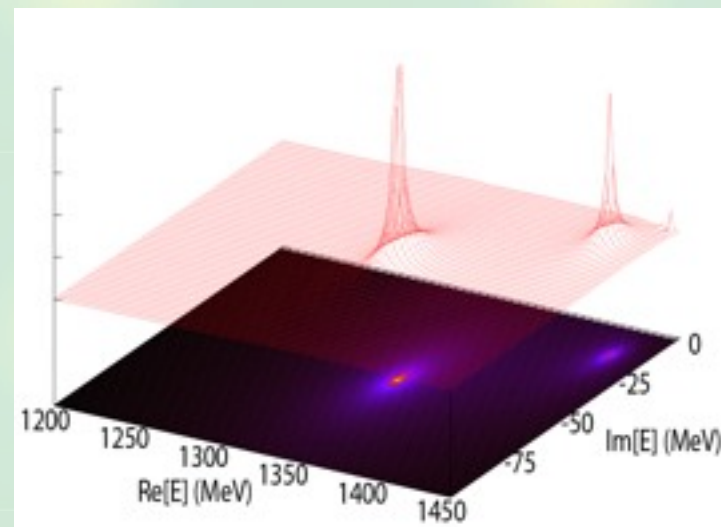
Two-pole structure for $\bar{K}NN$

Y. Ikeda, H. Kamano, T. Sato

Y. Ikeda, RCNP workshop, Dec. 25, 2008

“Chiral unitary like” model
: two poles in $\bar{K}N$ - $\pi\Sigma$ amplitude

Solving Faddeev equation, they find
two poles in $\bar{K}NN$ - $\pi\Sigma N$ amplitude



	pole I	pole II
\bar{K} - and N -exchanges in Z	$-14.5 - i28.7$	$-36.7 - i109.3$
\bar{K} -, N - and π -exchanges in Z	$-13.6 - i27.8$	$-45.8 - i104.0$
Full mechanism	$-13.7 - i29.0$	$-37.2 - i93.3$

--> Two poles for three-body system?

Two-pole structure for $\bar{K}NN$

Two poles in two-body scattering:

higher energy pole --> $\bar{K}N$ bound state

lower energy pole --> $\pi\Sigma$ resonance



If there are two poles in three-body system, “single-channel” approach of DHW focuses on the **higher energy pole of $\bar{K}NN$** , since the $\pi\Sigma N$ channel has been eliminated.

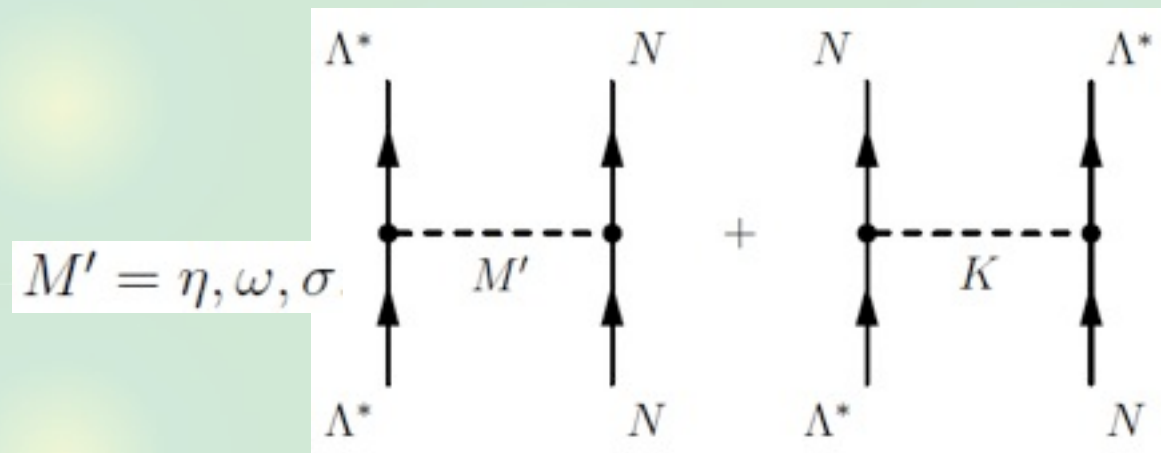
Then, where is the **lower energy state** in three-body system?

--> Schematic model calculation

Λ^* hypernuclei model

Treating $\Lambda(1405)$ as an elementary field,
 construct “ Λ^*N potential” through meson exchange

A. Arai, M. Oka and S. Yasui, Prog. Theor. Phys. 119, 103 (2008)



This approach may be justified by the observation that the Λ^* seems to be surviving in K-pp system.

Attractive interaction (mainly from σ exchange)
 --> bound Λ^*N , Λ^*NN systems

Λ^* coupling constant : unknown (<-- FINUDA data).

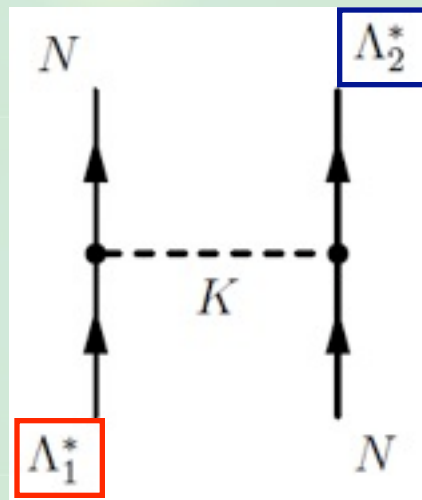
Λ^*N state in chiral model

Chiral dynamics \rightarrow two Λ^* states : Λ^*_1 , Λ^*_2

With sufficient attraction (σ exchange),

\rightarrow two Λ^*N bound states in $B=2$ system : Λ^*_1N , Λ^*_2N

In addition, mixing of $\Lambda^*_1N \leftrightarrow \Lambda^*_2N$: level repulsion

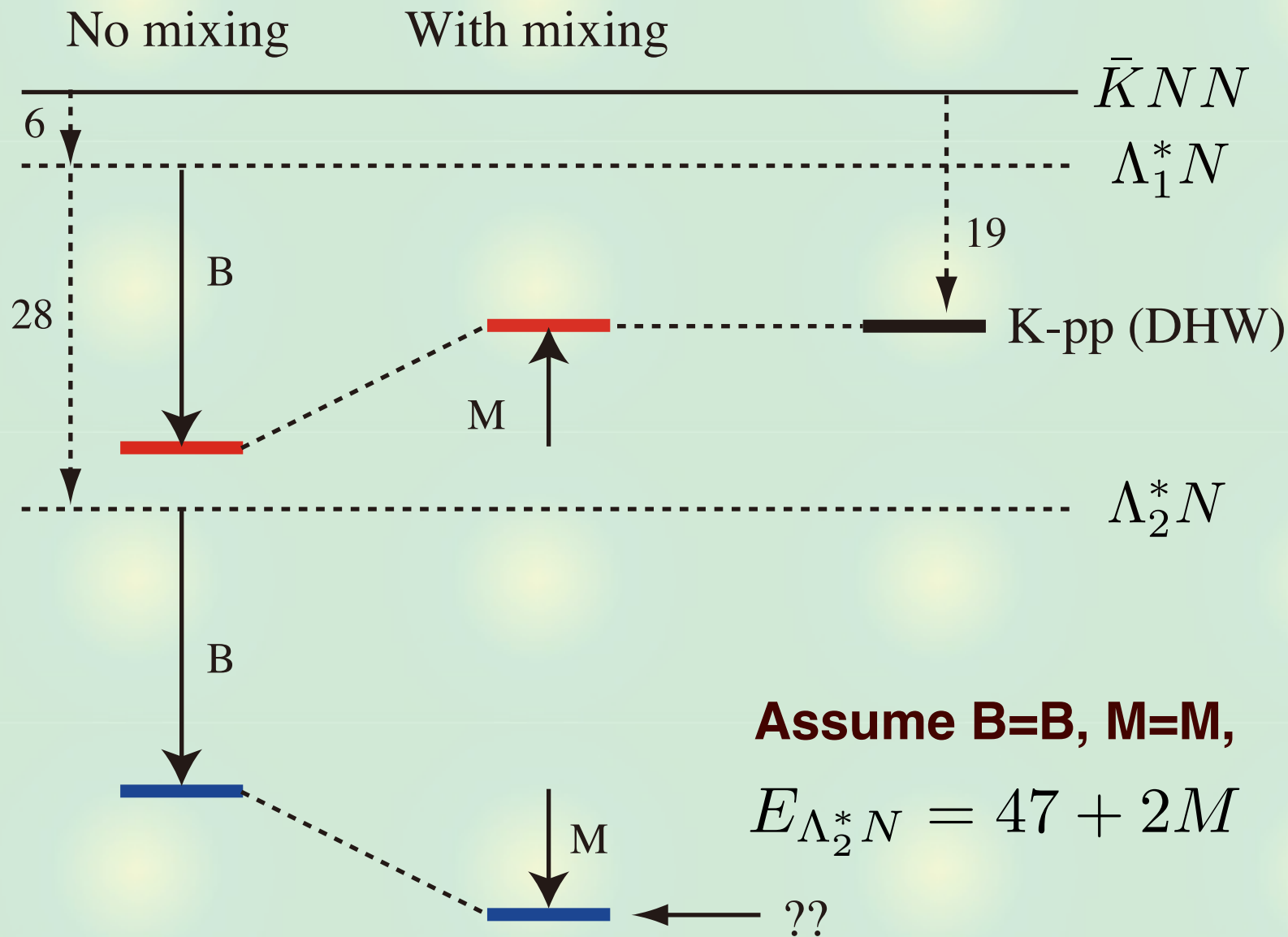


Λ^* coupling constant : unknown

We consider this model simulates the three-body calculation

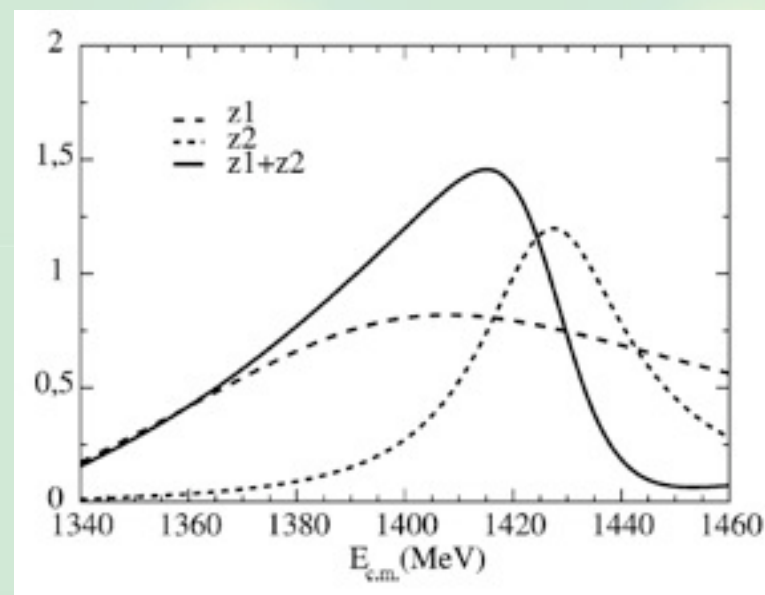
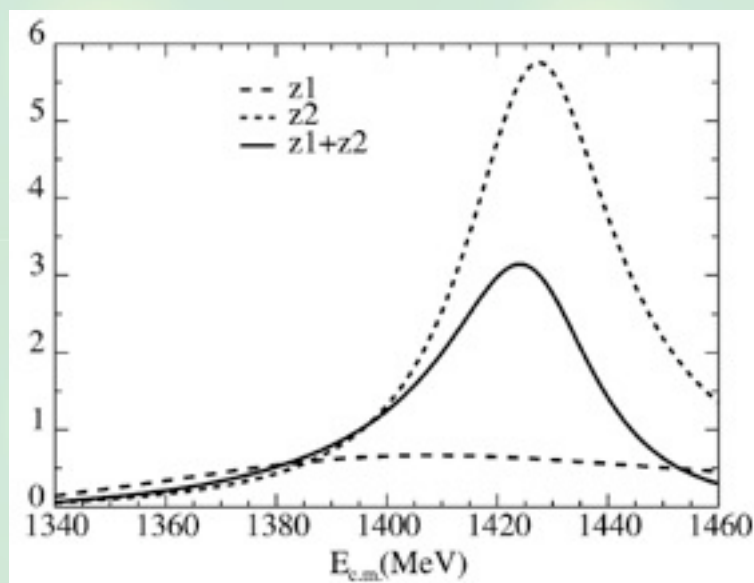
\rightarrow DHW result = Λ^*_1N

Λ^*N state in chiral model : result would be...



Λ^*N state in chiral modelChiral dynamics --> two Λ^* states : Λ^*_1 , Λ^*_2

$$|\Lambda(1405)\rangle = a \underbrace{|\Lambda^*_1\rangle}_{\bar{K}N} + b \underbrace{|\Lambda^*_2\rangle}_{\pi\Sigma}$$




D. Jido, et al, Nucl. Phys. Rev. A725, 181 (2003)

B=2 system : Λ^*_1N , Λ^*_2N

$$|B = 2, S = -1\rangle = a' |\Lambda^*_1 N\rangle + b' |\Lambda^*_2 N\rangle$$

Summary 3 : $\bar{K}NN$ system

B=2 and S=-1 system in chiral dynamics


 Possibility of **two poles** for three-body $\bar{K}NN$ state in Faddeev approach.

Y. Ikeda, RCNP workshop, Dec. 25, 2008

 If this is the case, DHW result may corresponds to higher energy Λ^*_1N state.

estimate of $\Lambda^*_2N \sim 47 + 2M$ MeV?

T. Uchino, T. Hyodo, M. Oka, in preparation

 When the lower energy channel is strongly interacting, coupled-channel approach would be mandatory.