Determining the Θ^+ quantum numbers through $K^+p \rightarrow \pi^+KN$





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Motivation : Spin parity determination

LEPS, T. Nakano, et al., Phys. Rev. Lett. 91, 012002 (2003)

 $|\Theta\rangle = |uudd\bar{s}\rangle$



S = +1, manifestly exotic

Spin and parity are not determined.

Find a reaction where qualitatively different results depending on the quantum numbers are observed.

Utilize the polarization of particles

Motivation : Advantage of hadronic process

We propose

$$K^+p \rightarrow \pi^+\Theta^+ \rightarrow \pi^+K^+n(K^0p)$$

at threshold of π and Θ .



- Low energy (k_{in} ~ 350 MeV in c.m. frame)
- Decay is considered -> background, interference,...
- Hadronic process : clear mechanism
 - : large cross section ~ $10^2 \,\mu b$

Chiral model for the reaction: Background

E. Oset and M. J. Vicente Vacas, PLB386, 39 (1996)

Threshold production of π and Θ Vertices <- chiral Lagrangian



Dominant

Proportional to $S \cdot p_{\pi^+}$ vanishes

Assume the final π^+ is almost at rest

Chiral model for the reaction: Resonance term



Chiral model for the reaction: Resonance term



Proportional to $\sigma \cdot p_{\pi^+}$ -> vanishes

Spin and parity : $KN \rightarrow \Theta \rightarrow KN$

1/2⁻ (KN s-wave resonance) $\stackrel{\bigvee}{\sim} M_R = 1540 \text{ MeV}$ $\Gamma_R = 20 \text{ MeV}$ 1/2⁺, 3/2⁺ (KN p-wave resonance)

 $t_{K^+n(K^0p)\to K^+n}^{(s)} = \frac{(\pm)g_{K^+n}^2}{M_I - M_R + i\Gamma/2}$ $t_{K^+n(K^0p)\to K^+n}^{(p,1/2)} = \frac{(\pm)\bar{g}_{K^+n}^2(\boldsymbol{\sigma}\cdot\boldsymbol{q}')(\boldsymbol{\sigma}\cdot\boldsymbol{q})}{M_I - M_R + i\Gamma/2}$ $t_{K^{+}n(K^{0}p)\to K^{+}n}^{(p,3/2)} = \frac{(\pm)\tilde{g}_{K^{+}n}^{2}(\boldsymbol{S}\cdot\boldsymbol{q}')(\boldsymbol{S}^{\dagger}\cdot\boldsymbol{q})}{M_{I}-M_{R}+i\Gamma/2}$ $g_{K+n}^2 = \frac{\pi M_R \Gamma}{Mq} , \quad \overline{g}_{K+n}^2 = \frac{\pi M_R \Gamma}{Mq^3} , \quad \widetilde{g}_{K+n}^2 = \frac{3\pi M_R \Gamma}{Mq^3}$

Spin and parity : Resonance amplitude

Resonance term for $K^+p \rightarrow \pi^+K^+n$

$$-i\tilde{t}_{i}^{(s)} = \frac{g_{K+n}^{2}}{M_{I} - M_{R} + i\Gamma/2} \left\{ G(M_{I})(a_{i} + c_{i}) - \frac{1}{3}\bar{G}(M_{I})b_{i} \right\} \boldsymbol{\sigma} \cdot \boldsymbol{k}_{in} S_{I}(i)$$

$$-i\tilde{t}_{i}^{(p,1/2)} = \frac{\bar{g}_{K+n}^{2}}{M_{I} - M_{R} + i\Gamma/2} \bar{G}(M_{I}) \left\{ \frac{1}{3}b_{i}\boldsymbol{k}_{in}^{2} - a_{i} + d_{i} \right\} \boldsymbol{\sigma} \cdot \boldsymbol{q}' S_{I}(i)$$

$$-i\tilde{t}_{i}^{(p,3/2)} = \frac{\tilde{g}_{K+n}^{2}}{M_{I} - M_{R} + i\Gamma/2} \bar{G}(M_{I}) \frac{1}{3}b_{i} \left\{ (\boldsymbol{k}_{in} \cdot \boldsymbol{q}')(\boldsymbol{\sigma} \cdot \boldsymbol{k}_{in}) - \frac{1}{3}\boldsymbol{k}_{in}^{2}\boldsymbol{\sigma} \cdot \boldsymbol{q}' \right\} S_{I}(i)$$



Numerical results : Mass distributions

$$I,J^{P}=0,1/2^{-}$$

--- $I,J^{P}=0,1/2^{+}$ $k_{in}(Lab) = 850 \text{ MeV/c}$
--- $I,J^{P}=0,3/2^{+}$ $\theta = 0 \text{ deg}$



Numerical results : Mass distributions



Deviation of the peak from B.W. mass -> interference effect size ~ width

Production 1 : $K^+p \rightarrow \pi^+K^+n$

Numerical results : Angular dependence



Numerical results : Polarization test





Z direction

$\langle -1/2 | \boldsymbol{\sigma} \cdot \boldsymbol{k}_{in} | 1/2 \rangle = 0$ $\langle -1/2 | \boldsymbol{\sigma} \cdot \boldsymbol{q}' | 1/2 \rangle \propto q' \sin \theta$

Same result is obtained for final pK^0

Numerical results : Mass distributions



Numerical results : Angular dependence



Numerical results : Incomplete polarization



Conclusion

We calculate the $K^+p \rightarrow \pi^+K^+n$ reaction using a chiral model, assuming the possible quantum numbers of Θ^+ baryon.

If we find the resonance in the polarization test, the quantum numbers of Θ⁺ can be determined as I=0, J^P=1/2⁺

T. Hyodo, A. Hosaka, E. Oset, Phys. Lett. B579, 290 (2004) T. Hyodo, Doctor Thesis (2006) **Problems & future work**

Problems 0 momentum π polarization of final N



E. Oset, M. J. Vicente Vacas, PLB386, 39

 $p_{lab} (MeV/c)$