

Phenomenology of spin $3/2$ baryons with pentaquarks



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Introduction : Flavor SU(3) symmetry

Existence of Θ^+ + Flavor SU(3) symmetry

➔ Existence of **flavor partners** of Θ^+

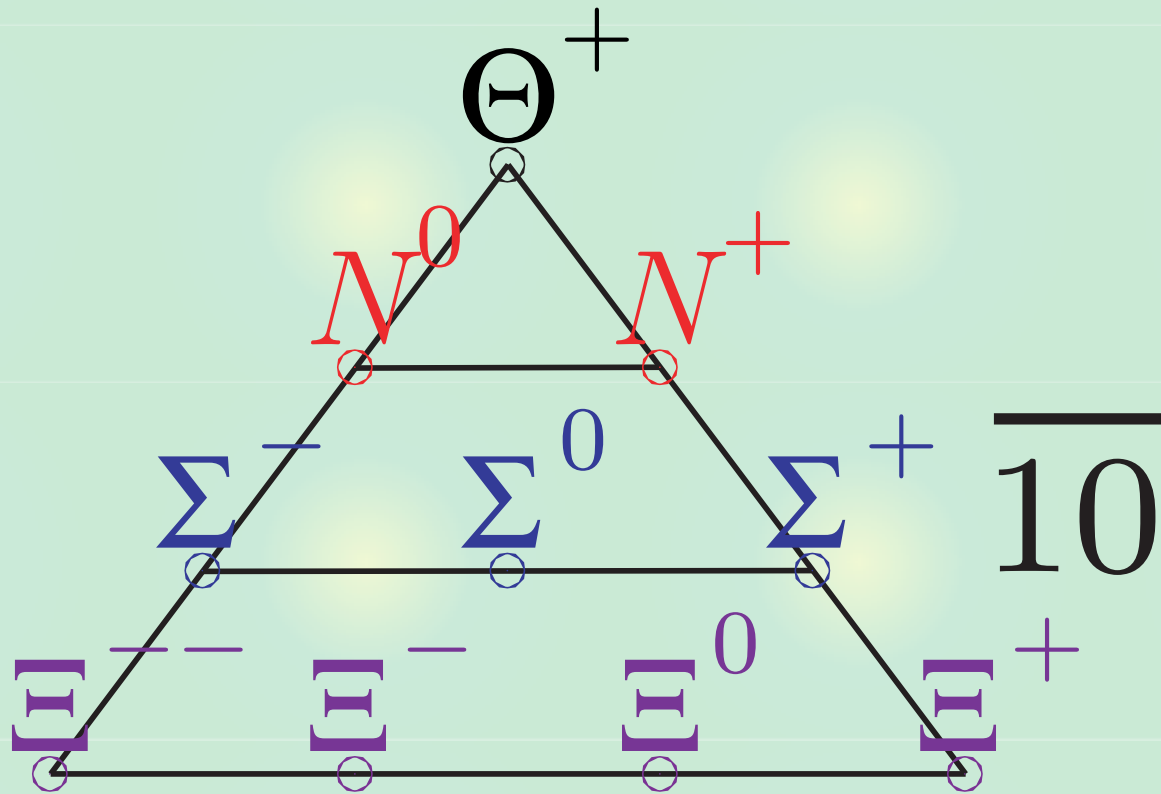
Assuming the flavor multiplet that Θ^+ belongs to, we examine its properties by symmetry relation, in connection with known baryon resonances.

➔ to determine the J^P of Θ^+

Phenomenological but model independent analysis up to $O(m_s)$

Pure antidecuplet case

Simplest assignment for Θ^+



Test the masses and widths of partners via flavor SU(3) symmetry relations

Pure antidecuplet case

Mass and decay width [MeV]

$$M(\overline{10}; Y) = M_{\overline{10}} - aY$$

$$g_{\Theta KN} = \sqrt{6}g_{N^* \pi N}$$

J^P	M_{Θ}	M_N	M_{Σ}	M_{Ξ}	Γ_{Θ}
$1/2^-$ exp.	1540 $\Theta(1540)$	1647 N(1650)	1753 $\Sigma(1750)$	1860 $\Xi(1860)$	156.1
$1/2^+$ exp.					
$3/2^+$ exp.					
$3/2^-$ exp.					

Pure antidecuplet case

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$3/2^+$ exp.					
$3/2^-$ exp.					

Pure antidecuplet case

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Pure antidecuplet case

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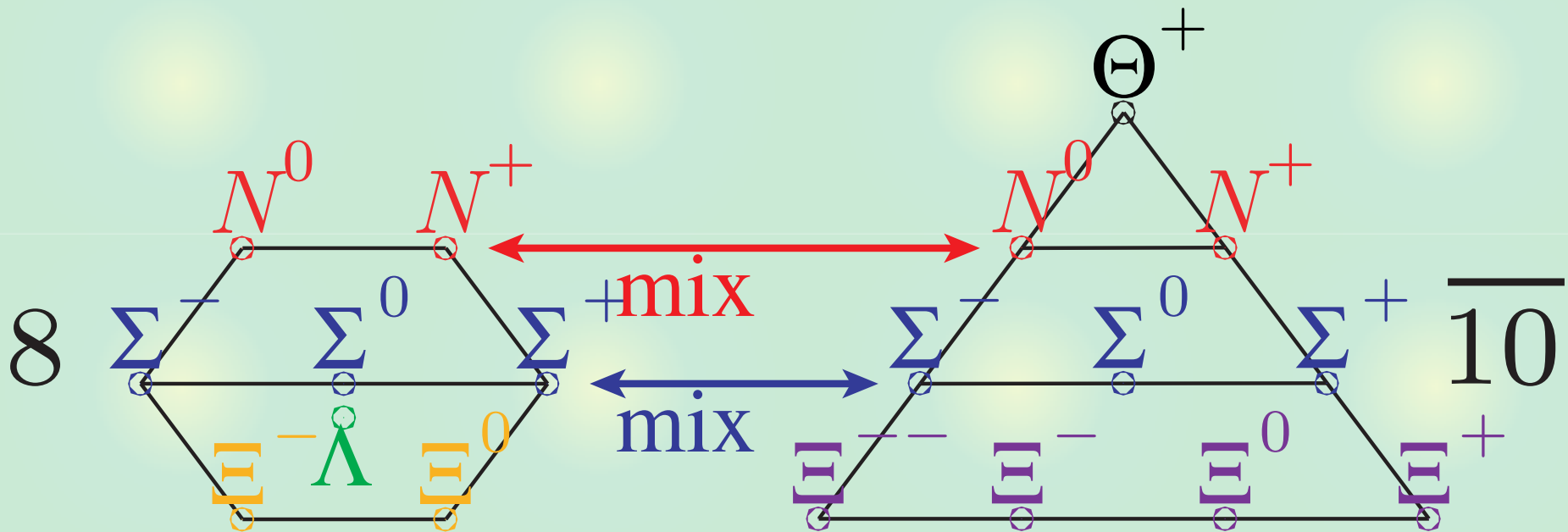
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$3/2^-$ exp.	1540 $\Theta(1540)$	1700 N(1700)	1860	2020 $\Xi(2030)$	1.3

are not reproduced simultaneously.

Octet-antidecuplet mixing

Second simplest assignment for Θ^+



Mixing is induced by the $SU(3)$ breaking in mass term.

Octet-antidecuplet mixing

Mass formulae

$$M_{\Theta} = M_{\overline{10}} - 2a$$

$$M_{\Xi_{\overline{10}}} = M_{\overline{10}} + a$$

$$M_{\Lambda} = M_{\mathbf{8}}$$

$$M_{\Xi_{\mathbf{8}}} = M_{\mathbf{8}} + b + \frac{1}{2}c$$

$$M_{N_1} = \left(M_{\mathbf{8}} - b + \frac{1}{2}c \right) \cos^2 \theta_N + (M_{\overline{10}} - a) \sin^2 \theta_N - \delta \sin 2\theta_N$$

$$M_{N_2} = \left(M_{\mathbf{8}} - b + \frac{1}{2}c \right) \sin^2 \theta_N + (M_{\overline{10}} - a) \cos^2 \theta_N + \delta \sin 2\theta_N$$

$$M_{\Sigma_1} = (M_{\mathbf{8}} + 2c) \cos^2 \theta_{\Sigma} + M_{\overline{10}} \sin^2 \theta_{\Sigma} - \delta \sin 2\theta_{\Sigma}$$

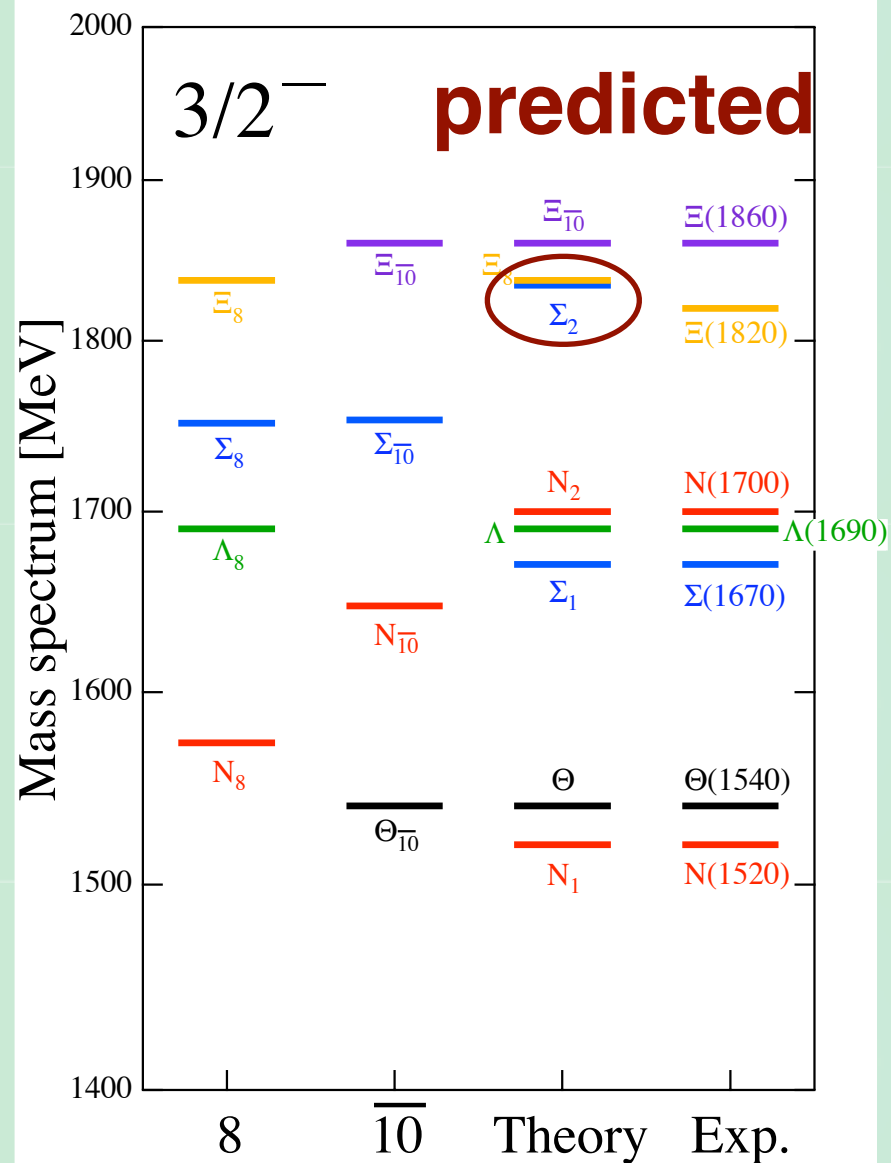
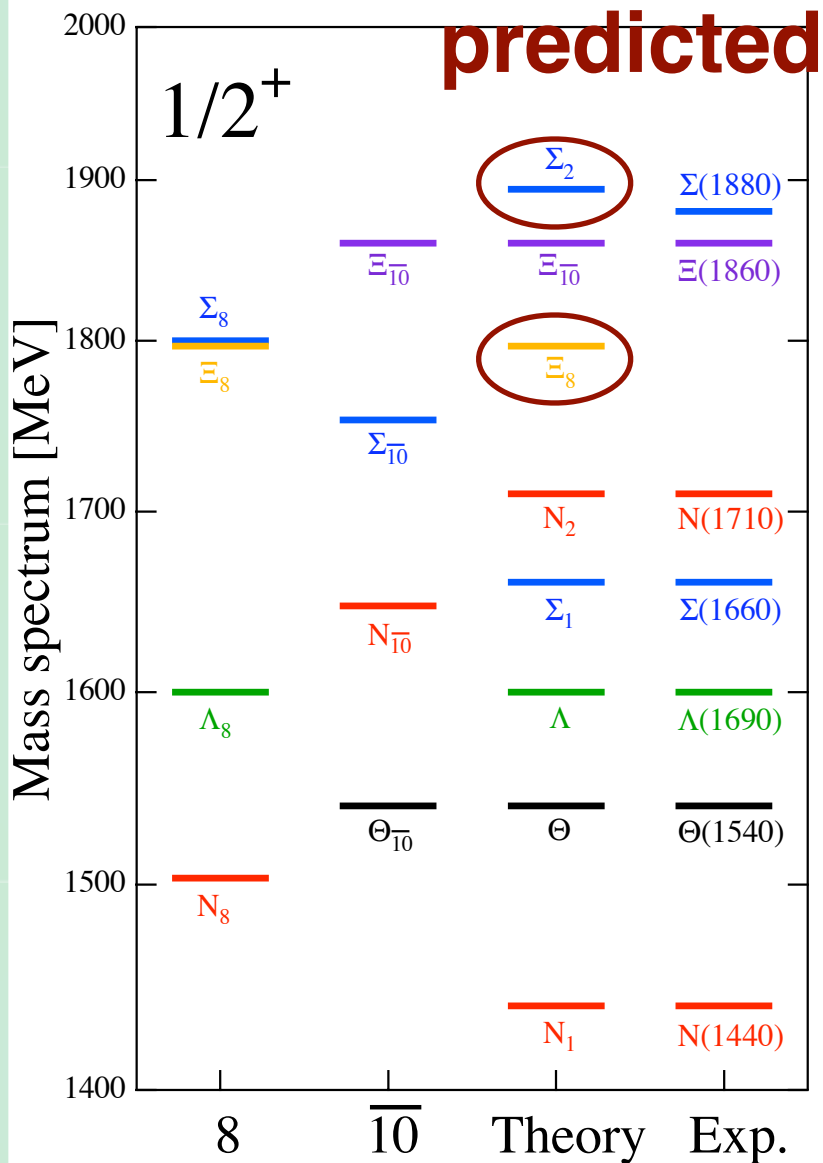
$$M_{\Sigma_2} = (M_{\mathbf{8}} + 2c) \sin^2 \theta_{\Sigma} + M_{\overline{10}} \cos^2 \theta_{\Sigma} + \delta \sin 2\theta_{\Sigma}$$

8 masses v.s. 6 parameters

$J^P = 1/2^-$: too wide width

$J^P = 3/2^+$: states are not well established

Mass spectra



Decay width of Θ

Relation between coupling constants

$$g_{N_1} = g_{N_8} \cos \theta_N - \frac{g_{10}}{\sqrt{6}} \sin \theta_N$$

$$g_{N_2} = \frac{g_{10}}{\sqrt{6}} \cos \theta_N + g_{N_8} \sin \theta_N$$

J^P	θ_N [deg]	Γ_Θ [MeV]
$1/2^+$	29	29.1
$3/2^-$	33	3.1

c.f. ideal mixing ~ 35 deg

Two-meson coupling

Contact interaction :



A. Hosaka, T. H., et al., Phys. Rev. C71 045205 (2005)

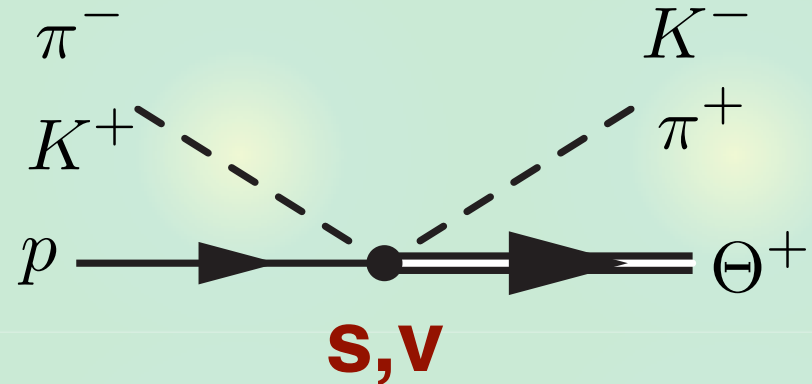
Branching fraction [%]

J^P	state	πN	$\pi\pi N(s)$	$\pi\pi N(v)$
$1/2^+$	N(1440)	65	7.5	<8
	N(1710)	15	25	15
$3/2^-$	N(1520)	55	25	20
	N(1700)	10	-	<35

Two-meson coupling

SU(3) relation enable us to calculate

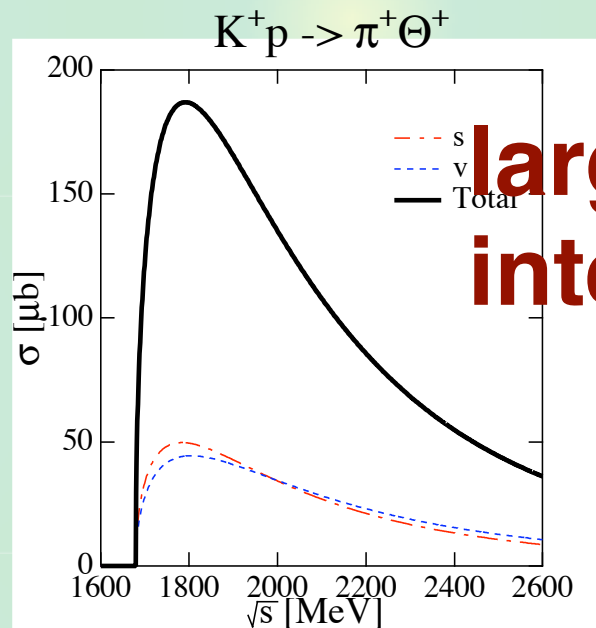
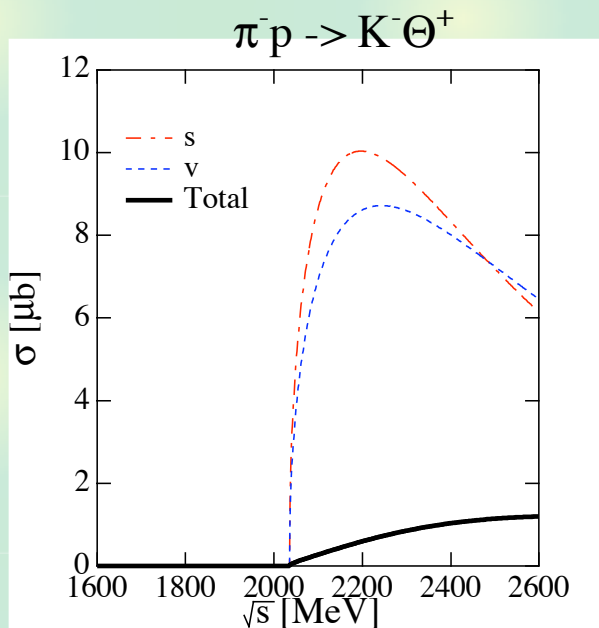
the cross section of



Two different channels should be summed coherently
-> interference effect

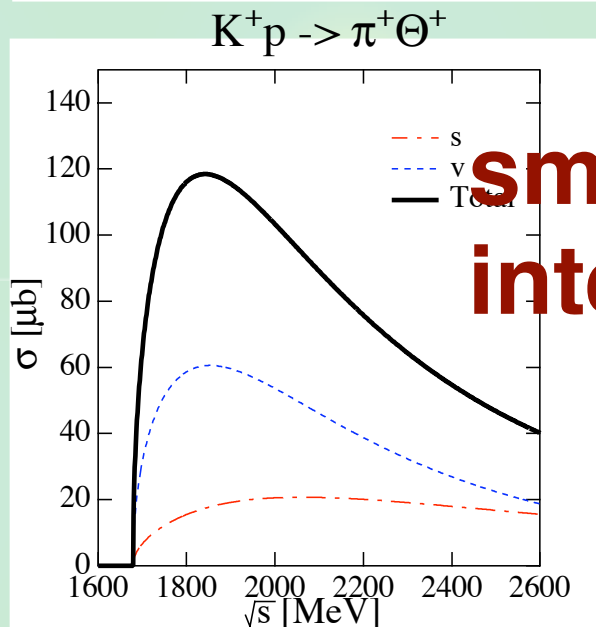
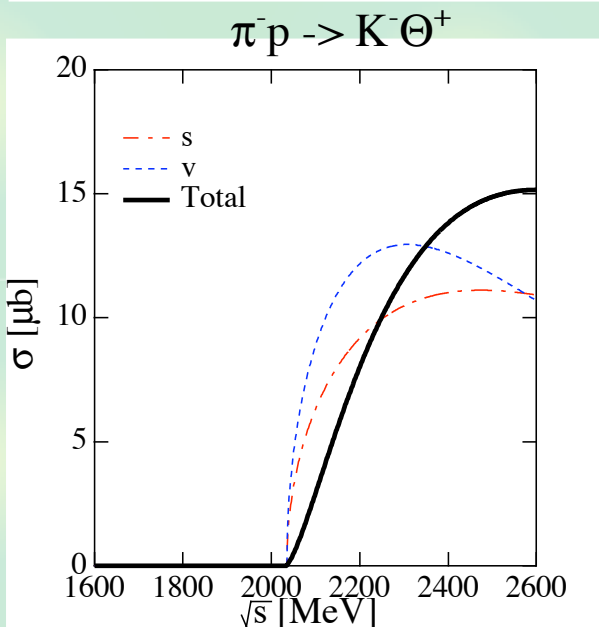
Two-meson coupling

$1/2^+$



**large
interference**

$3/2^-$



**small
interference**

Conclusion 1 : mixing scheme

We examine $8-\overline{10}$ mixing scheme for the exotic and non-exotic baryon resonances.

- Masses of $\Theta(1540)$ and $\Xi(1860)$ are well fitted in the $8-\overline{10}$ mixing scheme with $J^P = 1/2^+$ or $3/2^-$ baryons.
- The width of Θ is **very narrow** for the $J^P = 3/2^-$ case.
- For both cases, the mixing angle is close to the **ideal angle**.

Conclusion 2 : Θ production

Based on the mixing scheme, we evaluate the two-meson coupling of Θ , and calculate the reaction process for Θ production



There is an **interference effect** between two amplitudes, which is prominent for $1/2+$ case and rather moderate for $3/2-$ case