Photoproduction of K* for the study of $\Lambda(1405)$





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Motivation : Two poles?

There are two poles of the scattering amplitude around nominal $\Lambda(1405)$ energy region.

- <u>Cloudy bag model</u> (1990) J. Fink, *et al.* PRC41, 2720
- <u>Chiral unitary model</u> (2001~)

J. A. Oller, *et al.* PLB500, 263 E. Oset, *et al.* PLB527, 99 D. Jido, *et al.* PRC66, 025203 T. Hyodo, *et al.* PRC68, 018201 C. Garcia-Recio, *et al.*, PRD67, 076009 D. Jido, *et al.*, NPA725, 181 T. Hyodo, *et al.* PRC68, 065203 $\Lambda(1405):J^P=1/2^-,I=0$



ChU model, T. Hyodo

Flavor SU(3) meson-baryon scatterings (s-wave)



Dynamical generation

 $J^P = 1/2^-$ resonances

 $egin{aligned} \Lambda(1405), \Lambda(1670), \ \Sigma(1620), \Xi(1620), \ N(1535) \end{aligned}$

Framework of the chiral unitary model



Total cross sections of K-p scattering



Trajectories of the poles with SU(3) breaking (S=-1)



D. Jido, et al., Nucl. Phys. A 723, 205 (2003)

$\Lambda(1405)$ in the chiral unitary model



D. Jido, et al., Nucl. Phys. A 723, 205 (2003)

Photoproduction of K* and $\Lambda(1405)$

In order to study



we calculate

$p \rightarrow K^* \Lambda(1405)$

Photoproduction of K* and $\Lambda(1405)$



Advantage of this reaction



(1)
$$\varepsilon_{\mu}(K^{*}) // \varepsilon_{\mu}(\gamma)$$
 : J^{P} = natural

(2)
$$\varepsilon_{\mu}(K^{*}) \perp \varepsilon_{\mu}(\gamma) : J^{P} = \text{unnatural}$$

With polarized photon beam, the exchanged particle can be identified. Clear mechanism

Effective interaction for meson part

 $\widehat{\boldsymbol{\chi}} \quad \widehat{\boldsymbol{\chi}} \quad \mathbf{K} \quad \mathbf{K} \quad \mathbf{Coupling}$ $\mathcal{L}_{K^*K\gamma} = g_{K^*K\gamma} \epsilon^{\mu\nu\alpha\beta} \partial_{\mu} A_{\nu} (\partial_{\alpha} K_{\beta}^{*-} K^+ + \partial_{\alpha} \bar{K}_{\beta}^{*0} K^0) + \text{h.c.}$



VPP coupling

$$\mathcal{L}_{VPP} = -\frac{ig_{VPP}}{\sqrt{2}} \operatorname{Tr}(V^{\mu}[\partial_{\mu}P, P])$$



Effective interaction for baryon part





ChU amplitude

Σ(1385) pole term

$\simeq \Sigma(1385)$ MB coupling

$$\overleftarrow{\quad} -it_{\Sigma^*i} = c_i \frac{12}{5} \frac{D+F}{2f} \mathbf{S} \cdot \mathbf{k}_i$$

$\begin{array}{l} \checkmark \quad \text{form factor} \\ F_f(k_1) = \frac{\Lambda^2 - m_K^2}{\Lambda^2 - (k_1)^2} \end{array}$

Total cross sections



Invariant mass distributions



Isospin decomposition of \pi\Sigma states

Since initial state is KN, we neglect the I=2.

$$rac{d\sigma(\pi^0\Sigma^0)}{dM_I} \propto rac{1}{3} |T^{(0)}|^2$$
 • Pure I=0 amplitude

$$\frac{d\sigma(\pi^{\pm}\Sigma^{\mp})}{dM_{I}} \propto \frac{1}{3} |T^{(0)}|^{2} + \frac{1}{2} |T^{(1)}|^{2} \pm \frac{2}{\sqrt{6}} \operatorname{Re}(T^{(0)}T^{(1)*})$$

- Difference among charged states
 -> when summed up, this term vanishes
- No p-wave contribution
 -> I=1 s-wave amplitude

Invariant mass distributions 2



I=1, s-wave amplitude



Summary and conclusions 1

We study the structure of $\Lambda(1405)$ using the chiral unitary model.

There are two poles of the scattering amplitude around nominal $\Lambda(1405)$. Pole 1 (1426+16i) : strongly couples to KN state Pole 2 (1390+66i) : strongly couples to $\pi\Sigma$ state By observing the charged $\pi\Sigma$ states in the $\gamma p \rightarrow K^* \Lambda(1405)$ reaction, it is possible to isolate the higher energy pole.

Summary and conclusions 2

If we observe neutral πΣ state, clear
I=0 distribution is obtained.

Combining three πΣ states, we can also study the s-wave I=1 amplitude, where the existence of another pole is argued.

T. H., A. Hosaka, E. Oset, M. J. Vicente Vacas, nucl-th/0401051

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Appendix : other processes

 $\gamma p o K^- \pi \Sigma$

J.C. Nacher, et al., PLB445, 55



