


Photoproduction of K^ for the study of $\Lambda(1405)$*



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Motivation : Two poles?

There are two poles of the scattering amplitude around nominal $\Lambda(1405)$ energy region.

- Cloudy bag model
(1990)

J. Fink *et al.* PRC41, 2720

- Chiral unitary model
(2001~)

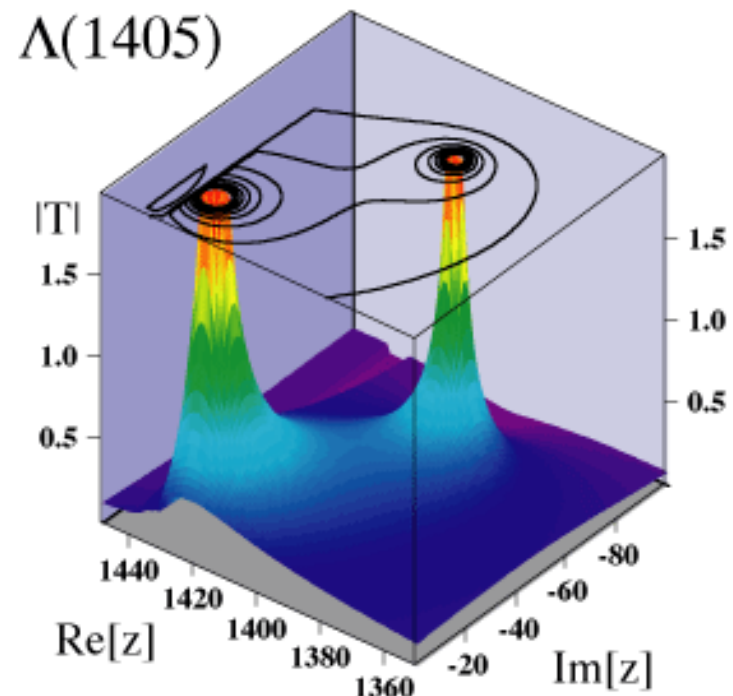
J. A. Oller *et al.* PLB500, 263

E. Oset *et al.* PLB527, 99

D. Jido *et al.* PRC66, 025203

T. Hyodo *et al.* PRC68, 018201

$\Lambda(1405) : J^P=1/2^-, I=0$



ChU model, T. Hyodo

Chiral unitary model

Flavor SU(3) meson-baryon scatterings (s-wave)

Chiral symmetry

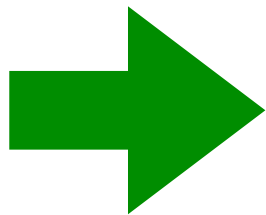
**Low energy
behavior**



Unitarity of S-matrix

**Non-perturbative
resummation**

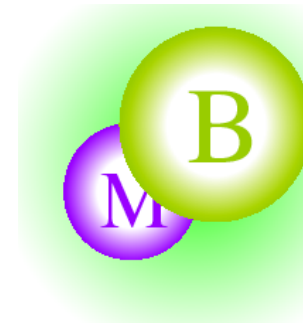
**Dynamical
generation**



$$J^P = 1/2^-$$

Resonances

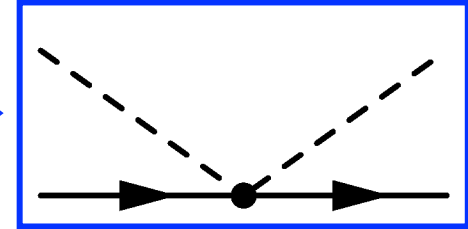
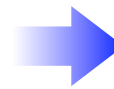
$\Lambda(1405)$, $\Lambda(1670)$, $N(1535)$,
 $\Sigma(1620)$, $\Xi(1620)$



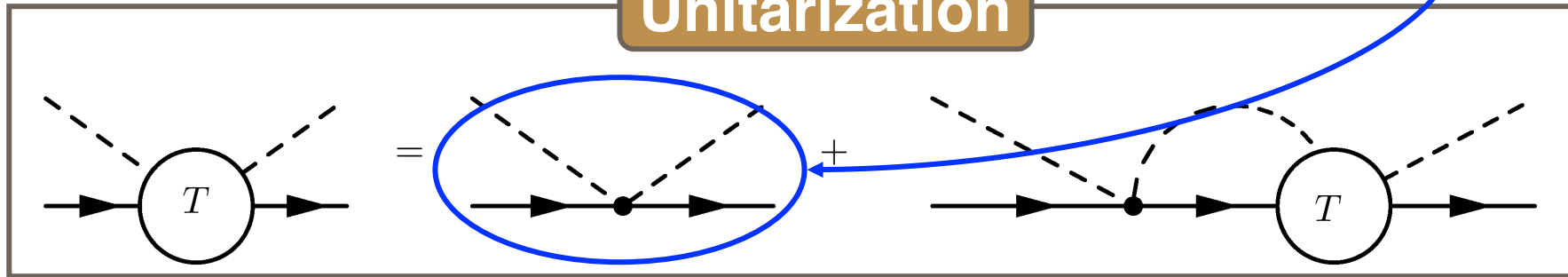
Framework of the chiral unitary model

Chiral perturbation theory

$$\mathcal{L}_{WT} = \frac{1}{4f^2} \text{Tr}(\bar{B}i\gamma^\mu[(\Phi\partial_\mu\Phi - \partial_\mu\Phi\Phi), B])$$

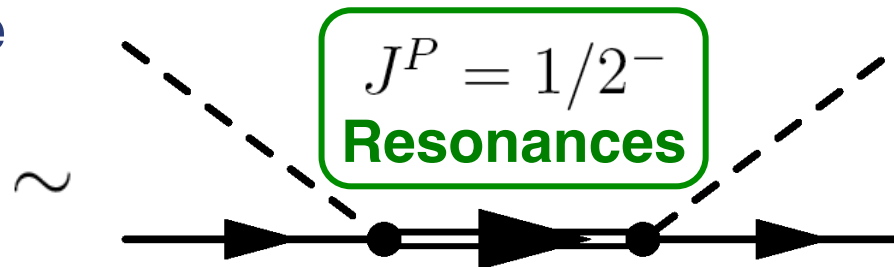


Unitarization

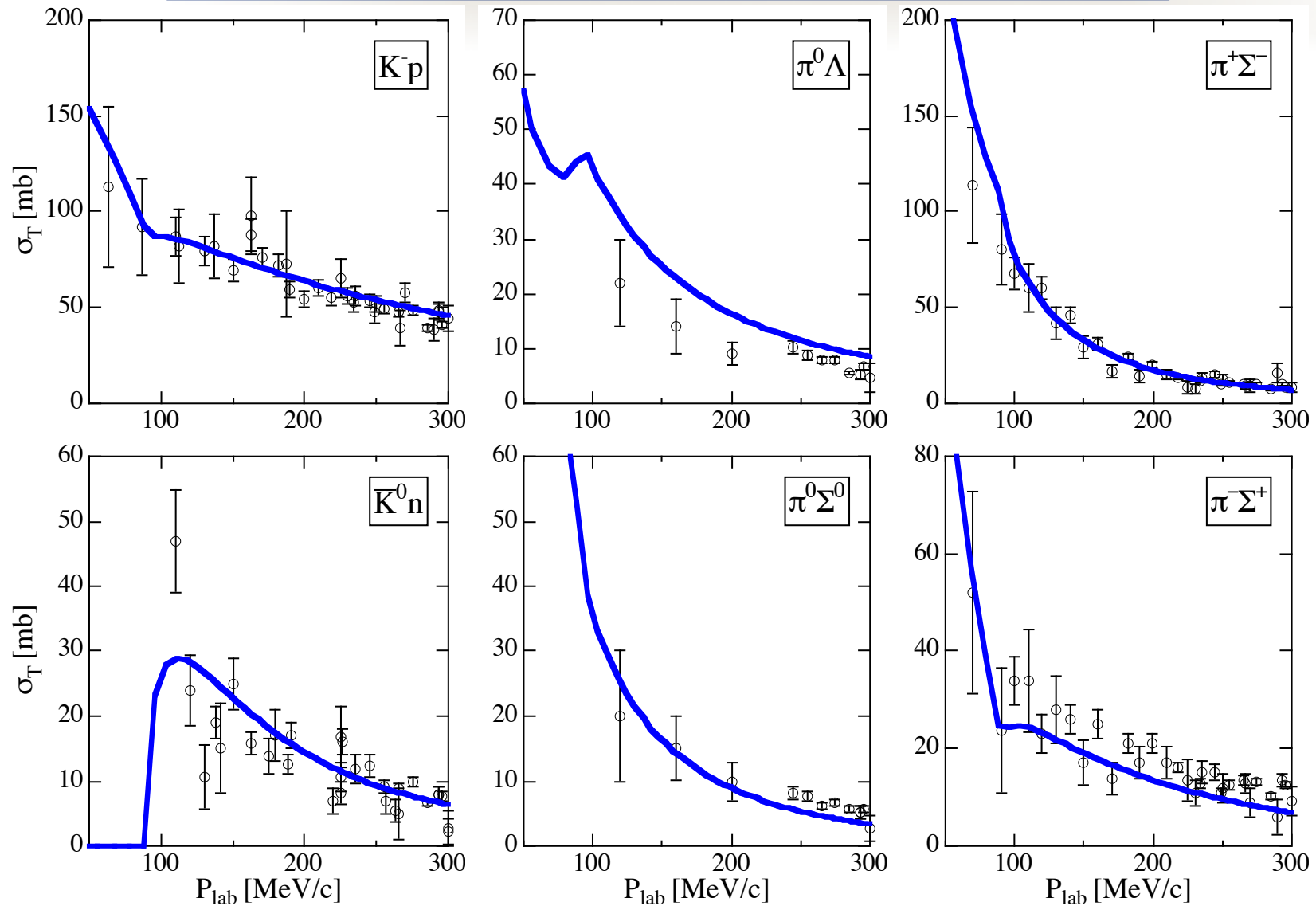


$$T_{ij}(\sqrt{s}) \sim \frac{g_i g_j}{\sqrt{s} - M_R + i\Gamma_R/2} + T_{ij}^{BG}$$

Generated resonances are expressed as poles of the scattering amplitude.

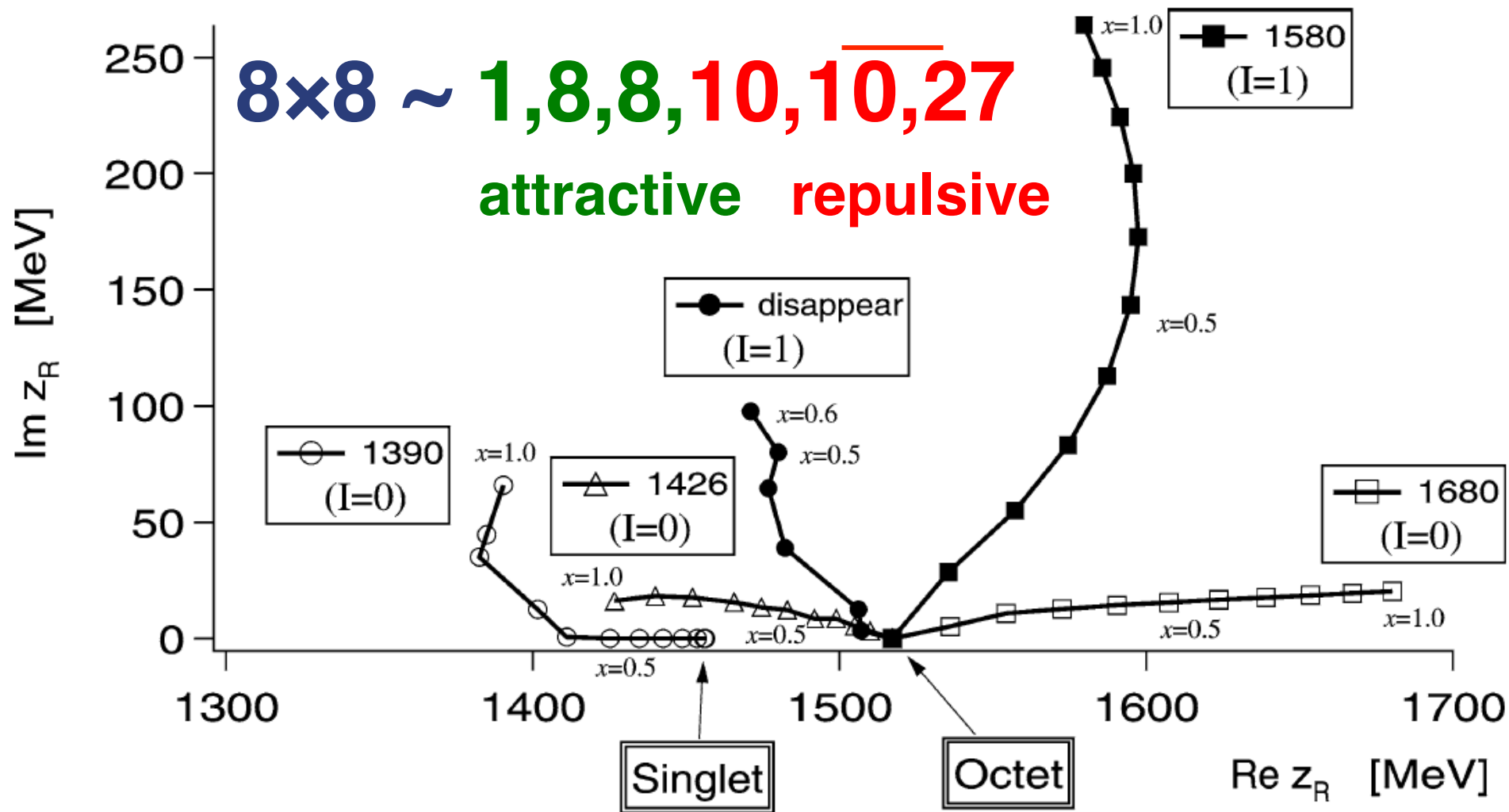


Total cross sections of K^-p scatterings



T. Hyodo, et al., Phys. Rev. C 68, 018201 (2003)

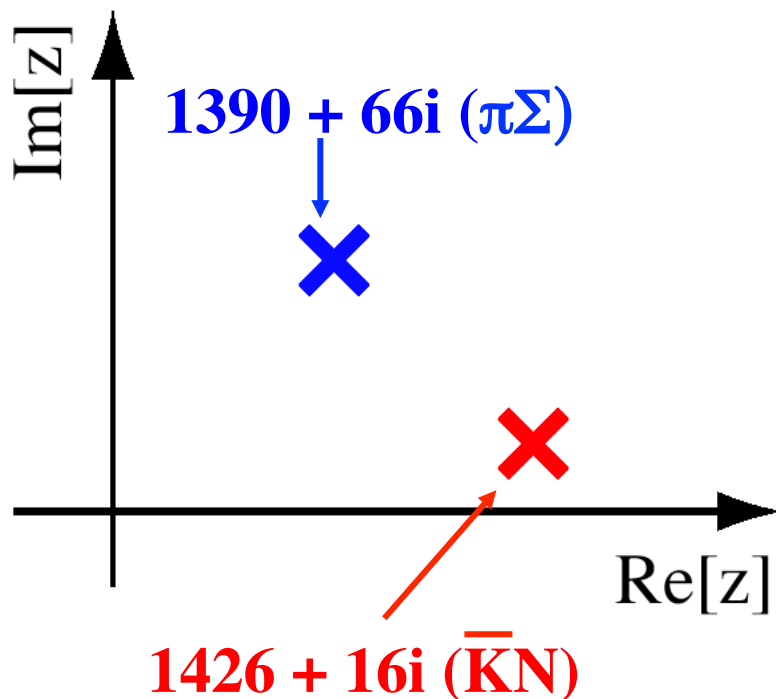
Poles of amplitude in $S=-1$ channel



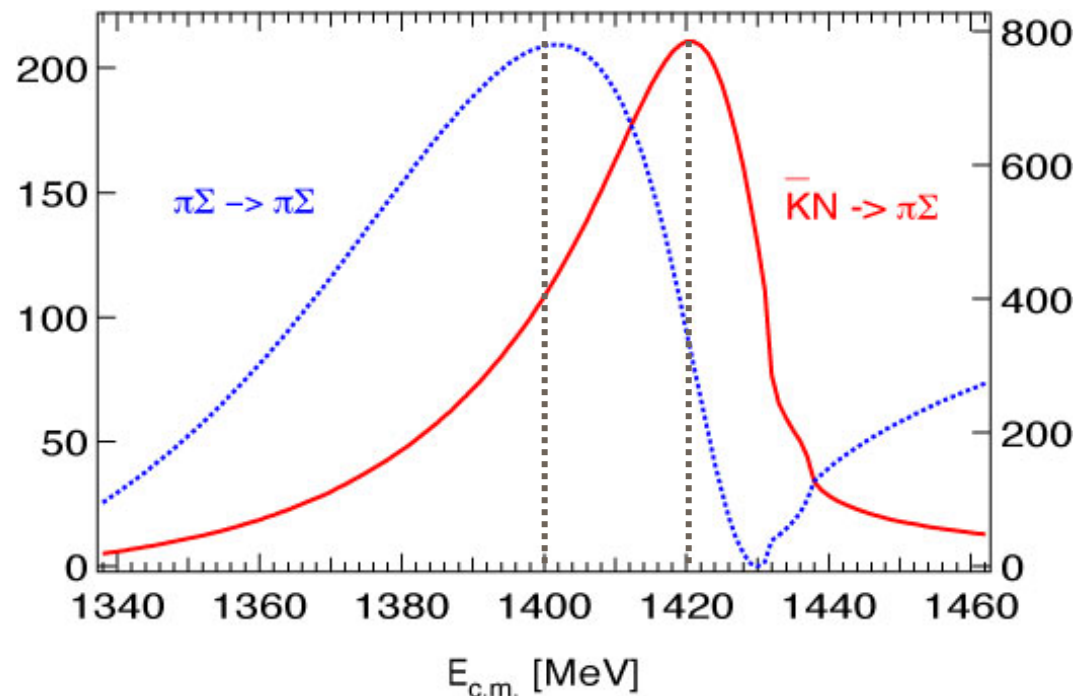
D. Jido, et al., Nucl. Phys. A 723, 205 (2003)

$\Lambda(1405)$ in the chiral unitary model

position of poles



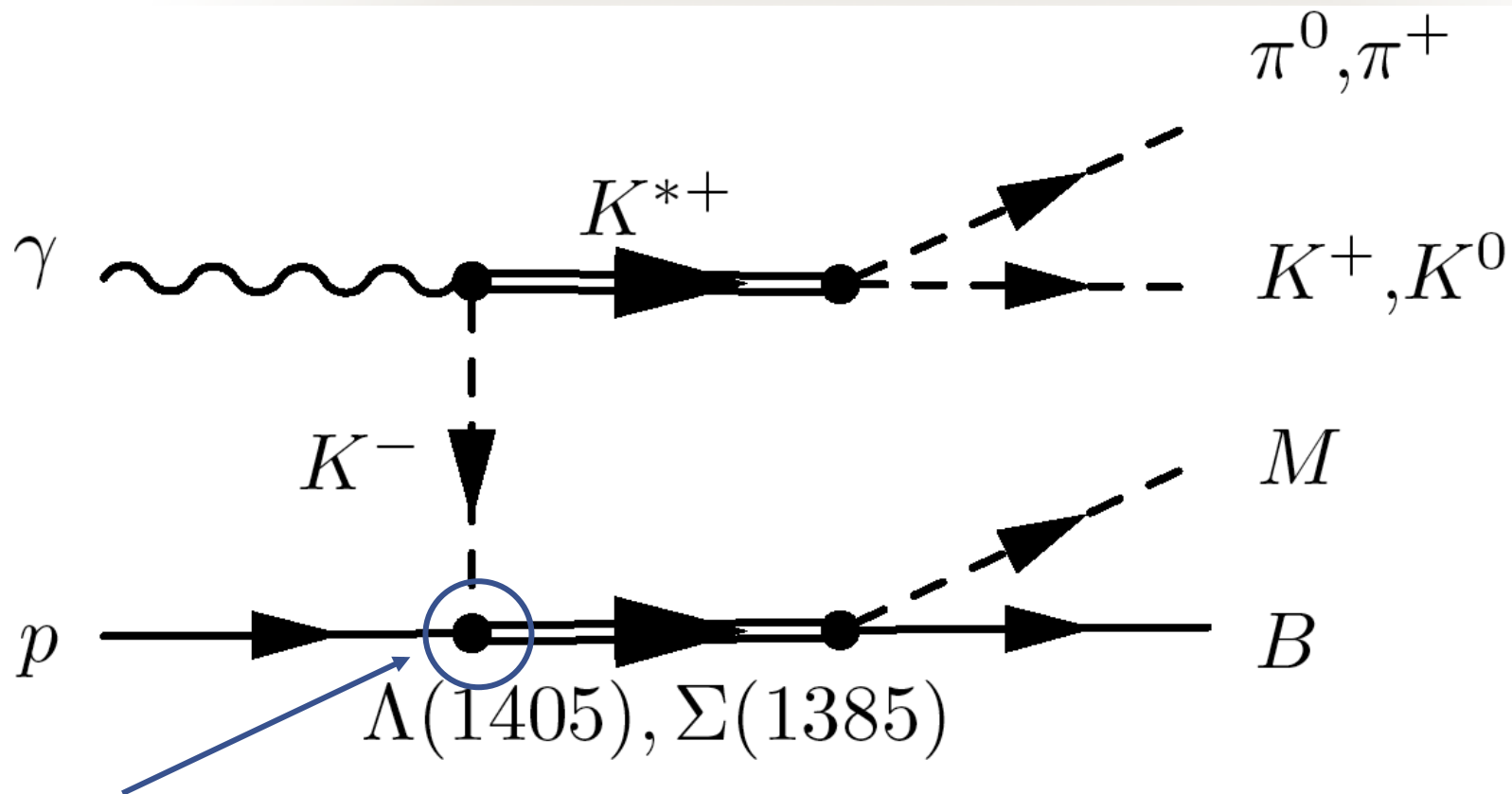
$\pi\Sigma$ mass distribution



$$\frac{d\sigma}{dM_I} = C |t_{\pi\Sigma \rightarrow \pi\Sigma}|^2 p_{CM} \quad \longrightarrow \quad \frac{d\sigma}{dM_I} = \left| \sum_i C_i t_{i \rightarrow \pi\Sigma} \right|^2 p_{CM}$$

D. Jido, et al., Nucl. Phys. A 723, 205 (2003)

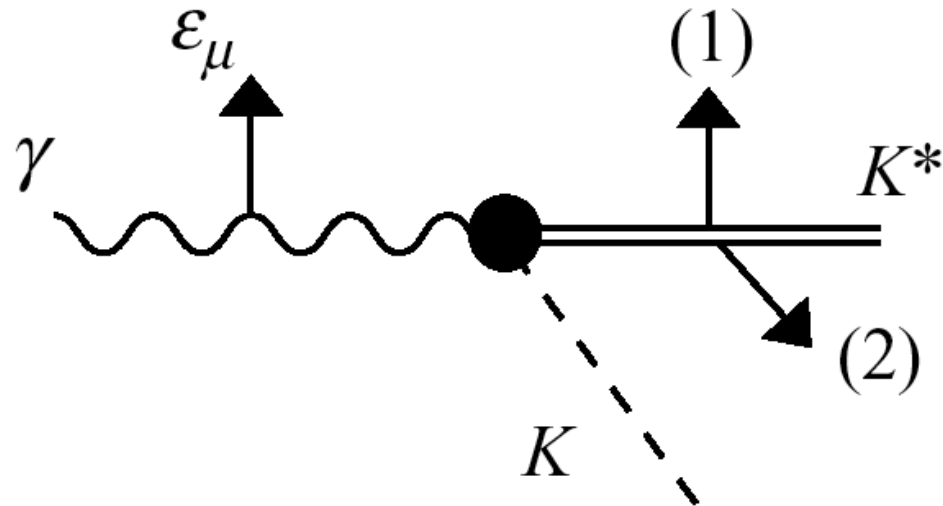
Photoproduction of $K^* \Lambda(1405)$



Only K^-p channel appears at the initial stage

Higher pole ??

Advantage of this reaction



$$(1) \quad \epsilon_\mu(K^*) \parallel \epsilon_\mu(\gamma) : J^P = \text{natural}$$

$$(2) \quad \epsilon_\mu(K^*) \perp \epsilon_\mu(\gamma) : J^P = \text{unnatural}$$

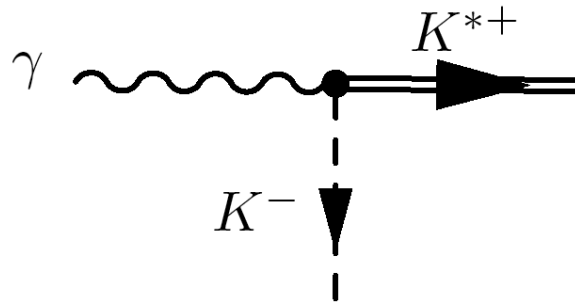
With polarized photon beam, the exchanged particle can be identified.

Clear mechanism

Effective interactions for vector meson

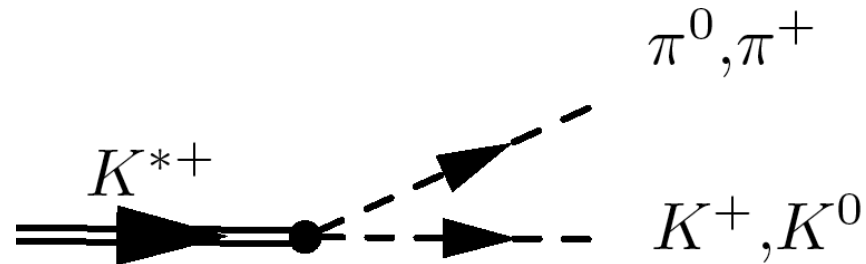
1. γVP coupling

$$\mathcal{L}_{K^*K\gamma} = g_{K^*K\gamma} \epsilon^{\mu\nu\alpha\beta} \partial_\mu A_\nu (\partial_\alpha K_\beta^{*-} K^+ + \partial_\alpha \bar{K}_\beta^{*0} K^0) + \text{h.c.}$$

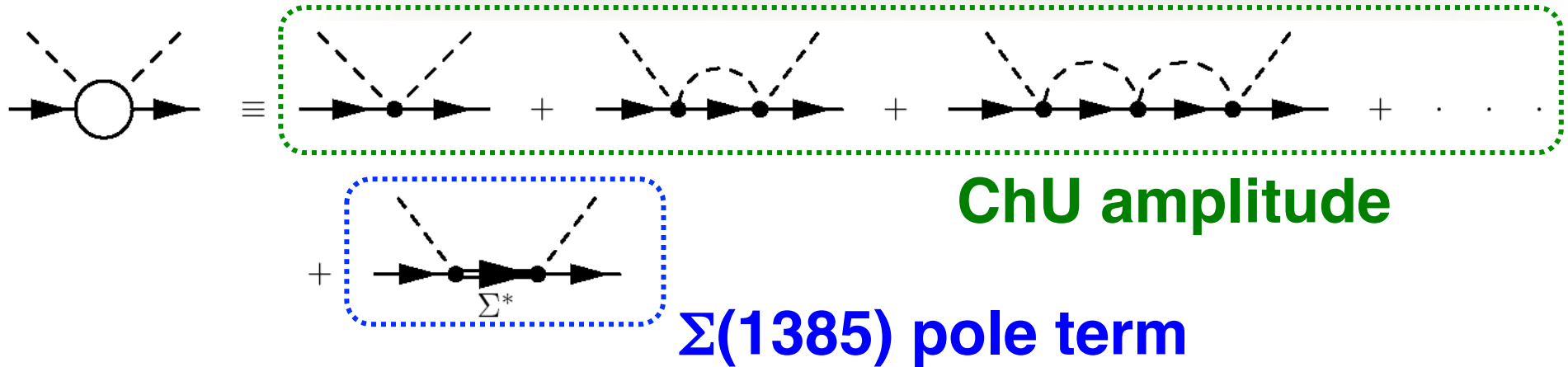


2. VPP coupling

$$\mathcal{L}_{VPP} = -\frac{ig_{VPP}}{\sqrt{2}} \text{Tr}(V^\mu [\partial_\mu P, P])$$



Effective interaction for baryon resonance



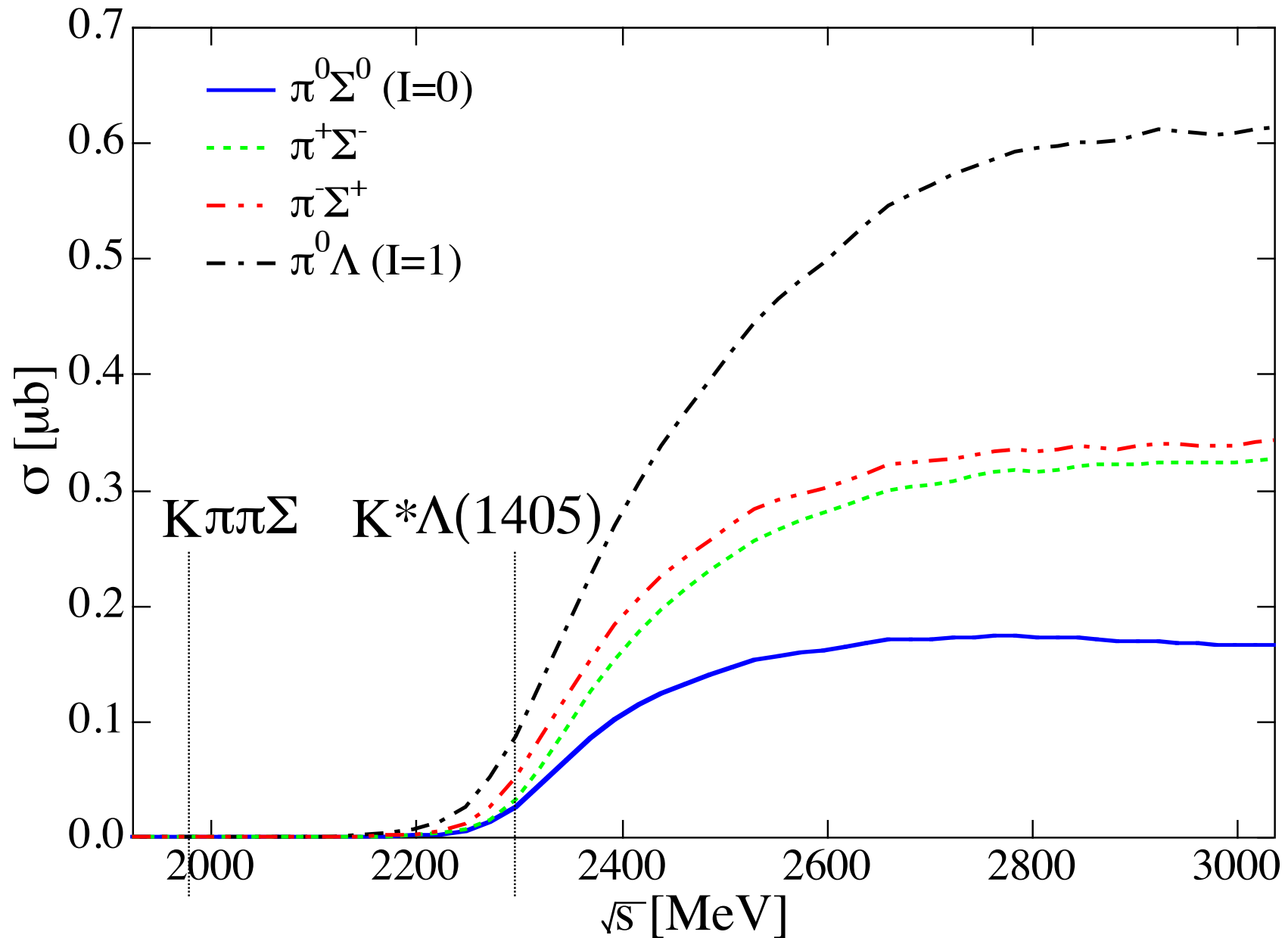
3. $\Sigma(1385)$ MB coupling

$$-it_{\Sigma^*i} = c_i \frac{12 D + F}{5} \frac{1}{2f} \mathbf{S} \cdot \mathbf{k}_i$$

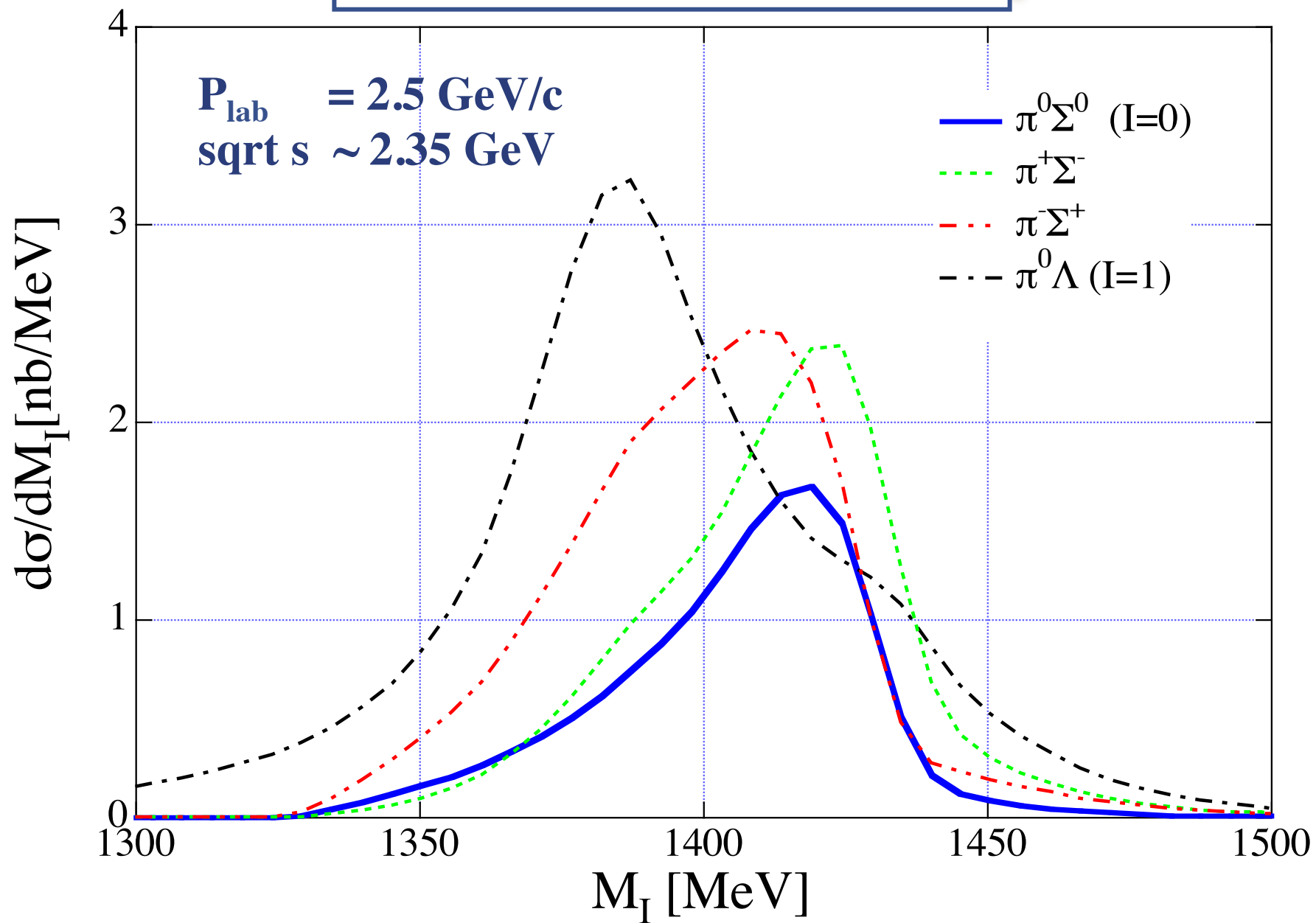
4. Form factor

$$F_f(k_1) = \frac{\Lambda^2 - m_K^2}{\Lambda^2 - (k_1)^2}$$

Total cross section



Invariant mass distribution



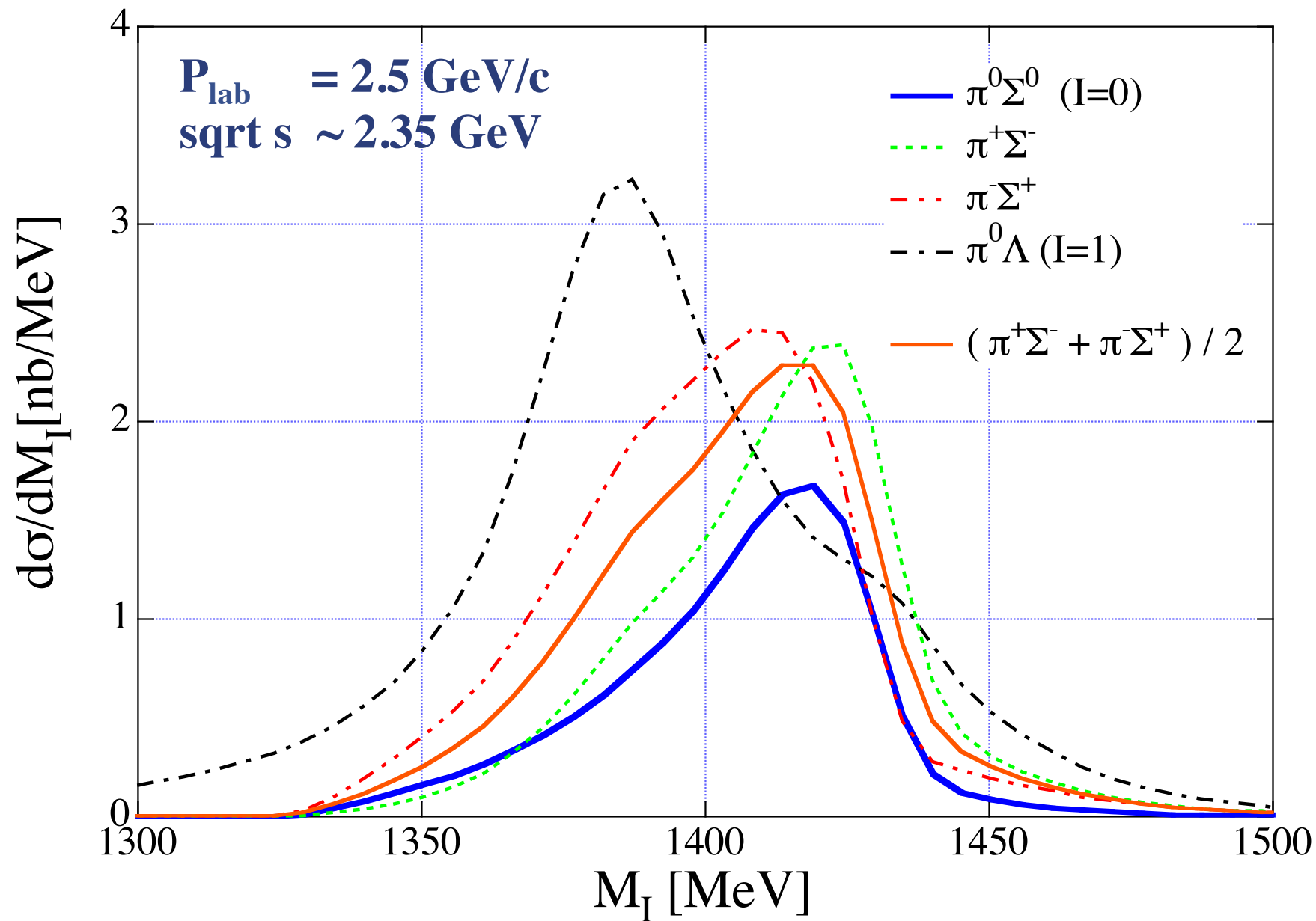
Isospin decomposition of $\pi\Sigma$ states

$$\frac{d\sigma(\pi^\pm\Sigma^\mp)}{dM_I} \propto \frac{1}{3}|T^{(0)}|^2 + \frac{1}{2}|T^{(1)}|^2 \pm \frac{2}{\sqrt{6}}\text{Re}(T^{(0)}T^{(1)*})$$

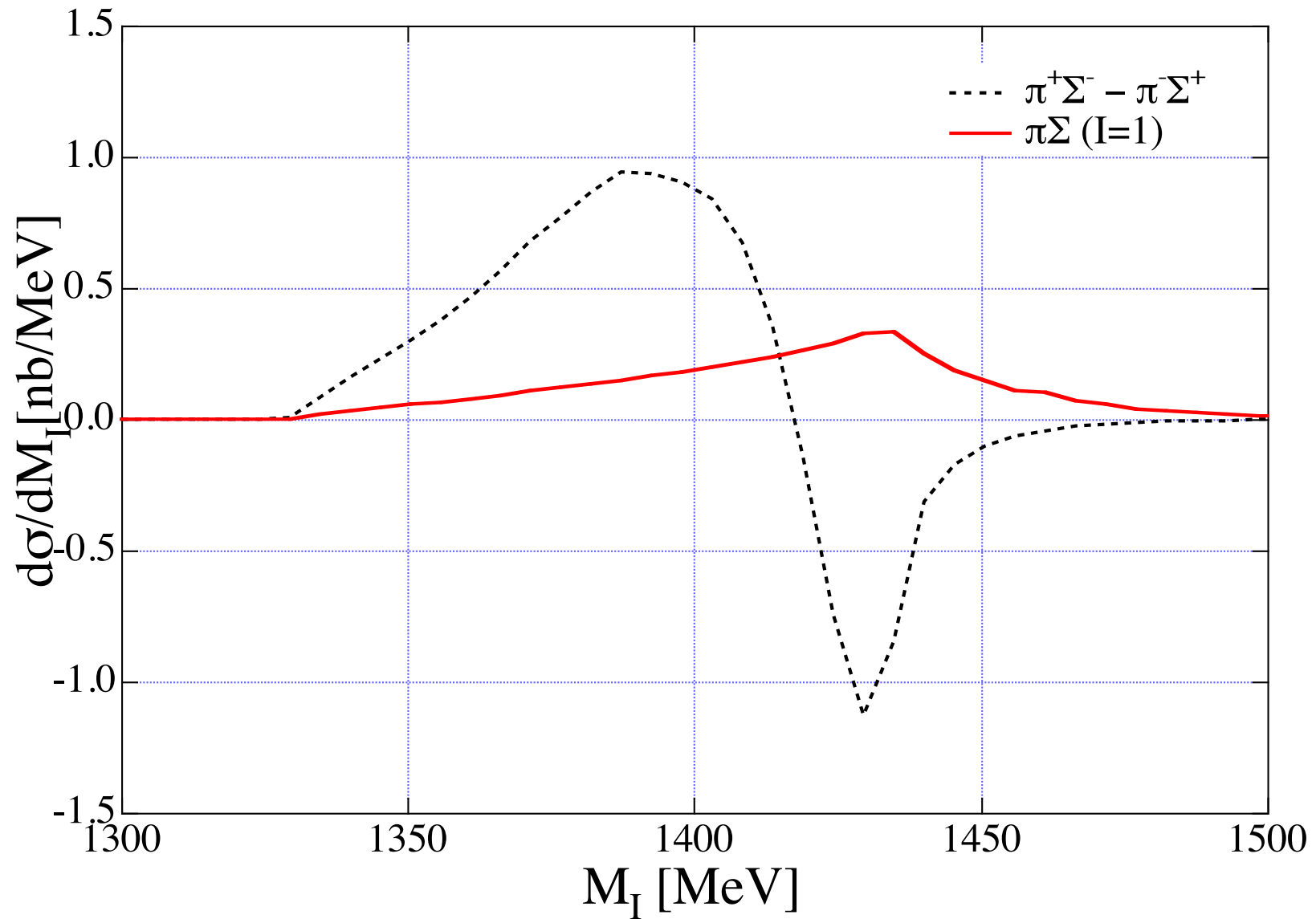
$$\frac{d\sigma(\pi^0\Sigma^0)}{dM_I} \propto \frac{1}{3}|T^{(0)}|^2$$

- **Difference among charged states**
-> when summed up, this term vanishes
- **No p-wave contribution**
-> **l=1 s-wave amplitude**

Invariant mass distribution 2



I=1, s-wave amplitude



Summary and conclusions 1

We study the **structure of $\Lambda(1405)$** using the chiral unitary model.

◆ There are **two poles** of the scattering amplitude around nominal $\Lambda(1405)$.

Pole 1 (1426+16i) : strongly couples to $\bar{K}N$ state

Pole 2 (1390+66i) : strongly couples to $\pi\Sigma$ state

◆ By observing the **charged $\pi\Sigma$ states** in the $\gamma p \rightarrow K^* \Lambda(1405)$ reaction, it is possible to isolate **higher energy pole.**

Summary and conclusions 2

- ◆ If we observe **neutral $\pi\Sigma$ state**, clear **$l=0$ distribution** is obtained.
- ◆ Combining three **$\pi\Sigma$ states**, we can also study the **s-wave $l=1$ amplitude**, where the existence of another pole is argued.

[T. Hyodo, et al, nucl-th/0401051](#)

<http://www.rcnp.osaka-u.ac.jp/~hyodo/>

Appendix : other processes

$\gamma p \rightarrow K^- \pi \Sigma$

J.C. Nacher, *et al.*, PLB445, 55(1999)

$K^- p \rightarrow \gamma \pi \Sigma$

J.C. Nacher, *et al.*, PLB461, 299(1999)

